#### **Finite-Element Methods and Numerical Linear Algebra**

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# Weighted Residual Methods

Idea:

- Assume that the solution can be represented in terms of analytic functions
- Express the approximate solution as a sum of such functions (rather than point-values)

Many subclasses of methods:

- finite-element methods (today)
- finite-volume methods (previous lecture)
- spectral methods (not covered here)
- boundary element methods (not covered here)

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## Weighted Residual Methods, cont'd

#### So far: finite-differences

- unknown function computed as a set of discrete nodal values
- differential formulation for each node
- Taylor series expansions on (structured) grids
- Increase accuracy by reducing local truncation error

#### Today: weighted-residual methods

- unknown function computed as a sum of continuous shape functions
- integral formulation of the equations
- minimize weighted residual for arbitrary control volume
- interpolation errors
- Increase accuracy by higherorder interpolation and optimized coefficients for minimum residuals

### **General Class of Problems**

Assume the PDE:  $\mathcal{L}(u(\mathbf{x})) = 0, \quad \mathbf{x} \in \Omega$ 

Example: Let us revisit the steady heat equation

$$-\nabla [K(x)\nabla u] = f(x), \qquad x \in \Omega$$
$$u(x) = g(x), \qquad x \in \partial \Omega$$

$$\longrightarrow \mathcal{L}(u(x)) = f(x) + \nabla [K(x)\nabla u]$$

Here

- K(x), f(x), and g(x) are known functions
- u(x) is the unknown function
- $\Omega$  is a domain with complete boundary  $\partial \Omega$

## Weighted Residual Methods, cont'd

Seek approximations of the form

$$\hat{u} = \sum_{j=1}^{M} u_j N_j(\mathbf{x})$$

where  $N_j(\mathbf{x})$  are prescribed functions and  $u_j$  are unknown coefficients. The  $N_j$ 's are called *basis functions* or *trial functions* 

An approximate solution  $\hat{u}$  should minimize the error  $u - \hat{u}$ :

- For special problems, minimize  $\|u \hat{u}\|$
- Not possible in general, since *u* is unknown

## How to Compute the Coefficients $u_j$ ?

In general, to determine  $u_j$  we must minimize the residual

 $R(u_1,\ldots,u_M;\mathbf{x})=\mathcal{L}(\hat{u})$ 

i.e., minimize how far  $\hat{u}$  is from satisfying the equation  $\mathcal{L}(u) = 0$ 

 Galerkin-type methods: the weighted residual should disappear over Ω for linearly independent weights W<sub>i</sub> or weighting functions W<sub>i</sub>(x)

$$\int_{\Omega} RW_i \, d\Omega = 0, \quad i = 1, \dots, M$$

Notice that  $R(\hat{u}) = 0$  in the *weak sense* and that  $R(\hat{u}) \neq 0$  pointwise

## How to Compute the Coefficients $u_j$ , cont'd

• The least-squares method: minimize the average square residual  $\int_{\Omega} R^2 d\Omega$ 

$$\int_{\Omega} R \frac{\partial R}{\partial u_i} \, d\Omega = 0, \quad i = 1, \dots, M$$

i.e.,  $W_i = \partial R / \partial u_i$ .

• The collocation method: choose  $W_i = \delta(\mathbf{x} - \mathbf{x}^{[i]})$ , where  $\delta$  is the Dirac delta function,

$$R(u_1, \dots, u_M; \mathbf{x}^{[i]}) = 0, \quad i = 1, \dots, M.$$

• The subdomain collocation method: decompose  $\Omega$  into M subdomains,  $\Omega = \bigcup_{\ell=1}^{M} \Omega_{\ell}$  (equivalent to a finite-volume method)

$$\int_{\Omega_i} L(\hat{u}) d\Omega = 0, \quad i = 1, \dots, M.$$

### **Example: 1D Poisson equation**

$$\mathcal{L}(u) = u''(x) + f(x) = 0, \qquad x \in [0, 1]$$

Discretization:

$$u(x) \approx \hat{u}(x) = \sum_{j=1}^{M} u_j N_j(x)$$

The residual:

$$R(\hat{u}(x)) = f(x) + \sum_{j=1}^{M} u_j N_j''(x)$$

Using the least-squares approach:

$$\frac{\partial R}{\partial u_i} = \sum_{j=1}^M \frac{\partial u_j}{\partial u_i} N_j''(x) = N_i''(x)$$

## Example, cont'd

The system of equations becomes

$$\int_0^1 \left( f(x) + \sum_{j=1}^M u_j N_j(x) \right) N_i''(x) \, dx = 0$$

$$-\sum_{j=1}^{M} \left( \int_{0}^{1} N_{i}''(x) N_{j}''(x) \, dx \right) u_{j} = \int_{0}^{1} f(x) N_{i}''(x) \, dx$$

 $\longrightarrow$  linear system of equations  $\mathbf{A}\mathbf{u}=\mathbf{b},$  where

- $A_{i,j} = \int_0^1 N_i''(x) N_j''(x) dx$
- $b_i = \int_0^1 f(x) N_i''(x) \, dx$

## Example, cont'd

Using the Galerkin approach:

$$-\sum_{j=1}^{M} \left( \int_{0}^{1} W_{i}(x) N_{j}''(x) \, dx \right) u_{j} = \int_{0}^{1} f(x) W_{i}(x) \, dx$$

Again a system of equations Au = b.

Two type of methods:

- $W_i = N_i$ , Galerkin method
- $W_i \neq N_i$ , *Petrov–Galerkin* method

## **How to Choose Test Functions?**

For simplicity, assume that u = 0 on  $\partial \Omega$ .

The functions  $N_k(x)$  can in principle be choosen almost arbitrarily:

- power series:  $N_k(x) = x^k$
- fourier series:  $N_k(x) = \{\sin(kx), \cos(kx)\}$
- Lagrange, Hermite, Chebychev polynomials
- •

In general, to get a well-behaved method we require that:

- $N_k = 0$  on the boundary
- $N_k$  almost orthogonal, to avoid numerical instabilities

## **The Finite-Element Method (FEM)**

*Finite elements* is a way of fulfilling these two requirements:

- divide the domain into non-overlapping *elements*
- let  $N_k$  be a simple polynomial over each element
- the global  $N_k$  is a picewise polynomial that vanishes except on a local patch of elements

Features:

- A very flexible approach
- Straightforward handling of complicated geometries
- Easy to construct higher-order approximations
- A broad spectrum of applications
- An engineering method
- Has a strong mathematical foundation

## **Piecewise Polynomial Basis Functions**

Define elements  $\Omega_e$  and nodes  $x^{[i]}$ 

The  $N_i$ 's have the properties:

- N<sub>i</sub> is a polynomial over each element, determined uniquely by the *nodal values*
- $N_i(x^{[j]}) = \delta_{i,j}$ , i.e., 1 if i = j and zero otherwise
- Hence,  $\hat{u}(x^{[i]}) = \sum_{j} u_{j} N_{j}(x^{[i]}) = u_{i}$

## **Examples of Basis Functions**

#### Linear basis functions



Each element has two nodes

#### Quadratic basis functions



Each element has three nodes

## **Essential Boundary Conditions**

Boundary-value problem

 $-u'' = f, x \in (0,1), \quad u(0) = u_L, u(1) = u_R$ 

Expansion with  $u_i = \hat{u}(x^{[i]})$ :

$$\hat{u}(x) = \psi(x) + \sum_{j=2}^{n-1} \hat{u}_j N_j(x), \quad \psi(x) = u_L N_1(x) + u_R N_n(x)$$

Alternative: skip  $\psi$  and enforce  $a_1 = u_L$  and  $a_n = u_R$  directly in the linear system

This is a general procedure

### **A Worked Example with Linear Elements**

Boundary-value problem

$$-u'' = f, x \in (0,1), \quad u(0) = u_L, u(1) = u_R$$

Galerkin's method (using integration by parts):

$$\sum_{j=1}^{n} A_{i,j} u_j = b_i, \quad i = 1, \dots n$$

where

$$A_{i,j} = \int_{0}^{1} N'_{i}(x)N'_{j}(x)dx, \quad b_{i} = \int_{0}^{1} f(x)N_{i}(x)dx$$

## Worked Example, cont'd

Observation:  $N_i(x)$  and  $N'_i(x)$  vanish over large parts of the domain ("nearly" orthogonal functions)



## **Example: Evaluation of the Coefficients**

**Direct computations:** 

$$A_{i,i-1} = \int_0^1 N'_{i-1} N'_i dx = -\frac{1}{h} \qquad A_{1,1} = A_{n,n} = \frac{1}{h}$$
$$A_{i,i} = \int_0^1 N'_i N'_i dx = \frac{2}{h}, \qquad A_{1,2} = A_{n,n-1} = -\frac{1}{h}$$
$$A_{i,i+1} = \int_0^1 N'_i N'_{i+1} dx = -\frac{1}{h} \qquad b_i = \int_0^1 f(x) N_i(x) dx$$

Numerical integration by trapezoidal rule:

$$\int_0^1 f(x)N_i(x) \, dx \approx \frac{1}{2} f(x^{[1]})N_i(x^{[1]})h + \sum_{j=1}^n f(x^{[j]})N_i(x^{[j]})h + \frac{1}{2} f(x^{[n]})N_i(x^{[n]})h$$
$$= 0 + \dots + 0 + f(x^{[i]})h + 0 + \dots + 0$$

## **Example: The Linear Equations**

Replace eq. no. 1 and *n* by boundary conditions The linear system:

$$u_{1} = u_{L},$$
  
$$-\frac{1}{h}u_{i-1} + \frac{2}{h}u_{i} - \frac{1}{h}u_{i+1} = f(x^{[i]})h, \quad i = 2, \dots, n-1,$$
  
$$u_{n} = u_{R}$$

Same result as from the finite difference method!

Exact or more accurate numerical integration: different right-hand side term

## **Element by Element Computations**

Split integral into a sum over each element:

$$A_{i,j} = \int_{0}^{1} N'_{i} N'_{j} dx = \sum_{e=1}^{m} A^{(e)}_{i,j}, \qquad A^{(e)}_{i,j} = \int_{\Omega_{e}} N'_{i} N'_{j} dx$$
$$b_{i} = \int_{0}^{1} f N_{i} dx = \sum_{e=1}^{m} b^{(e)}_{i}, \qquad b^{(e)}_{i} = \int_{\Omega_{e}} f N_{i} dx$$

 $A_{i,j}^{(e)} \neq 0$  iff *i* and *j* are nodes in element *e*  $b_i^{(e)} \neq 0$  iff *i* is node in element *e* 

## **Element by Element Computations, cont'd**

Collect nonzero  $A_{i,j}^{(e)}$  in a  $2 \times 2$  <u>elemental matrix</u> (for piecewise linear elements):

$$\tilde{A}_{r,s}^{(e)}, \quad r,s=1,2$$

r,s: local node numbers

Similar strategy for  $b_i^{(e)}$  give the elemental vector  $\tilde{b}_r^{(e)}$ 

Algorithm:

- compute all  $\tilde{A}_{r,s}^{(e)}$  and  $\tilde{b}_{r}^{(e)}$ ,
- combine them to a linear system

# **Local Coordinates**

- Map element  $\Omega_e = [x^{[e]}, x^{[e+1]}]$  to [-1,1]
- Define  $N_i$  in local  $\xi$  coordinates
- Perform all computations in local coordinates



- Local node r (=1,2) in element e corresponds to global node i = q(e, r)
- Local linear basis functions:

$$\tilde{N}_1(\xi) = \frac{1}{2}(1-\xi), \quad \tilde{N}_2(\xi) = \frac{1}{2}(1+\xi)$$

## Local Coordinates, cont'd

- Jacobian matrix of mapping: J
- Change integration variable from x to  $\xi$ :

$$\int_{x^{[e]}}^{x^{[e+1]}} N'_i(x) N'_j(x) \, dx = \int_{-1}^{1} J^{-1} \frac{d\tilde{N}_r(\xi)}{d\xi} J^{-1} \frac{d\tilde{N}_s(\xi)}{d\xi} \det J \, d\xi$$

- Uniform partition in 1D:  $J = \{h/2\}$
- We often write

$$\int_{\Omega_e} \frac{d\tilde{N}_r}{dx} \frac{d\tilde{N}_s}{dx} \det J \, d\xi$$

as the expression in local coordinates, knowing that

$$\frac{d\tilde{N}_r}{dx} = J^{-1}\frac{d\tilde{N}}{d\xi} = \frac{2}{h}\frac{d\tilde{N}}{d\xi}$$

## **1D Poisson Equation, revisited**

• 
$$-u''(x) = f(x)$$

• Elemental matrix and vector:

$$\tilde{A}_{r,s}^{(e)} = \int_{-1}^{1} \frac{2}{h} N_r'(\xi) \frac{2}{h} N_s'(\xi) \frac{h}{2} d\xi$$
$$\tilde{b}_r^{(e)} = \int_{-1}^{1} f(x^{(e)}(\xi)) \tilde{N}_r(\xi) \frac{h}{2} d\xi$$

• Example: 
$$r = s = 1$$
,

$$\tilde{A}_{1,1}^{(e)} = \frac{2}{h} \int_{-1}^{1} (-\frac{1}{2})(-\frac{1}{2}) \, d\xi = \frac{1}{h}$$

## **1D Poisson Equation, cont'd**

• Elemental matrix and vector:

$$\left\{ \tilde{A}_{r,s}^{(e)} \right\} = \frac{1}{h} \begin{pmatrix} 1 & -1 \\ -1 & 1 \end{pmatrix}$$
$$\left\{ \tilde{b}_{r}^{(e)} \right\} = \frac{h}{2} \begin{pmatrix} f(x^{(e)}(-1)) \\ f(x^{(e)}(1)) \end{pmatrix}$$

where numerical integration is used:

$$\int_{-1}^{1} g(\xi) \, d\xi \approx g(-1) + g(1)$$

# **Numerical Integration**

Integration rules are normally tabulated for integrals on [-1, 1]:

$$\int_{-1}^{1} g(\xi) d\xi \approx \sum_{k=1}^{n_I} g(\xi_k) w_k$$

 $\xi_k$ : integration points,  $w_k$ : integration weights

Some rules:

name	$n_I$	p	weights	points
Gauss-Legendre	1	1	(1)	(0)
Gauss-Legendre	2	3	(1,1)	$(-1/\sqrt{3}, 1/\sqrt{3})$
Gauss-Legendre	3	5	(5/9, 8/9, 5/9)	$(-\sqrt{3/5}, 0, \sqrt{3/5})$
Gauss-Lobatto	2	2	(1,1)	(-1, 1)
Gauss-Lobatto	3	3	(1/3, 4/3, 1/3)	(-1,0,1)

The rules integrate polynomials of degree p exactly





Elemental matrices and vectors must be assembled in the global system. Essential: local  $\rightarrow$  global mapping, q(e, r) Algorithm:

$$A_{q(e,r),q(e,s)} := A_{q(e,r),q(e,s)} + \tilde{A}_{r,s}^{(e)}, \quad r,s = 1,2$$

$$b_{q(e,r)} := b_{q(e,r)} + \tilde{b}_r^{(e)}, \quad r = 1, 2$$

# **Summing up the Procedures**

- Weighted residual formulation, often Galerkin's choice with  $W_i = N_i$
- Integration by parts
- Derivative boundary conditions in boundary terms
- Compute <u>elemental</u> matrices and vectors
  - Local coordinates with local numbering
  - Numerical integration
  - Enforce essential boundary conditions
  - Assemble local contributions
- Solve linear system

#### The Elementwise Algorithm:

initialize global linear system: set  $A_{i,j} = 0$  for i, j = 1, ..., n,  $b_i = 0$  for i = 1, ..., nloop over all elements: for  $e = 1, \ldots, m$ set  $\tilde{A}_{r,s}^{(e)} = 0, r, s = 1, \dots, n_e$ , set  $\tilde{b}_r^{(e)} = 0, r = 1, \dots, n_e$ loop over numerical integration points: for  $k = 1, ..., n_I$ evaluate  $\tilde{N}_r(\xi_k)$ , derivatives of  $\tilde{N}_r$  wrt.  $\xi$  and x, Jcontribution to elemental matrix and vector from the current integration point for  $r = 1, \ldots, n_e$ for  $s = 1, \ldots, n_e$  $\tilde{A}_{r,s}^{(e)} := \tilde{A}_{r,s}^{(e)} + \frac{d\tilde{N}_r}{dr} \frac{\tilde{N}_s}{dr} \det J w_k$  $\tilde{b}_{r}^{(e)} := \tilde{b}^{(e)} + f(x^{(e)}(\xi_{k}))N_{r} \det Jw_{k}$ incorporate essential boundary conditions: for  $r = 1, ..., n_e$ if node r has an essential boundary condition then modify  $\tilde{A}_{r,s}^{(e)}$  and  $\tilde{b}_{r}^{(e)}$  due to this condition assemble: for  $r = 1, ..., n_e$ for  $s = 1, ..., n_{e}$  $A_{q(e,r),q(e,s)} := A_{q(e,r),q(e,s)} + \tilde{A}_{r,s}^{(e)}$  $b_{q(e,r)} := b_{q(e,r)} + \tilde{b}_{r}^{(e)}$ 

## **2D Domains**

Strength of the finite element method: easy to work with geometrically complicated domains



#### Lake Superior with 6 islands, 2330 triangles