



# Bandgap References and Discrete Time Signals (chapter 8 + 9)

Tuesday 9th of February, 2010

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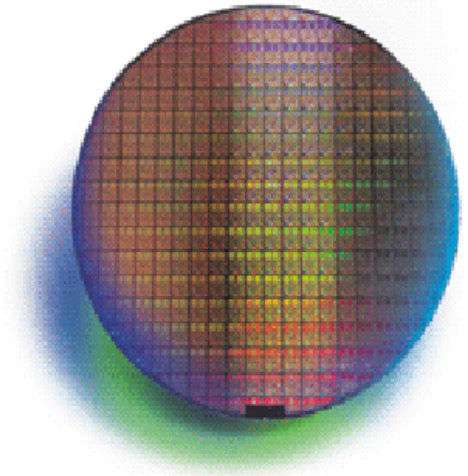
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Last time – Tuesday 2nd of February, and today, February the 9th:

- 8.1 performance of Sample-and-Hold Circuits
- 8.2 MOS Sample-and-Hold circuits
- 8.3 Examples of CMOS S/H circuits
- 8.5 Bandgap Voltage Reference Basics
- 8.6 Circuits for Bandgap References
- Chapter 9 Discrete-Time Signals
- 9.1 Overview of some signal spectra
- 9.2 Laplace Transforms of Discrete-Time Signals
- 9.3 Z-transform



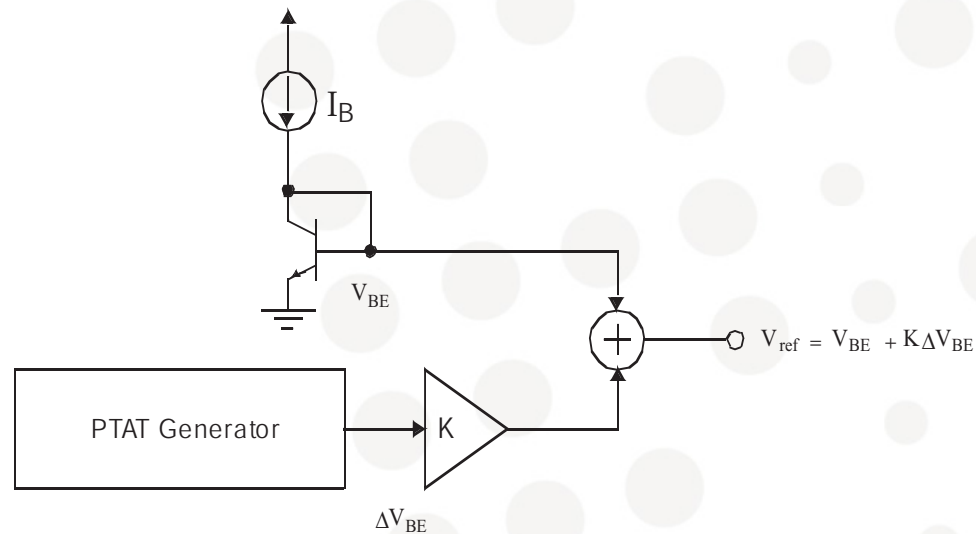
# Voltage references (chapter 8.5)

- Purpose:
  - Generate a **constant on-chip voltage** which is independent of temperature, supply voltage, aging etc.
- Different approaches:
  - 1) **Breakdown** voltage of a reverse-biased **zener** diode
    - **Too high voltage** for CMOS
  - 2) Threshold voltage difference between CMOS enhancement and **depletion transistors**
    - Depletion-mode transistors **unavailable** in most CMOS processes
  - 3) **Bandgap references**: Canceling the negative temperature dependence of a forward-biased pn-junction (CTAT) with a positive temperature dependence (PTAT) (proportional-to-absolute-temperature) circuit
    - Most commonly used
    - CTAT: Conversely proportional to temperature
    - PTAT: Proportional to temperature

# More about today's Bandgap Reference agenda (including, but not limited to):

- About the fundamental equations giving the relationship between the output voltage of a bandgap reference and temperature.
- How to design a bandgap reference for a "most stable" reference voltage at a particular temperature.
- How to estimate temperature dependence at another temperature that the BG reference was designed for.
- Practical implementations

# Basic principle



- The voltage  $V_{BE}$  is CTAT
- The voltage  $\Delta V_{BE}$  is PTAT
- $\Delta V_{BE}$  is scaled by  $K$  to get the same slope as  $V_{BE}$
- By adding  $V_{BE}$  and  $K \Delta V_{BE}$ , the output  $V_{ref}$  becomes independent of temperature



# Bandgap reference example

## A High Precision Curvature Compensated Bandgap Reference without Resistors

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 Institute of Microelectronics, Peking University, Beijing 100871, P. R. China  
 \* Email: chenjianghua@ime.pku.edu.cn

**A High Precision Curvature Compensated Bandgap Reference with**  
 Jianghua Chen; Xuewen Ni; Bangxian Mo; Zhanfei Wang;  
[Solid-State and Integrated Circuit Technology, 2006. ICSICT '06. 8th Inte](#)  
 23-26 Oct. 2006 Page(s):1748 - 1750  
 Digital Object Identifier 10.1109/ICSICT.2006.306414  
[AbstractPlus](#) | Full Text: [PDF](#)(124 KB) [IEEE CNF](#)  
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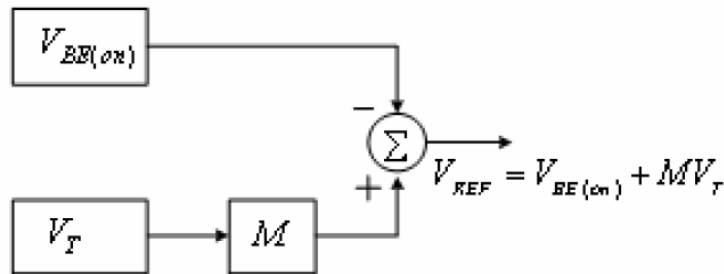


Figure 1 General bandgap reference architecture.

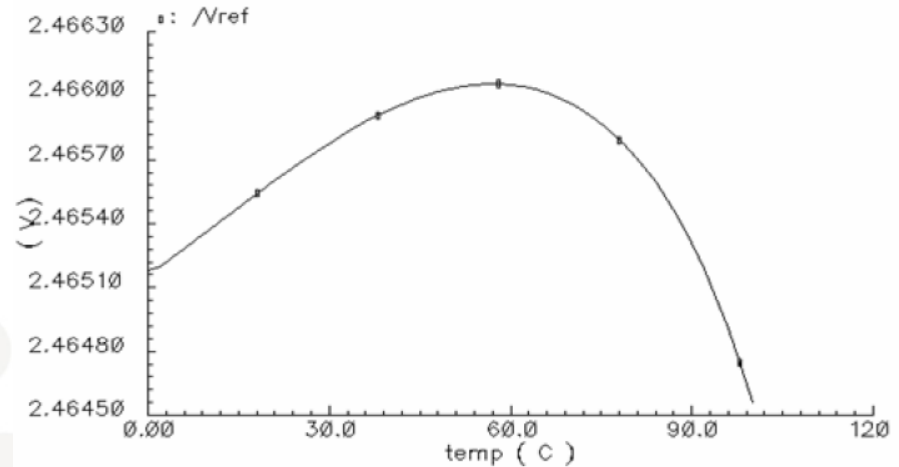


Figure 4 Output reference Vref vs. temperature.

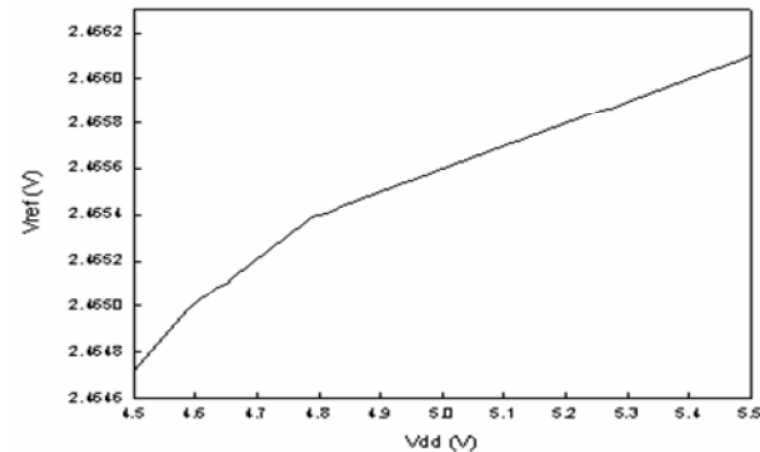


Figure 5 Output reference Vref vs. power supply Vdd.

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# Theory

- Collector current

$$I_C = I_s e^{V_{BE}/((kT)/q)}$$

- Solved with respect to  $V_{BE}$ :

$$V_{BE} = V_{G0} \left(1 - \frac{T}{T_0}\right) + V_{BE0} \frac{T}{T_0} + \frac{mkT}{q} \ln\left(\frac{T}{T_0}\right) + \frac{kT}{q} \ln\left(\frac{J_C}{J_{C0}}\right)$$

- The junction current equals the **effective area of the base-emitter junction** times the **junction current density,  $J_c$** :

-

$$I_C = A_E J_C$$

The difference between two base-emitter junctions biased at different densities (proportional to temperature):

$$\Delta V_{BE} = V_2 - V_1 = \frac{kT}{q} \ln\left(\frac{J_2}{J_1}\right)$$

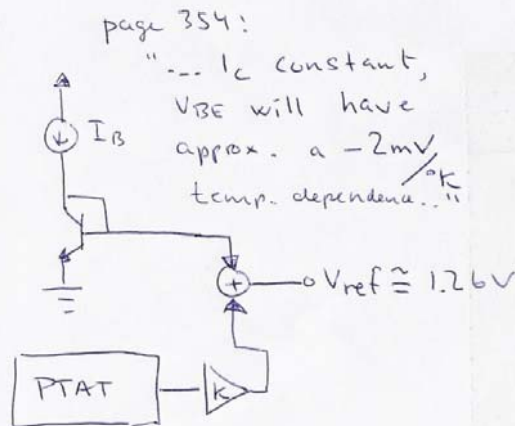
# Example 8.3

Ex. 8.3 Assume two trans. biased at current-density ratio of 10:1 at 300 °K. What is the difference in their base-emitter voltages and what is its temperature dependence?

likn. 8.12 : 
$$\Delta V_{BE} = \frac{kT}{q} \ln\left(\frac{J_2}{J_1}\right)$$

← eq. 8.12  
p. 354

$$= \frac{1.38 \cdot 10^{-23} (300)}{1.602 \cdot 10^{-19}} \ln(10) = \underline{\underline{59.5 \text{ mV}}}$$



$$V_{ref} = V_{BE2} + K \Delta V_{BE}$$

Since this voltage is proportional to absolute temperature, after a  $1^\circ\text{K}$  temp. increase, the voltage difference will be  $\Delta V_{BE} = 59.5 \text{ mV} \cdot \frac{301}{300} = 59.7 \text{ mV}$

Thus, the voltage dependence is  $59.5 \text{ mV}/300^\circ\text{K}$  or  $0.198 \text{ mV}/^\circ\text{K}$ .

Since the temperature dependence of a single  $V_{BE}$  is  $-2 \text{ mV}/^\circ\text{K}$ , if it is desired to cancel the temp. dependence of a single  $V_{BE}$  then  $\Delta V_{BE}$  should be amplified by about a factor of 10.



# Theory

- Assuming that:

$$\frac{J_i}{J_{i0}} = \frac{T}{T_0}$$

- $V_{\text{ref}}$  can then be written as:

$$\begin{aligned} V_{\text{ref}} &= V_{\text{BE2}} + K \Delta V_{\text{BE}} \\ &= V_{\text{G0}} + \frac{T}{T_0} (V_{\text{BE0-2}} - V_{\text{G0}}) + (m-1) \frac{kT}{q} \ln\left(\frac{T_0}{T}\right) + K \frac{kT}{q} \ln\left(\frac{J_2}{J_1}\right) \end{aligned}$$

- For a given temperature  $V_{\text{ref}}$  may be independent of changes in the temperature if a proper value of  $K$  is assigned
- This (equation 8.16) is the **fundamental equation giving the relationship between the output voltage of a bandgap voltage reference and temperature.**

# From $V_{BE}$ as a function of collector current and temperature to $V_{out}$ for BG ref. (part 1 of 2)

PP 354-355

$$V_{BE} = V_{G0} \left(1 - \frac{T}{T_0}\right) + V_{BE0} \frac{T}{T_0} + \frac{m k T}{q} \ln\left(\frac{T_0}{T}\right) + \frac{k T}{q} \ln\left(\frac{J_c}{J_{c0}}\right) \quad (8.10)$$

$$\uparrow \quad \text{(using } \frac{J_c}{J_{c0}} = \frac{T}{T_0} \text{)}$$

$$V_{BE} = V_{G0} - V_{G0} \frac{T}{T_0} + V_{BE0} \frac{T}{T_0} + \frac{m k T}{q} \ln\left(\frac{T_0}{T}\right) + \frac{k T}{q} \ln\left(\frac{T}{T_0}\right)$$

$$\uparrow$$

$$V_{BE} = V_{G0} + \frac{T}{T_0} (V_{BE0} - V_{G0}) + \frac{m k T}{q} [\ln T_0 - \ln T] + \frac{k T}{q} [\ln T - \ln T_0]$$

$$\uparrow$$

$$V_{BE} = V_{G0} + \frac{T}{T_0} (V_{BE0} - V_{G0}) + \frac{m k T}{q} \ln T_0 - \frac{m k T}{q} \ln T + \frac{k T}{q} \ln T - \frac{k T}{q} \ln T_0$$

$$\uparrow$$

$$V_{BE} = V_{G0} + \frac{T}{T_0} (V_{BE0} - V_{G0}) + \frac{m k T}{q} \ln T_0 - \frac{m k T}{q} \ln T - \left(\frac{k T}{q} \ln T_0 - \frac{k T}{q} \ln T\right) \quad \boxed{\ln\left(\frac{a}{b}\right) = \ln a - \ln b}$$

$$\uparrow$$

$$V_{BE} = V_{G0} + \frac{T}{T_0} (V_{BE0} - V_{G0}) + \frac{m k T}{q} \ln T_0 - \frac{m k T}{q} \ln T - \frac{k T}{q} \ln \frac{T_0}{T}$$

$$\uparrow$$

$$V_{BE} = V_{G0} + \frac{T}{T_0} (V_{BE0} - V_{G0}) + \frac{m k T}{q} \ln \frac{T_0}{T} - \frac{k T}{q} \ln \frac{T_0}{T}$$

$$\uparrow$$

$$V_{BE} = V_{G0} + \frac{T}{T_0} (V_{BE0} - V_{G0}) + (m-1) \frac{k T}{q} \ln \frac{T_0}{T}$$

OBS! "Unnecessary rewriting from 2nd to 3rd eq."

$\Rightarrow$  SIMILAR TO ALL BUT LAST PART OF EQ 8.16.

10

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From  $V_{BE}$  as a function of collector current and temperature to  $V_{out}$  for BG ref. (part 2 (of 2))

$$\begin{aligned} 8.16 \quad V_{ref} &= V_{BE2} + K \Delta V_{BE} \\ &= V_{60} + \frac{T}{T_0} (V_{BE0-2} - V_{60}) + (m-1) \frac{kT}{q} \ln\left(\frac{T_0}{T}\right) \quad (8.16) \\ &\quad + K \left[ \frac{kT}{q} \ln\left(\frac{J_2}{J_1}\right) \right] \quad (\text{using } 8.12) \end{aligned}$$

This equation (8.16) is the fundamental equation giving the relationship between the output voltage of a bandgap reference and temperature.

Differentiating eq. 8.16 with respect to temperature, getting eq. 8.17

$\left[ V_{G0} + \frac{T}{T_0} (V_{BE0-2} - V_{G0}) + (m-1) \frac{kT}{q} \ln \left( \frac{T_0}{T} \right) + K \frac{kT}{q} \ln \left( \frac{J_2}{J_1} \right) \right]'$ $= \cancel{(V_{G0})'} + \left( \frac{T}{T_0} \right)' (V_{BE0-2} - V_{G0}) + \cancel{\frac{T}{T_0} (V_{BE0-2} - V_{G0})'} + (m-1) \left[ \frac{kT}{q} \ln \left( \frac{T_0}{T} \right) \right]' + (m-1) \left[ \frac{kT}{q} \ln \left( \frac{T_0}{T} \right) \right]'$ $+ \left( K \frac{kT}{q} \right)' \ln \left( \frac{J_2}{J_1} \right) + K \frac{kT}{q} \left( \ln \left( \frac{J_2}{J_1} \right) \right)'$ $\frac{\partial V_{ref}}{\partial T} = \underbrace{\left( \frac{T}{T_0} \right)' (V_{BE0-2} - V_{G0}) + (m-1) \left[ \frac{kT}{q} \ln \left( \frac{T_0}{T} \right) \right]'}_A + \underbrace{\left( K \frac{kT}{q} \right)' \cdot \ln \left( \frac{J_2}{J_1} \right)}_B$	$y = u+v, y' = u'+v'$ $y = uv, y' = u'v + uv'$ $y = \ln x, y' = \frac{1}{x}$
---	--

Noen mellomregninger: A

$$\left( \frac{T}{T_0} \right)' = \left( T \cdot \frac{1}{T_0} \right)' = 1 \cdot \frac{1}{T_0} + T \cdot 0 = \frac{1}{T_0}$$

$$y = \ln \frac{T_0}{T} = \ln T_0 - \ln T$$

$$y' = (\ln T_0)' - (\ln T)' = -\frac{1}{T}$$

$$A = (m-1) \cdot \left( \frac{kT}{q} \right)' \cdot \ln \left( \frac{T_0}{T} \right) + \frac{kT}{q} \left( \ln \frac{T_0}{T} \right)' = (m-1) \frac{k}{q} \cdot \ln \left( \frac{T_0}{T} \right) + \frac{kT}{q} \left( -\frac{1}{T} \right) = (m-1) \frac{k}{q} \cdot \ln \left( \frac{T_0}{T} \right) - \frac{k}{q}$$

$$= \frac{(m-1) \frac{k}{q} \left( \ln \left( \frac{T_0}{T} \right) - 1 \right)}{}$$

$$B = \left[ \left( K \frac{k}{q} \right)' \cdot T + K \frac{k}{q} \cdot T' \right] \cdot \ln \left( \frac{J_2}{J_1} \right) = K \frac{k}{q} \cdot \ln \left( \frac{J_2}{J_1} \right)$$

∴  $\frac{\partial V_{ref}}{\partial T} = \frac{1}{T_0} (V_{BE0-2} - V_{G0}) + K \frac{k}{q} \cdot \ln \left( \frac{J_2}{J_1} \right) + (m-1) \frac{k}{q} \left[ \ln \left( \frac{T_0}{T} \right) - 1 \right]$  eq. 8.18

Setting equation 8.17 = 0, and  $T = T_0$  getting eq. 8.18, giving the needs for zero temperature dependence at the reference temp.

$$\text{Set } \frac{\partial V_{ref}}{\partial T} = 0 \quad \wedge \quad T = T_0$$

$$\ln 1 = 0$$

$$0 = \frac{1}{T_0} (V_{BE0-2} - V_{G0}) + K \frac{k}{q} \ln \left( \frac{J_2}{J_1} \right) + (m-1) \frac{k}{q} \left[ \ln \left( \frac{T_0}{T_0} \right) - 1 \right]$$

$$0 = \frac{1}{T_0} (V_{BE0-2} - V_{G0}) + K \frac{k}{q} \ln \left( \frac{J_2}{J_1} \right) + (m-1) \frac{k}{q} (-1)$$

$$0 = (V_{BE0-2} - V_{G0}) + K \frac{k T_0}{q} \ln \left( \frac{J_2}{J_1} \right) + (m-1) \frac{k T_0}{q} (-1)$$

$$V_{BE0-2} + K \frac{k T_0}{q} \ln \left( \frac{J_2}{J_1} \right) = V_{G0} - (m-1) \frac{k T_0}{q} (-1)$$

$$V_{BE0-2} + K \frac{k}{q} T_0 \ln \left( \frac{J_2}{J_1} \right) = V_{G0} + (m-1) \frac{k T_0}{q} \quad (8.18)$$



Setting  $T=T_0$  in eq. 8.16 gives the left side of eq. 8.18

$$\begin{aligned} \underline{\text{8.16}} : \quad V_{\text{ref}} &= V_{\text{BE}2} + K \Delta V_{\text{BE}} \\ &= V_{\text{G}0} + \frac{T}{T_0} (V_{\text{BE}0-2} - V_{\text{G}0}) + (m-1) \frac{kT}{q} \ln\left(\frac{T_0}{T}\right) + K \frac{kT}{q} \ln\left(\frac{J_2}{J_1}\right) \end{aligned}$$

Letting  $T=T_0$  :

$$\begin{aligned} V_{\text{ref}} &= V_{\text{G}0} + \frac{T_0}{T_0} (V_{\text{BE}0-2} - V_{\text{G}0}) + (m-1) \frac{kT_0}{q} \ln\left(\frac{T_0}{T_0}\right) + K \frac{kT_0}{q} \ln\left(\frac{J_2}{J_1}\right) \\ &= V_{\text{G}0} + (V_{\text{BE}0-2} - V_{\text{G}0}) + K \frac{kT_0}{q} \ln\left(\frac{J_2}{J_1}\right) \\ &= \underline{V_{\text{BE}0-2} + K \frac{kT_0}{q} \ln\left(\frac{J_2}{J_1}\right)} \end{aligned}$$

For  $T=T_0$ , the left side of eq. 8.18 equals the result above.

$$\underline{\text{8.18}} : \quad \underline{V_{\text{BE}0-2} + K \frac{kT_0}{q} \ln\left(\frac{J_2}{J_1}\right) = V_{\text{G}0} + (m-1) \frac{kT_0}{q}}$$

For zero temperature dependence at  $T=T_0$ . At **300 K** (8.18, 8.19, 8.20):

$$V_{BE0-2} + K \frac{kT_0}{q} \ln\left(\frac{J_2}{J_1}\right) = V_{G0} + (m-1) \frac{kT_0}{q} \quad (8.18) \quad \begin{aligned} [^{\circ}\text{C}] &= [\text{K}] - 273.15 \\ 300[\text{K}] &- 273.15[\text{K}] \\ &\approx 27^{\circ}\text{C} \end{aligned}$$

The left side of eq 8.18 is the output voltage  $V_{\text{ref}}$  at  $T=T_0$  (as we have shown).

For zero temperature dependence at  $T=T_0$ , we need

$$V_{\text{ref-0}} = V_{G0} + (m-1) \frac{kT_0}{q} \quad (8.19)$$

For the special case of  $T_0 = 300 \text{ K}$  and  $m = 2.3$  (8.19) implies that, for zero temperature dependence

$$\begin{aligned} V_{\text{ref-0}} = V_{G0} &= 1.206 \text{ V} + (2.3 - 1) \cdot \frac{1.38 \times 10^{-23} (300)}{1.602 \times 10^{-19}} \\ &= 1.206 \text{ V} + 1.3 \times 25.8 \text{ mV} \\ &= \underline{1.24 \text{ V}} \end{aligned}$$

Note that this value is independent of the current densities chosen.

## Required value for K at 300K (eq. 8.21):

From eq (8.20) we got  $V_{ref-0} = 1.24 \text{ V}$  for zero temp. dependence at  $300^\circ\text{K}$ . This value is independent of the current densities chosen.

K from equation 8.18:

$$K = \frac{V_{G0} + (m-1) \frac{kT_0}{q} - V_{BE0-2}}{\frac{kT_0}{q} \ln\left(\frac{J_2}{J_1}\right)} = \frac{1.24 \text{ V} - V_{BE0-2}}{25.8 \text{ mV} \cdot \ln\left(\frac{J_2}{J_1}\right)}$$

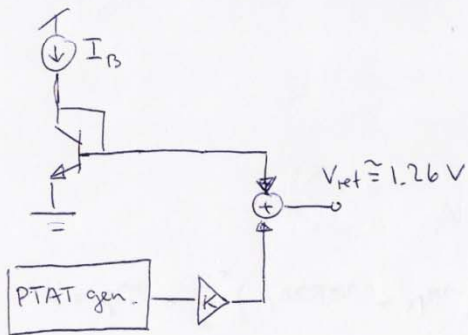
The output of a bandgap reference is given by the bandgap voltage,  $V_{G0}$  plus a small correction to account for 2nd-order effects.

## Output voltage for temperatures different from the reference; get (8.22) and then differentiate ...

The fundamental equation giving the relationship between the output voltage of a bandgap reference and temperature is equation 8.16, page 355:

$$V_{ref} = V_{BE2} + K \Delta V_{BE}$$

$$= V_{G0} + \frac{T}{T_0} (V_{BE0-2} - V_{G0}) + (m-1) \frac{kT}{q} \ln\left(\frac{T_0}{T}\right) + K \frac{kT}{q} \ln\left(\frac{J_2}{J_1}\right)$$



The output voltage of the reference for temperatures different from the reference is found after backsubstituting (8.18) and (8.19) into (8.16) and some manipulation:

$$V_{ref} = V_{G0} + (m-1) \frac{kT}{q} \left[ 1 + \ln\left(\frac{T_0}{T}\right) \right] \quad (8.22)$$

(8.22) differentiated with respect to T, getting (8.23):

Differentiating 8.22 :

$$\log_a \left( \frac{u}{v} \right) = \log_a u - \log_a v$$

$$y = uv \\ y' = u'v + uv'$$

$$y = \ln x \\ y' = \frac{1}{x}$$

$$\begin{aligned} \frac{\partial V_{ref}}{\partial T} &= (V_{ref})' + \underbrace{\left[ (m-1) \frac{kT}{q} \right]'}_{u'} \underbrace{\left[ 1 + \ln \left( \frac{T_0}{T} \right) \right]}_v + \underbrace{(m-1) \frac{kT}{q}}_u \underbrace{\left[ 1 + \ln \left( \frac{T_0}{T} \right) \right]'}_{v'} \\ &= (m-1) \frac{k}{q} \left[ 1 + \ln \left( \frac{T_0}{T} \right) \right] + (m-1) \frac{kT}{q} \cdot \left( \frac{-1}{T} \right) \\ &= \text{---} \parallel \text{---} + (-m+1) \frac{k}{q} \\ &= \text{---} \parallel \text{---} \div (m-1) \frac{k}{q} \\ &= (m-1) \frac{k}{q} \left[ \cancel{1 + \ln \left( \frac{T_0}{T} \right)} - \cancel{1} \right] = (m-1) \frac{k}{q} \ln \left( \frac{T_0}{T} \right) \quad (8.23) \end{aligned}$$

Equations 8.22 and 8.23 may be used to estimate the temperature dependence at temperatures different from the reference temperature



# Example 8.4

Ex. 8.4 Estimate the temperature dependence at  $0^\circ\text{C}$   
p. 357 for a bandgap reference that was designed  
to have zero temperature dependence at  $20^\circ\text{C}$ .  
Present the result as  $\text{ppm}/^\circ\text{K}$ .

Using eq. 8.23 : 
$$\frac{\partial V_{\text{ref}}}{\partial T} = (m-1) \frac{k}{q} \ln\left(\frac{T_0}{T}\right)$$

$0^\circ\text{K}$  corresponds to  $-273^\circ\text{C}$  :  $T_0 = 293^\circ\text{K}$  ,  $T = 273^\circ\text{K}$

$$\frac{\partial V_{\text{ref}}}{\partial T} = (2.3-1) \frac{1.38 \cdot 10^{-23}}{1.6 \cdot 10^{-19}} \ln\left(\frac{293}{273}\right) \left[\frac{\text{mV}}{^\circ\text{K}}\right] = \underline{8 \text{ mV}/^\circ\text{K}}$$

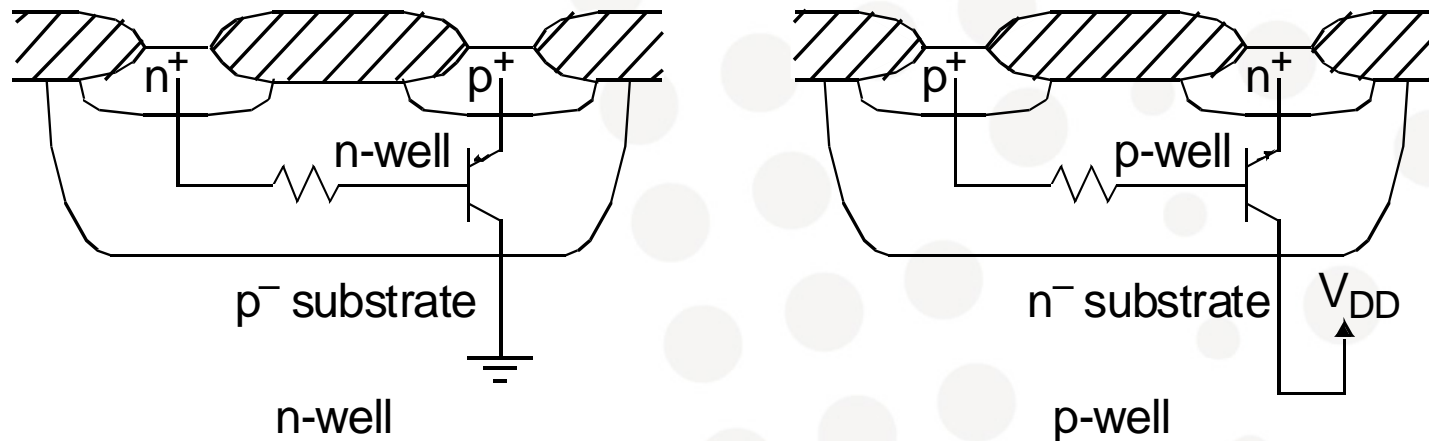
For a reference voltage of  $1.24 \text{ V}$ , a dependency of  
 $8 \text{ mV}/^\circ\text{K}$  results in  $\frac{8 \text{ mV}/^\circ\text{K}}{1.24 \text{ V}} = 6.5 \cdot 10^{-6} \text{ parts}/^\circ\text{K}$

$$= \underline{6.5 \text{ ppm}/^\circ\text{K}}$$

"ppm" : parts per million

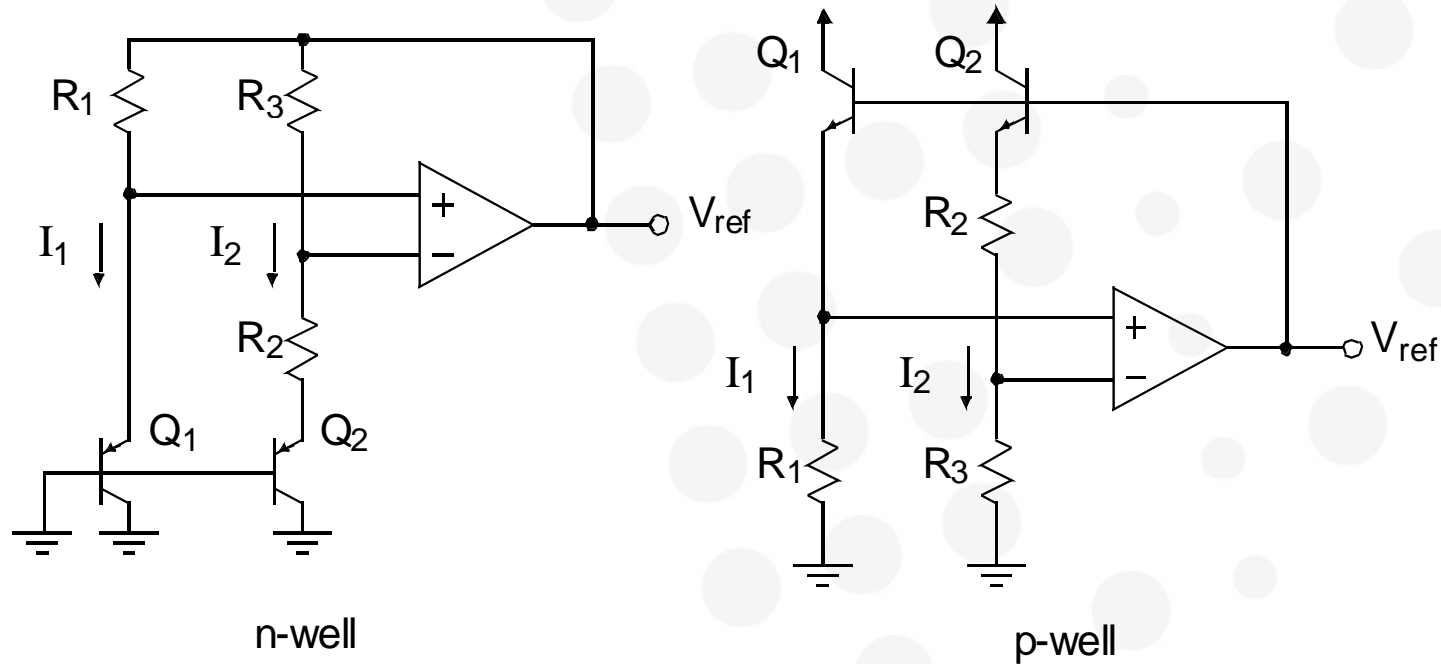
NB! Ideally:  $0 \text{ ppm}/^\circ\text{K}$  , typically  $4 \cdot 10 \times$  the  
value here.

# CMOS Bandgap References



- Vertical CMOS well transistors in an n-well and p-well process (pnp in n-well, npn in p-well)

# CMOS BGR Circuits



- CMOS bandgap references implemented with well transistors

# Design equations, BG ref.

## DESIGN EQUATIONS

The voltage drop due to  $I_1$ :

$$V_{ref} = V_{EB1} + V_{R1} \quad (8.35)$$

Since  $V_{R1} = V_{R3}$ , assuming an ideal OP-AMP:

$$V_{R2} = V_{EB1} - V_{EB2} = \Delta V_{EB} \quad (8.36)$$

$I_2$  runs through both  $R_3$  and  $R_2$ , so that

$$(8.37) \quad V_{R3} = \frac{R_3}{R_2} V_{R2} = \frac{R_3}{R_2} \cdot \Delta V_{EB} \quad (V_{R2} = \Delta V_{EB})$$

Substituting (8.37) into (8.35):

$$V_{ref} = V_{EB1} + V_{R1} = V_{EB1} + V_{R3} = V_{EB1} + \frac{R_3}{R_2} \Delta V_{EB} \quad (8.38)$$

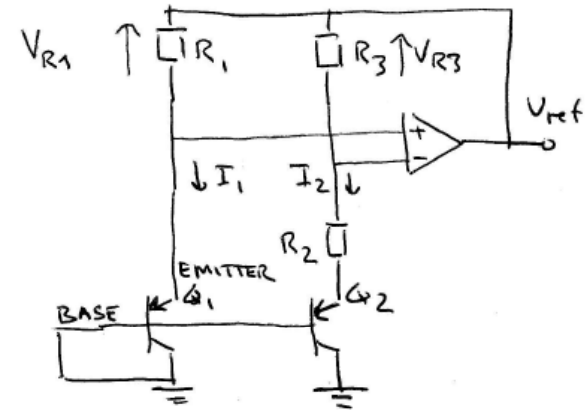
In integrated realizations the bipolar transistors are often taken the same size, and different current densities are realized by taking  $R_3 > R_1$ , causing  $I_1 > I_2$ . In this case

$$\frac{I_1}{I_2} = \frac{J_1}{J_2} = \frac{R_3}{R_1} \quad (8.39)$$

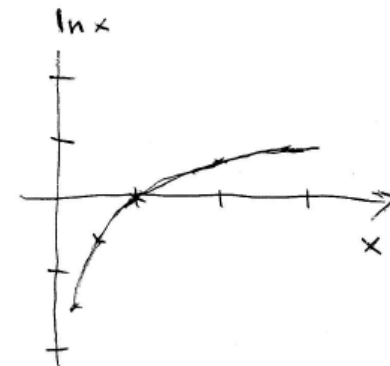
Recalling from 8.12 that  $\Delta V_{EB} = V_{EB1} - V_{EB2} = \frac{kT}{q} \ln \frac{J_1}{J_2} \quad (8.40)$

$$(8.39) \text{ into } (8.38): V_{ref} = V_{EB1} + \frac{R_3}{R_2} \frac{kT}{q} \ln \left( \frac{R_3}{R_1} \right) \quad K = \frac{R_3}{R_2} \quad (8.42)$$

Fig. 8.27



ln x	x	0,2	0,5	0	1	2	e
ln x	-1,61	-0,69	inv- alid	0	0,69	1	



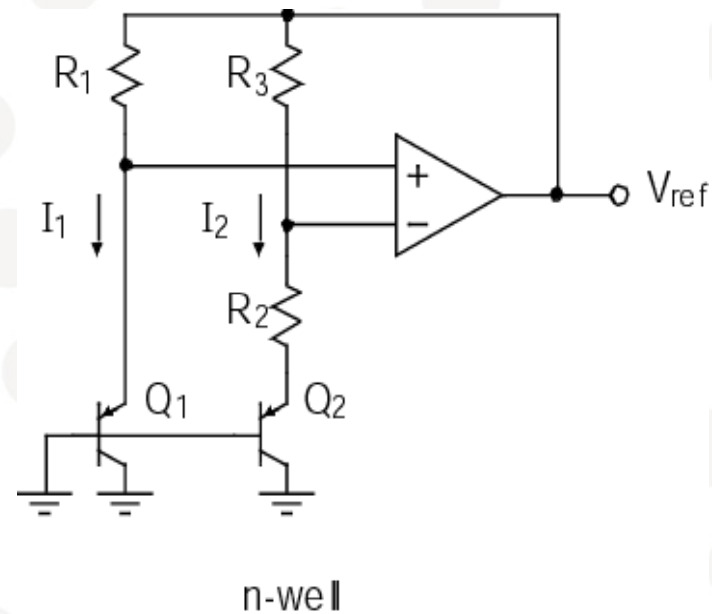
# Design Equations

$$V_{\text{ref}} = V_{\text{EB1}} + V_{\text{R1}}$$

$$V_{\text{R2}} = V_{\text{EB1}} - V_{\text{EB2}} = \Delta V_{\text{EB}}$$

$$V_{\text{R3}} = \frac{R_3}{R_2} V_{\text{R2}} = \frac{R_3}{R_2} \Delta V_{\text{EB}}$$

$$V_{\text{ref}} = V_{\text{EB1}} + \frac{R_3}{R_2} \Delta V_{\text{EB}}$$





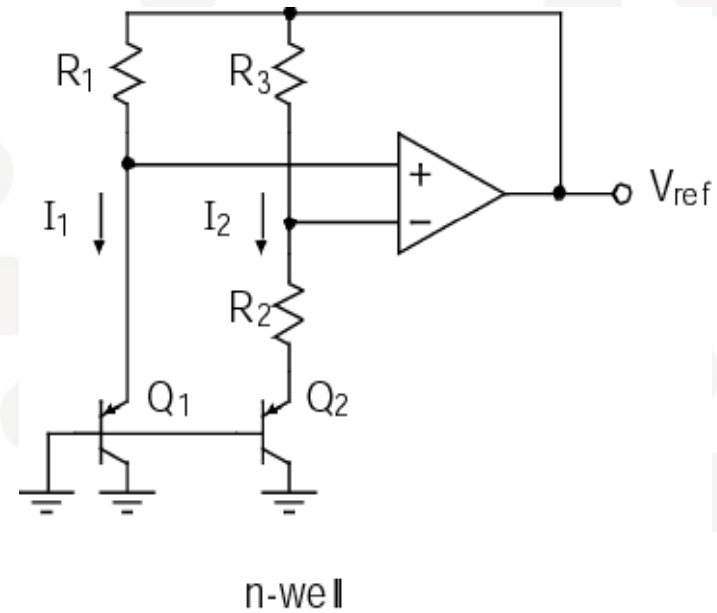
# Design Equations

$$\frac{J_1}{J_2} = \frac{R_3}{R_1}$$

$$\Delta V_{EB} = V_{EB1} - V_{EB2} = \frac{kT}{q} \ln\left(\frac{J_1}{J_2}\right)$$

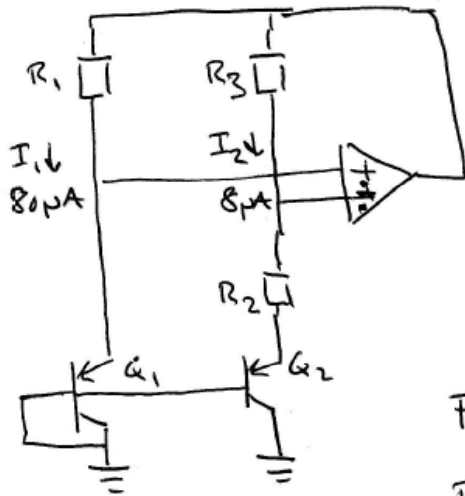
$$V_{ref} = V_{EB1} + \frac{R_3 kT}{R_2 q} \ln\left(\frac{R_3}{R_1}\right)$$

$$K = \frac{R_3}{R_2}$$



### EXAMPLE 8.5

Find the resistances of a bandgap reference based on fig. 8.27 a), where  $I_1 = 80 \mu\text{A}$ ,  $I_2 = 8 \mu\text{A}$ , and  $V_{EB1-0} = 0.65\text{V}$  at  $T = 300^\circ\text{K}$



Assuming that the sizes of  $Q_1$  and  $Q_2$  are the same

$$\Delta V_{EB} = \frac{kT_0}{q} \ln \frac{I_1}{I_2} = \frac{kT_0}{q} \cdot 2.30259 = 59.4 \text{ mV}$$

$$R_2 = \frac{V_{R2}}{I_2} = \underline{\underline{7.44 \text{ k}\Omega}} = \frac{59.4 \text{ mV}}{8 \mu\text{A}}$$

From (8.20) we know that  $V_{ref-0} = 1.24 \text{ V}$  at  $T_0 = 300^\circ\text{K}$

From (8.35) we get  $V_{R1} = V_{ref-0} - V_{EB1-0} = 1.24\text{V} - 0.65\text{V}$

$$\underline{\underline{V_{R1} = 0.59 \text{ V}}} \quad \underline{\underline{R_1 = \frac{0.59 \text{ V}}{80 \mu\text{A}} = 7.38 \text{ k}\Omega}}, \quad \underline{\underline{R_3 = \frac{0.59 \text{ V}}{8 \mu\text{A}} = 73.8 \text{ k}\Omega}}$$

## Example 8.5 (2)

$V_{R2} = \Delta V_{EB} = \frac{kT}{q} \ln(10)$ , since the sizes of  $Q_1$  and  $Q_2$  are assumed to be the same. (equations 8.36 and 8.40)

$$\Rightarrow V_{R2} = \frac{1.38 \cdot 10^{-23}}{1.62 \cdot 10^{-19}} (300) \cdot \ln(10) = 59.5 \text{ mV} \quad (V_{R2} = \Delta V_{EB})$$

$$U = RI \Leftrightarrow R = \frac{59.5 \text{ mV}}{8 \mu\text{A}} = 7.44 \text{ k}\Omega \quad \underline{R_2 = 7.44 \text{ k}\Omega}$$

Remember from 8.20 that for  $300^\circ\text{K}$  and  $m=2.3$ , (8.19) implies that  $V_{ref-0} = V_{G0} + (m-1) \frac{kT_0}{q} = 1.24 \text{ V}$  ( $V_{G0} = 1.206 \text{ V}$ )

Then we can get  $K$  from (8.21):

$$K = \frac{V_{G0} + (m-1) \frac{kT_0}{q} - V_{EB1-0}}{\frac{kT_0}{q} \ln\left(\frac{J_1}{J_2}\right)} = \frac{1.24 \text{ V} - 0.65 \text{ V}}{0.0258 \ln(10)} = 9.93$$

$$8.42 \text{ says that } K = \frac{R_3}{R_2} \Leftrightarrow \underline{R_3 = K \cdot R_2 = 73.9 \text{ k}\Omega}$$

$$8.39 : \frac{J_1}{J_2} = \frac{R_3}{R_1} \left( = \frac{I_1}{I_2} \right) \Leftrightarrow \frac{R_3}{R_1} = \frac{8 \mu\text{A}}{8 \mu\text{A}} \Leftrightarrow \underline{R_1 = \frac{R_3}{10} = 7.39 \text{ k}\Omega}$$



9. februar 2010

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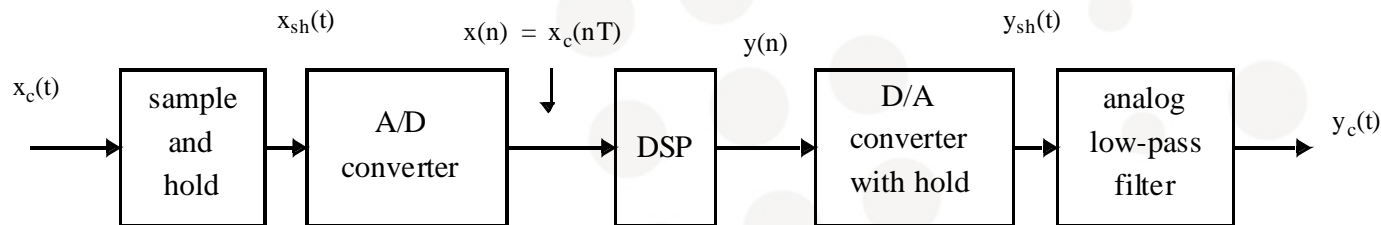
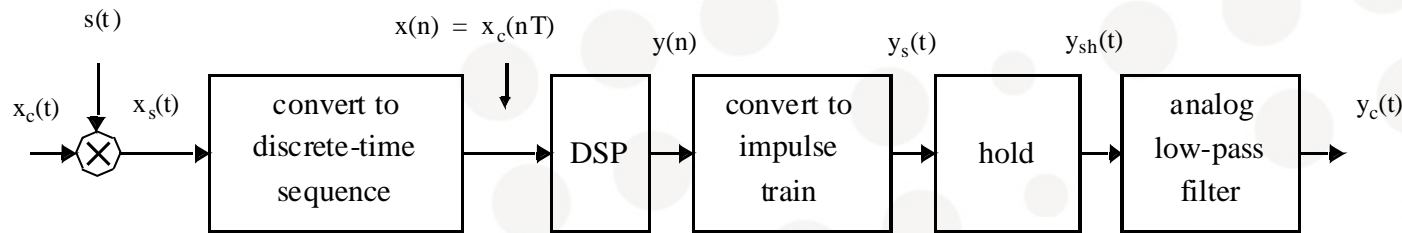
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# Chapetr 9; Discrete-time signals

- Discrete-time signal processing is heavily used in the design and analysis of **oversampling A/D and D/A converters** as well as **switched capacitor filtering** ;"SC-circuits".
- Switched Capacitor filters are classified as analog, since they use continous time analog values.

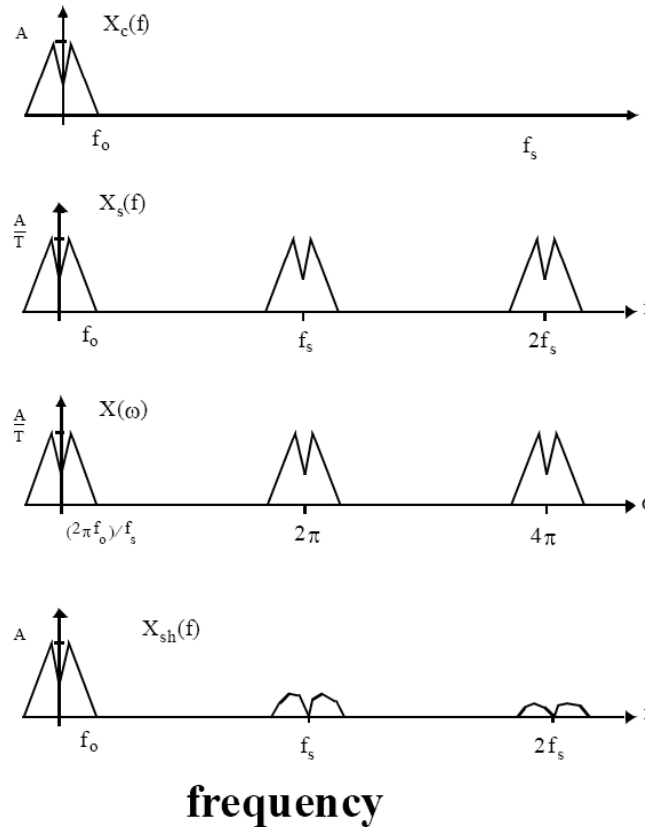
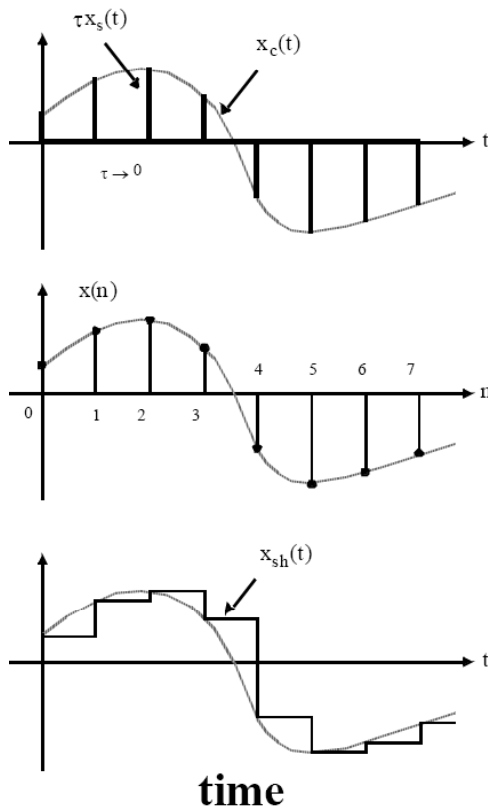


# Overview of signal spectra – conceptual and physical realizations



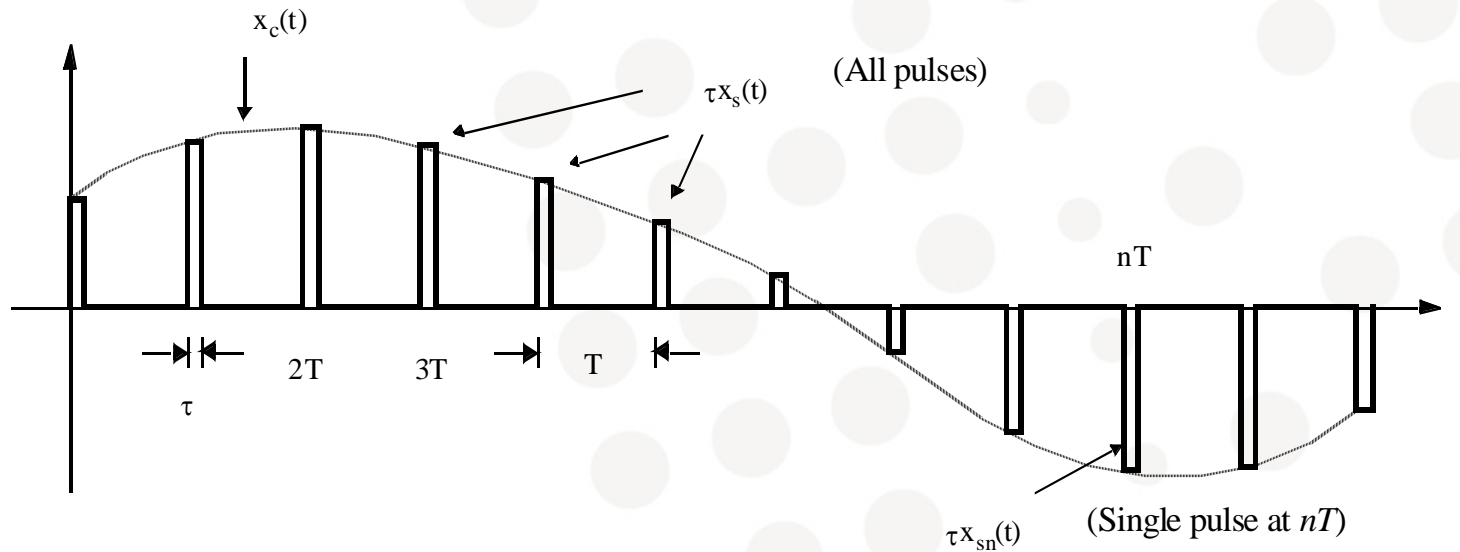
- An **anti-aliasing** filter (not shown) is assumed to band limit the continuous time signal,  $x_c(t)$ .
- **DSP** ("discrete-time signal processing") may be accomplished using fully digital processing or discrete-time analog circuits (ex.: SC-circ.).

# Signals in time, and frequency spectra



- $S(t)$ : periodic impulse train with period  $T$  ( $T=1/f_s$ )
- $x_s(t)$  has the same frequency spectrum as  $x_c(t)$ , but the baseband spectrum repeats every  $f_s$  (assuming no aliasing)
- $x(n)$  has the same frequency spectrum as  $x_c(t)$ , but the sampling frequency is normalized to 1
- The frequency spectrum of  $x_{sh}(t)$  is equal to that of  $x_s(t)$  multiplied by the  $\sin(x)/x$  response of the S/H.

# Laplace Transform of Discrete-Time Signals (1/3)



- The signal must be defined for all time
- For  $t=nT$ :
 
$$x_s(nT) = \frac{x_c(nT)}{\tau}$$
- $\tau$  is chosen such that the area under  $x_s(nT)$  equals the value of  $x_c(nT)$
- As  $\tau$  approaches 0, the height of  $x_s(nT)$  goes to  $\infty$

# Laplace Transform of Discrete-Time Signals (2/3)

- A single pulse at  $t=nT$  may be defined as:

- $\mathfrak{g}(t)$  is the step function: 
$$\mathfrak{g}(t) \equiv \begin{cases} 1 & (t \geq 0) \\ 0 & (t < 0) \end{cases}$$

- $x_s(t)$  may then be rewritten as a linear combination of a series of pulses,  $x_{sn}(t)$ , where  $x_{sn}(t)$  is zero everywhere except for a single pulse at  $nT$ :

$$x_{sn}(t) = \frac{x_c(nT)}{\tau} [\mathfrak{g}(t - nT) - \mathfrak{g}(t - nT - \tau)]$$

$x_s(t)$  is now defined for all time:

$$x_s(t) = \sum_{n=-\infty}^{\infty} x_{sn}(t)$$

# Laplace Transform of Discrete-Time Signals (3/3)

- The Laplace transform for  $x_{sn}(t)$  is:

$$X_{sn}(s) = \frac{1}{\tau} \left( \frac{1 - e^{-s\tau}}{s} \right) x_c(nT) e^{-snT}$$

- Since there is a linear relationship between  $x_s(t)$  and  $x_{sn}(t)$ , the Laplace transform of  $x_s(t)$  is:

$$X_s(s) = \frac{1}{\tau} \left( \frac{1 - e^{-s\tau}}{s} \right) \sum_{n=-\infty}^{\infty} x_c(nT) e^{-snT}$$

- When  $\tau$  approaches 0, the term before the sum equals 1 (eq. 9.7):

$$X_s(s) = \sum_{n=-\infty}^{\infty} x_c(nT) e^{-snT}$$

$$z \equiv e^{sT}$$

$$X(z) \equiv \sum_{n=-\infty}^{\infty} x_c(nT) z^{-n}$$

# Spectra of Discrete-Time Signals (1/2)

- The frequency spectrum of  $x_s(t)$  may be found by replacing  $s$  by  $j\omega$  in the Laplace transform (eq. 9.7).
- Another more intuitive approach is to use the property that **multiplication in the time domain equals convolution in the frequency domain**. Using this and  $\tau \rightarrow 0$ ,  $X_s(t)$  can be rewritten

$$x_s(t) = x_c(t)s(t)$$

- Define a pulse-train:

- The sampled signal is now:  $s(t) = \sum_{n=-\infty}^{\infty} \delta(t - nT)$

- The Fourier-transform of  $s(t)$  is:  $S(j\omega) = \frac{2\pi}{T} \sum_{k=-\infty}^{\infty} \delta(\omega - k\frac{2\pi}{T})$



# Spectra of Discrete-Time Signals (2/2)

- Writing (9.8) in the frequency domain:

$$X_s(j\omega) = \frac{1}{2\pi} X_c(j\omega) \otimes S(j\omega)$$

- The frequency spectrum of  $x_s(t)$  is then (eq. 9.12):

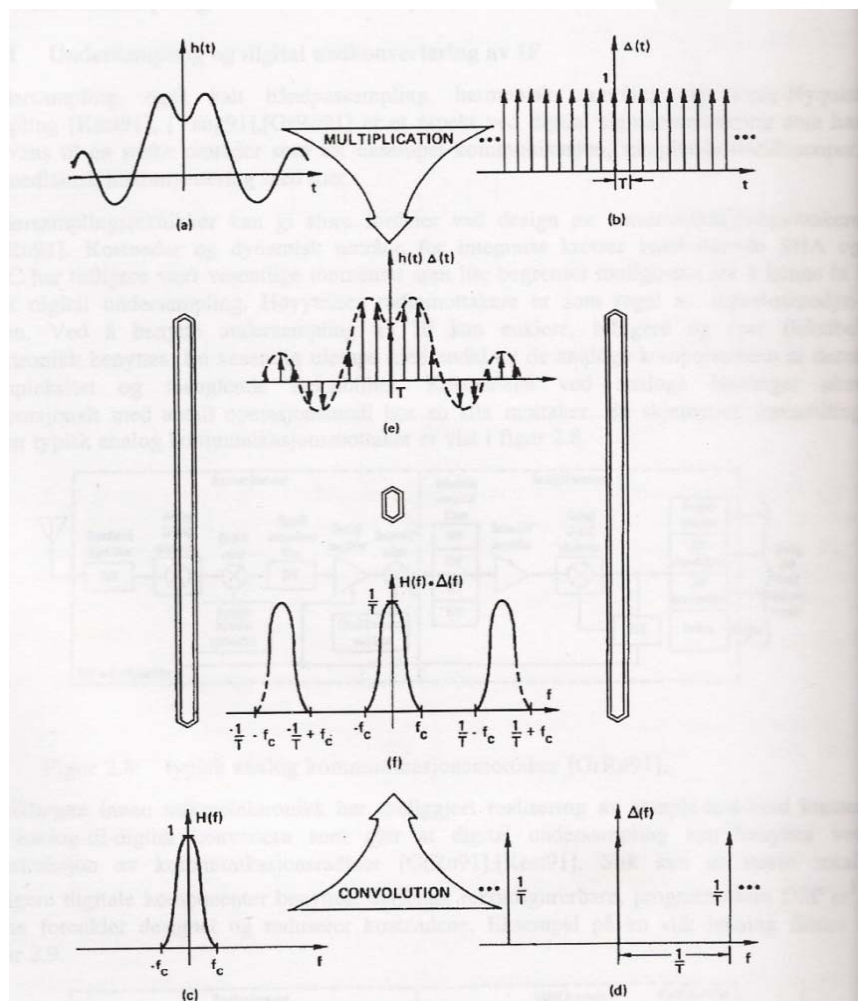
$$X_s(j\omega) = \frac{1}{T} \sum_{k=-\infty}^{\infty} X_c(j\omega - \frac{jk2\pi}{T})$$

which is periodic with period  $f_s$  ( 9.13:).

No aliasing occurs if  $f < f_s/2$

$$X_s(f) = \frac{1}{T} \sum_{k=-\infty}^{\infty} X_c(j2\pi f - jk2\pi f_s)$$

## Multiplication in the time domain equals convolution in the frequency domain

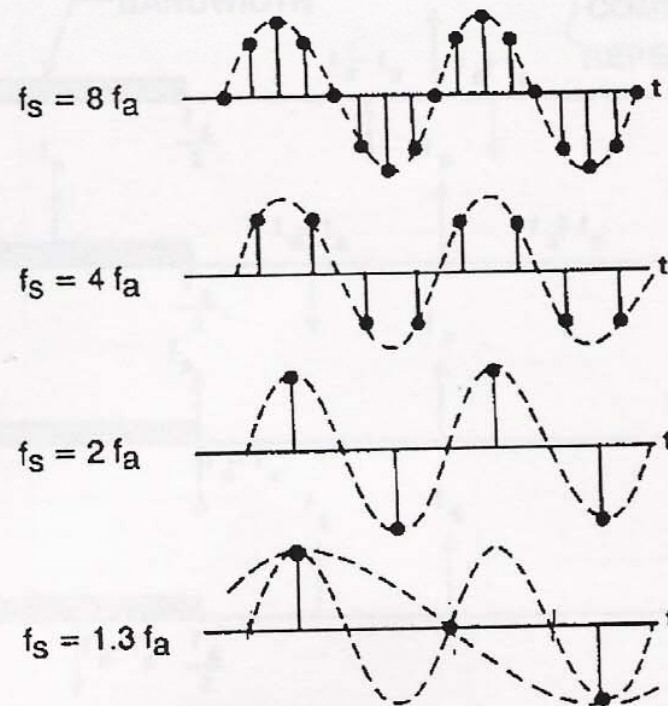


Figur 2.7: Sammenhengen mellom sampling i tids- og frekvensdomenet. Produktet av  $h(t)$ , figur 2.7 a), og  $\Delta(t)$ , i figur b), er lik den samplede kurveformen i figur c). Fouriertransformene av  $h(t)$  og  $\Delta(t)$  er gitt i henholdsvis figur 2.7 c) og d). Frekvenskonvolusjonsteoremet er illustrert ved at Fouriertransformen av  $h(t)*\Delta(t)$  er lik  $h(t)\Delta(t)$  [Brig74].

- Figure from E. O. Brigham: "The Fast Fourier Transform", Prentice Hall Inc., 1974., in S. Aunet: "BiCMOS sample-and-hold for satellitt-kommunikasjon", Cand. Scient. Thesis, University of Oslo, 1993.
- Wikipedia; Convolution:
- In [mathematics](#) and, in particular, [functional analysis](#), **convolution** is a mathematical [operation](#) on two [functions](#)  $f$  and  $g$ , producing a third function that is typically viewed as a modified version of one of the original functions. Convolution is similar to [cross-correlation](#).
- Computing the inverse of the convolution operation is known as [deconvolution](#).
- In [mathematics](#), the **Fourier transform** (often abbreviated **FT**) is an operation that [transforms](#) one [complex](#)-valued [function](#) of a [real variable](#) into another. In such applications as [signal processing](#), the domain of the original function is typically [time](#) and is accordingly called the [time domain](#). That of the new function is [frequency](#), and so the Fourier transform is often called the [frequency domain representation](#) of the original function. It describes which frequencies are present in the original function.

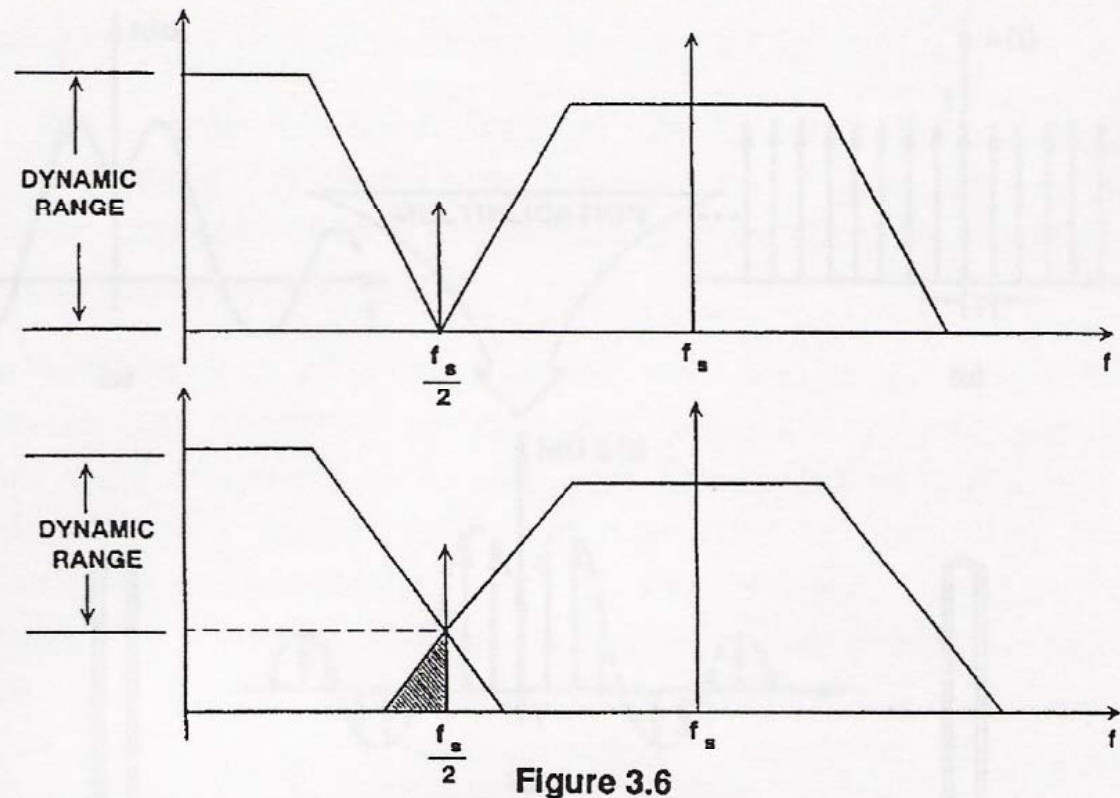
## Sampling at different frequencies

### 2.2 Signaler i tids- og frekvensdomenet, for ulike samplingsfrekvenser



Figur 2.4: Sampling ved ulike frekvenser, sett i tidsdomenet.  $f_s$  er samplingsfrekvensen, også kalt samplingsraten, mens  $f_a$  er frekvensen for det analoge signalet som samples. [Kest91].

## Aliasing and potential degrading of signal / noise



Figur 2.6: Aliasing og dynamisk område [Kest91] I det øverste tilfellet samples det slik at det dynamiske området beholdes. I det andre tilfellet overlapper frekvensspekterne slik at dynamisk område, eller signal/støy -forhold, reduseres.

# Z-Transform

- Discrete-time systems are most often analyzed using the z-transform which is equivalent to the Laplace-transform with the following substitution:
- Then the z-transform is defined as :

$$z \equiv e^{sT}$$

$$X(z) \equiv \sum_{n=-\infty}^{\infty} x_c(nT)z^{-n}$$



# Z-Transform

- Two important properties of the z-transform:
  - 1) If  $x(n) \leftrightarrow X(z)$ , then  $x(n-k) \leftrightarrow z^{-k}X(z)$
  - 2) Convolution in the time-domain is equal to multiplication in the freq. domain ( If  $y(n)=h(n) \otimes x(n)$ , then  $Y(z) = H(z)X(z)$ ). Similarly, multiplication in the time-domain equals convolution in the frequency domain
  - $X(z)$  is only related to the sampled sequence of numbers, while  $X_s(s)$  is the Laplace transform of  $x_s(t)$  when  $\tau \rightarrow 0$
  - The frequency response of  $X_s(f)$  is related to  $X(\omega)$  the following way:
$$X_s(f) = X\left(\frac{2\pi f}{f_s}\right)$$
  - Thus, the following scaling has been applied:

$$\omega = \frac{2\pi f}{f_s}$$

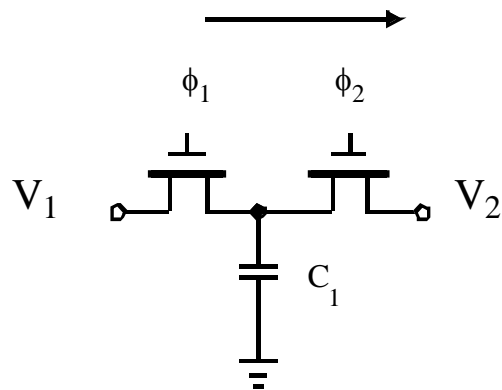


# Z-Transform

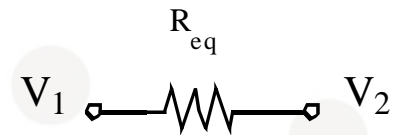
- Important observation:
  - Discrete-time signals have  $\omega$  in units of radians/sample
  - The original continuous-time signal have frequency units of cycles/second (Hertz) or radians / second. ( $2 \pi$  Radians  $\sim$  360 degrees)
- Example:
  - A continuous-time sinusoidal signal of 1kHz when sampled at 4 kHz will change by  $\pi/2$  radians between each sample. In such case the discrete time signal is defined to have a frequency of  $\pi/2$  radians per sample

# Next time, Tuesday 16th of February

- Chapter 9; 9.4 – 9.6
- Chapter 10; Switched Capacitor Circuits



$$\Delta Q = C_1(V_1 - V_2) \text{ every clock period}$$



$$R_{eq} = \frac{T}{C_1}$$