

The following people have participated in creating these solutions:
Nicolaas E. Groeneboom, Magnus Pedersen Lohne, Karl R. Leikanger

NOTE: There might be errors in the solution. If you find something which doesn't look right, please let me know

Partial solutions to problems: Part 1G

Problem 1G.1

In this exercise, we're asked to derive the expression for the mean kinetic energy of a particle in a degenerate gas. This gas no longer follows the normal M.B-distribution, which we have used in earlier exercises.

1. Let's summarize: We have a relation between $n(\vec{p})$ (the number density per volume per momentum space volume for particles with momentum \vec{p}) and $n(p)$ (the number density per real space volume for particles with absolute momentum p). This relation is given by $n(p)dp = 4\pi p^2 n(\vec{p})dp$, where we obtain the real-space volume element by integrating a sphere over the momentum-space for a fixed absolute momentum. We're now asked to find a relation between $n(p)$ and $n(E)$. We know that

$$E = \frac{p^2}{2m}$$

such that

$$p = \sqrt{2mE}$$

and

$$dp = \frac{1}{2\sqrt{2mE}} \cdot 2mdE = \sqrt{\frac{m}{2E}} dE$$

$$\frac{dp}{dE} = \sqrt{\frac{m}{2E}}$$

Now, we switch from $n(p)$ to $n(E)$ using the chain rule:

$$n(E) = n(p) \frac{dp}{dE} = n(p) \sqrt{\frac{m}{2E}}$$

and insert for $n(p) = 4\pi p^2 n(\vec{p})$:

$$n(E) = 4\pi p^2 n(\vec{p}) \sqrt{\frac{m}{2E}} = 4\sqrt{2}\pi m^{3/2} \sqrt{E} n(\vec{p})$$

where we substituted $p^2 = (2mE)$. We now insert for

$$n(\vec{p}) = \frac{2}{h^3} \frac{1}{e^{(p^2 - p_F^2)/(2mkT)} + 1}$$

$$n(E) = \frac{8\sqrt{2}\pi m^{3/2}}{h^3} \sqrt{E} \frac{1}{e^{(p^2 - p_F^2)/(2mkT)} + 1}$$

Rewriting, we find

$$n(E) = 4\pi \left(\frac{2m}{h^2}\right)^{3/2} \sqrt{E} \frac{1}{e^{(p^2 - p_F^2)/(2mkT)} + 1} = \frac{g(E)}{e^{(p^2 - p_F^2)/(2mkT)} + 1}$$

2. We continue by finding the mean kinetic energy of a particle in a degenerate gas:

$$\langle E \rangle = \int_0^\infty P(E) E dE$$

First, remember that the probability distribution is given by $n(E)$, but a probability distribution needs to be normalized such that

$$P(E) = Nn(E)$$

where N is found by

$$\int_0^\infty P(E) dE = N \int_0^{E_f} n(E) dE = 1$$

where the E_f -limit is because $n(E) = 0$ for $E > E_F$. The next thing we do is an approximation: in this energy range, the $e^{(p^2 - p_F^2)/(2mkT)}$ in $n(E)$ is much less than 1. We can then approximate $n(E) \approx g(E)$, and the integral becomes surprisingly simple. But first we need to normalize the distribution. For simplicity we define $K = 4\pi \left(\frac{2m}{\hbar^2}\right)^{3/2}$. Then

$$1 = N \int_0^{E_f} g(E) dE = NK \int_0^{E_f} E^{1/2} dE = NK \frac{2}{3} E_F^{3/2} = 1$$

such that $N = 3/2(K E_F^{3/2})$. The expectation value is thus

$$\begin{aligned} \langle E \rangle &= N \int_0^{E_f} g(E) E dE = \frac{3}{2} E_F^{-3/2} \int_0^{E_f} E^{3/2} dE \\ \langle E \rangle &= \frac{3}{2} E^{-3/2} \frac{2}{5} E_F^{5/2} = \frac{3}{5} E_F \end{aligned}$$