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Partial solutions to problems: part 2A

Exercise 2A.1

- 1. If cylinder B gets a smaller radius, it could pass through the inside of cylinder A.
- 2. From the reference of cylinder B, then cylinder A will have its radius shortened. Hence cylinder A will pass through the inside of cylinder B. This contradicts the results in the previous question.
- 3. $y \neq y'$ and $z \neq z'$ would give causal contradictions, as discussed in the previous question.

Exercise 2A.2

- 1. In the Earth-frame, a muon travelling 15km at 0.999c spends $15km/0.999c \approx 5.005 \cdot 10^{-5}$ s
- 2. Ignoring relativistic effects: the muons on average live for $2 \cdot 10^{-6}$ seconds each travelling at 0.999c and can therefore travel $2 \cdot 10^{-6} \cdot 0.999c \approx 600m$, which is less than 15km. We should therefore *not* expect to detect muons at the surface of Earth, when ignoring relativistic effects.
- 3. We will calculate the distance to the Earth and the time it takes to reach the Earth in both the Earth frame and the muon frame.
 - (a) In the Earth frame we have that $\Delta x = 15km = 5 \cdot 10^{-5} \text{s}$ and $\Delta t = 5.005 \cdot 10^{-5} \text{s}$.
 - (b) $\Delta x' = 0$
 - (c) $\Delta t' = \sqrt{\Delta t^2 \Delta x^2} = 2.2 \mu s.$

The muon decays in $2\mu s$ so some neutrinos which live a little longer than the average will reach the surface of the Earth (remember that $2\mu s$ is the mean life time).

4. In the reference frame of the relativistic particle, $\Delta s^2 = \Delta t^2 - 0^2 = \Delta t^2$. In the reference frame of the galaxy, $\Delta s'^2 = \Delta t'^2 - \Delta x'^2$. Equating the line elements and solving for Δt , we find

Whether this gives hope to future space travels or not: Transversing the galaxy in 50 days seems like a good thing, but there are a few negative sides:

- 1. The universe will have aged considerably during these 50 days (calculate for yourself)
- 3. The acceleration process would take much more than 50 days. A person can only handle a couple of G before turning liquid accelerating to almost c + decelerating + assuming constant 5G would take over 140 days alone.

Exercise 2A.3

- 1. In the reference frame of the clock, the time it takes for the light to travel between the two mirrors is L_0 (remember c=1).
- 2. We use this representation: (coordinate, time).

In the stationary frame:

A:(0,0) $B:(L_0,L_0)$

 $C:(L_0\cdot v,L_0)$

The line element in the stationary frame is given as $\Delta s_{AB}^2 = |(B - A)|^2 = (L_0 - 0)^2 - (L_0 - 0)^2 = 0$. This reflects that the line element for light is always zero.

In the train frame:

A:(0,0)

 $B:(x_B',t_B')$

 $C:(0,t'_{C})$

The line element in the train frame is given as $\Delta s_{AB}^{\prime 2} = |(B-A)|^2 = t_B^{\prime 2} - x_B^{\prime 2}$.

- 3. Setting the line elements equal $(\Delta s_{AB}^2 = \Delta s_{AB}'^2)$ shows that $t_B' = x_B'$. Since the speed of light is the same in all frames, light follows $\Delta x = \Delta t$ along all points of the tracetory.
- 4. In the stationary frame

$$\Delta s_{AC}^2 = |(C - A)|^2 = L_0^2 - (vL_0)^2$$

In the train frame

$$\Delta s_{AC}^{\prime 2} = |(C - A)|^2 = t_C^{\prime 2}$$

Invariance of the line element gives

$$t_C^{\prime 2} = L_0^2 - (vL_0)^2$$

such that

$$t_C' = L_0 \sqrt{1 - v^2} = L_0 / \gamma$$

5.

$$\Delta s_{BC}^2 = |(C - B)|^2 = (L_0 - L_0)^2 - (L_0 v - L_0)^2 = -(L_0 v - L_0)^2$$
$$\Delta (s_{BC}')^2 = |(C - B)|^2 = (t_c' - t_B')^2 - (0 - x_B')^2 = (L_0/\gamma - t_B')^2 - t_B'^2$$

Using the invariance of the line element:

$$(L_0/\gamma - t_B')^2 - (t_B')^2 = (L_0^2/\gamma^2 - 2L_0/\gamma t_B') = -L_0^2(1-v)^2$$

Using that $1/\gamma^2 = (1 - v^2)$ and solving for t_B' :

$$t'_B = \frac{\gamma L_0}{2} ((1 - v^2) + (1 - v)^2) = L_0 \gamma (1 - v)$$

6. In the stationary frame, $D = (0, 2L_0)$ while in the train frame, $D = (-v \cdot t'_D, t'_D)$. In the stationary frame,

$$\Delta s_{AD}^2 = |(D - A)|^2 = 4L_0^2$$

while in the train frame

$$\Delta(s'_{AD})^2 = |(D-A)|^2 = t_D^2 - v^2 t_D'^2$$

Equating these line elements gives:

$$4L_0^2 = |(D - A)|^2 = t_D^2 - v^2 t_D^2$$

Solving for t'_D gives

$$t'_D = \frac{2L_0}{\sqrt{1-v^2}} = 2L_0\gamma$$

7. In the frame of the train: $\Delta t'_{AB} = t'_B - t'_A = t'_B = L_0 \gamma (1-v)$ while $\Delta t'_{BD} = t'_D - t'_B = 2L_0 \gamma - L_0 \gamma (1-v) = L_0 \gamma (1+v)$. In the stationary reference frame, $\Delta t_{AB} = L_0$ while $\Delta t_{BD} = 2L_0 - L_0 = L_0$. As the speed of light is equal in all reference frames, the time experienced is different in the two coordinate systems.

Exercise 2A.4

1. We convert to the physical distance, using $d = r\theta$. We find

$$d = 2.6 \cdot 10^9 ly \cdot (2. \cdot 10^{-3}) \cdot 2\pi/60/60/360 \approx 2.7425.2ly'$$

Dividing by 3 years:

$$v = 25.2ly/3 \approx 8.4c$$

2. The observed time signal Δt_{obs} we observed equals the real time interval Δt minus the length the jet transversed in the radial coordinate. This length is given as $s = v * \Delta t$, and projecting the jet onto the radial coordinate gives $s = \Delta t v \cos \theta$, such that

$$\Delta t_{obs} = \Delta t - \Delta t v \cos \theta$$

3. We use that

$$v_{obs} = \frac{s_{obs}}{\Delta t_{obs}}$$

From the figure in the exercise description, we note that the transverse length is given as $\Delta tv \sin \theta$ such that

$$v_{obs} = \frac{\Delta t v \sin \theta}{\Delta t - \Delta t v \cos \theta} = \frac{v \sin \theta}{1 - v \cos \theta}$$

4. We use that $\sin \theta = \cos \theta = 1/\sqrt{2}$. Insert and ask what values of observed velocity will be greater than the speed of light:

$$\frac{v/\sqrt{2}}{1-v/\sqrt{2}}>1$$

$$\frac{v}{\sqrt{2}-v} > \sqrt{2}$$

Solve for v:

$$v > \frac{c\sqrt{2}}{2} \approx 0.7$$

Exercise 2A.5

1. The coordinates are defined as (position, time). Event A happens at (0,0) in both the laboratory and moving frame, while event B happens at (x,t) in the laboratory frame and (0,t') in the moving frame. Then $\Delta t_{AB} = t$ and $\Delta t'_{AB} = t'$. We then use that $\Delta t = \gamma \Delta t'$, or $t = \gamma t'$. This resembles the first equation (8) from the problem text:

$$t = f(v)x' + g(v)t' = \gamma t'.$$

Since x' = 0 we have that $g(v) = \gamma$.

2. Event B happens at position x = vt in the laboratory frame, or $x = vt = v\gamma t'$. Equation (9) then gives

$$x = h(x)x' + k(v)t' = 0 + k(v)t' = v\gamma t'$$

as x' = 0 in the moving frame. Hence $k(v) = v\gamma$.

3. Even A is still the same as in question 1, but event B is now (L_0, t') in the moving frame and (L, 0) in the laboratory frame. Equation (6) shows that $L = L_0/\gamma$ such that $x = x'/\gamma$. Then equation (11) reads

$$x' = h(-v)x - v\gamma t = h(-v)x = \gamma x$$

such that $h(v) = h(-v) = \gamma$ (we used that t = 0 since A and B are simultaneous in the lab frame and A happens at t = 0)

4. We only need to decide on f(v). The space-time interval Δs_{AB}^2 is then expressed as

$$\Delta s_{AB}^2 = |(B - A)|^2 = t^2 - x^2 = (f(v)x' + \gamma t')^2 - (\gamma x' + v\gamma t')^2$$

while

$$\Delta s_{AB}^{\prime 2} = |(B - A)|^2 = t^{\prime 2} - x^{\prime 2}$$

Equating these two line elements and expanding the square:

$$t'^2 - x'^2 = f(v)^2 x'^2 + \gamma^2 t'^2 + 2f(v)x'\gamma t' - \gamma^2 x'^2 - v^2 \gamma^2 t'^2 - 2\gamma^2 x' v t'$$

The simplest way to show $f(v) = \gamma v$ is by insertion:

$$t'^2 - x'^2 = \gamma^2 v^2 x'^2 + \gamma^2 t'^2 + 2v x' \gamma^2 t' - \gamma^2 x'^2 - v^2 \gamma^2 t'^2 - 2\gamma^2 x' v t'$$

Extract the γ^2 :

$$t'^{2} - x'^{2} = \gamma^{2}(v^{2}x'^{2} + t'^{2} + 2vx't' - x'^{2} - v^{2}t'^{2} - 2x'vt')$$

Notice how two of the terms cancel. So:

$$t'^2 - x'^2 = \gamma^2 (v^2 x'^2 + t'^2 - x'^2 - v^2 t'^2)$$

which equals

$$t'^2 - x'^2 = \gamma^2 (v^2 (x'^2 - t'^2) + (t'^2 - x'^2)) = \gamma^2 (t'^2 - x'^2)(1 - v^2)$$

But $\gamma^2 = 1/(1-v^2)$, and the equation holds. Thus, $f(v) = \gamma v$

Exercise 2A.6

1. This is done as in problem 4: We use this representation: (coordinate, time). In the stationary frame:

A:(0,0)

 $B:(L_0,L_0)$

 $D:(0,2L_0)$

In the train frame:

A:(0,0)

 $B: (x'_B, t'_B) \\ D: (x'_D, t'_D)$

2. The lorentz transformations states that

$$t'_B = f(-v)x_B + g(-v)t_B = -\gamma vL_0 + \gamma L_0 = \gamma L_0(1-v)$$

 t'_D is then

$$t'_{D} = f(-v)x_{D} + g(-v)t_{D} = 2\gamma L_{0}$$

as from problem 4. For the third thick we have $x_3=L_0$ and $t_3=3L_0$ giving

$$t_3' = -v\gamma L_0 + 3\gamma L_0 = \gamma L_0(3 - v).$$

Finally the fourth tick happening at $x_4 = 0$ and $t_4 = 4L_0$ gives

$$t_4' = 4\gamma L_0,$$

so that the time between the next to ticks is $\Delta t'_{34} = \gamma L_0(4 - (3 - v)) =$ $\gamma L_0(1+v)$ exactly as for the first two ticks.

Exercise 2A.7

See separate document for this exercise.