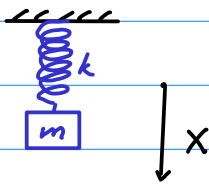


Svingninger

Eks: Fjær og masse.

Fjærkraft: $-kx$



$$mg - kx = m\ddot{x}. \quad \text{Likevekt (i ro): } mg = kx_0 \Rightarrow x_0 = \frac{mg}{k}$$

Def.: $x = X - X_0$ (avstand fra likevekt), $\ddot{x} = \ddot{X}$.

$$mg - kX = mg - k(x + X_0) = mg - kX_0 - kx = -kx$$

Gir: $\ddot{x} = -\frac{k}{m}x$ eller $\ddot{x} = -\omega^2 x$, der $\omega = \sqrt{\frac{k}{m}}$.

Løsning: $x(t) = C_1 e^{i\omega t} + C_2 e^{-i\omega t}$, $e^{i\omega t} = \cos \omega t + i \sin \omega t$ (Euler)

$$x(t) = C_1 \cos \omega t + C_2 \cos \omega t + i C_1 \sin \omega t + i C_2 \sin(-\omega t) \\ = (C_1 + C_2) \cos \omega t + i(C_1 - C_2) \sin \omega t$$

$$= B_1 \cos \omega t + B_2 \sin \omega t, \quad B_1 = C_1 + C_2, \quad B_2 = i(C_1 - C_2)$$

Kan også skrives $x(t) = D \cos(\omega t - \delta)$, for konst. D og δ (ukesoppgave)

Altså: Svingninger!

Kompleks representasjon: $x(t) = D \operatorname{Re} e^{i(\omega t - \delta)}$

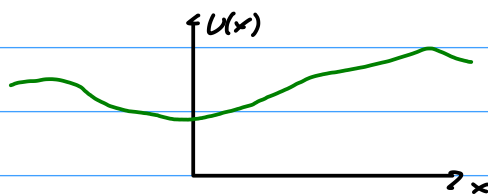
$$= \operatorname{Re} D e^{i(\omega t - \delta)}$$

$$= \operatorname{Re} D e^{-i\delta} e^{i\omega t}$$

$$= \operatorname{Re} C e^{i\omega t}, \quad C \text{ kompleks konst.}$$

Fordel: Lettere å addere...

Mer generelt: Potensiell energi $U(x) = U(0) + U'(0)x + \frac{1}{2}U''(0)x^2 + \dots$



effektivt en fjærkonst.

For små $|x|$: $m\ddot{x} = -\frac{dU}{dx} = \cancel{U'(0)} + U''(0)x \Rightarrow$ harm. osc.

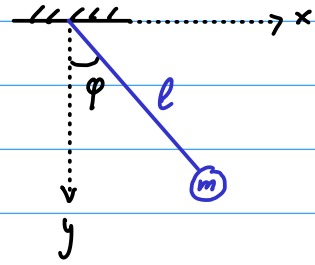
Eks: Pendel

$$L = \frac{1}{2} m l^2 \dot{\varphi}^2 - mgl(1 - \cos\varphi)$$

$$\text{Lagrange: } \frac{\partial L}{\partial \varphi} = \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{\varphi}} \right)$$

$$\Rightarrow -mgl \sin\varphi = \frac{d}{dt} m l^2 \dot{\varphi} = m l^2 \ddot{\varphi}$$

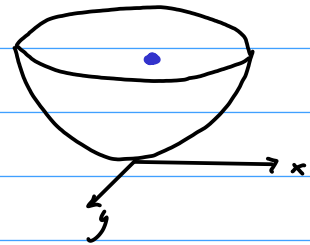
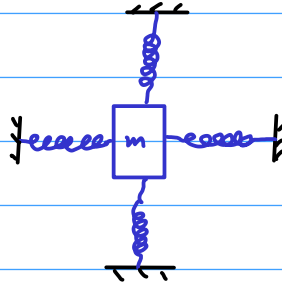
$$\Rightarrow \ddot{\varphi} = -\frac{g}{l} \sin\varphi = -\frac{g}{l} \left(\varphi - \frac{\varphi^3}{3!} + \dots \right) \approx -\frac{g}{l} \varphi = -\omega^2 \varphi, \quad \omega = \sqrt{\frac{g}{l}}$$



Eks: 2D-oscillator:

$$\ddot{x} = -\omega^2 x$$

$$\ddot{y} = -\omega^2 y$$

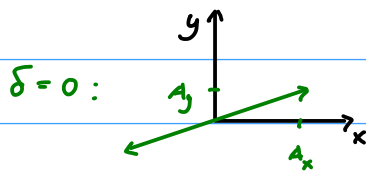


$$x(t) = A_x \cos(\omega t - \delta_x)$$

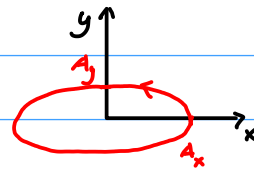
$$y(t) = A_y \cos(\omega t - \delta_y)$$

Redefinere t : $x(t) = A_x \cos(\omega t)$

$$y(t) = A_y \cos(\omega t - \delta)$$

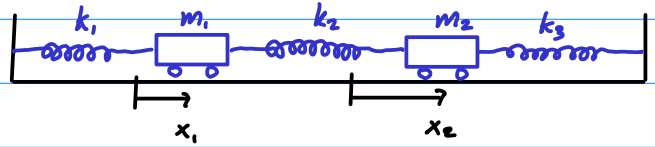


$\delta = \frac{\pi}{2}$:



Koblede oscillatorer (T. kap. 11)

$$\begin{aligned} \text{Bil 1: } m_1 \ddot{x}_1 &= -k_1 x_1 + k_2 (x_2 - x_1) \\ &= -(k_1 + k_2) x_1 + k_2 x_2 \end{aligned}$$



$$\begin{aligned} \text{Bil 2: } m_2 \ddot{x}_2 &= -k_2 (x_2 - x_1) - k_3 x_2 \\ &= k_2 x_1 - (k_2 + k_3) x_2 \end{aligned}$$

$$M \ddot{x} = -Kx, \quad M = \begin{bmatrix} m_1 & 0 \\ 0 & m_2 \end{bmatrix}, \quad x = \begin{bmatrix} x_1 \\ x_2 \end{bmatrix}, \quad K = \begin{bmatrix} k_1 + k_2 & -k_2 \\ -k_2 & k_2 + k_3 \end{bmatrix}$$

Ser etter løsninger på formen: $x(t) = \begin{bmatrix} \alpha_1 \cos(\omega t - \delta_1) \\ \alpha_2 \cos(\omega t - \delta_2) \end{bmatrix} = \begin{bmatrix} \alpha_1 \operatorname{Re} e^{i(\omega t - \delta_1)} \\ \alpha_2 \operatorname{Re} e^{i(\omega t - \delta_2)} \end{bmatrix}$

$$= \operatorname{Re} \underbrace{\begin{bmatrix} a_1 \\ a_2 \end{bmatrix}}_a e^{i\omega t}, \quad \begin{aligned} a_1 &= \alpha_1 e^{-i\delta_1} \\ a_2 &= \alpha_2 e^{-i\delta_2} \end{aligned}$$

$z(t)$

Hvis $z(t)$ er en løsning av $M\ddot{z} = -Kz$, da er $x(t) = \operatorname{Re} z(t)$ en løsning av $M\ddot{x} = -Kx$.

$$\begin{aligned} M\ddot{z} = -Kz &\Rightarrow M(i\omega)^2 a e^{i\omega t} = -K a e^{i\omega t} \\ &\Rightarrow (K - \omega^2 M) a = 0 \end{aligned}$$

Matriseteori: $Aa = 0$ og $a \neq 0 \Rightarrow \det A = 0$, fordi:

$$\begin{aligned} 0 &= \begin{bmatrix} A_{11} & A_{12} \\ A_{21} & A_{22} \end{bmatrix} \begin{bmatrix} a_1 \\ a_2 \end{bmatrix} \Rightarrow \begin{aligned} A_{11} a_1 + A_{12} a_2 &= 0 \\ A_{21} a_1 + A_{22} a_2 &= 0 \end{aligned} \Rightarrow \begin{aligned} a_2 &= -a_1 \frac{A_{11}}{A_{12}} \\ a_2 &= -a_1 \frac{A_{21}}{A_{22}} \end{aligned} \\ &\Rightarrow \frac{A_{11}}{A_{12}} = \frac{A_{21}}{A_{22}} \Rightarrow A_{11} A_{22} - A_{12} A_{21} = 0 \Rightarrow \det A = 0 \end{aligned}$$

Må ha $\det(K - \omega^2 M) = 0$.

Spesialtilfelle: $m_1 = m_2 = m$ og $k_1 = k_2 = k_3 = k$:

$$M = \begin{bmatrix} m & 0 \\ 0 & m \end{bmatrix}, \quad K = \begin{bmatrix} 2k & -k \\ -k & 2k \end{bmatrix}$$

$$\det(K - \omega^2 M) = \begin{vmatrix} 2k - m\omega^2 & -k \\ -k & 2k - m\omega^2 \end{vmatrix} = (2k - m\omega^2)^2 - k^2 \\ = (k - m\omega^2)(3k - m\omega^2) = 0$$

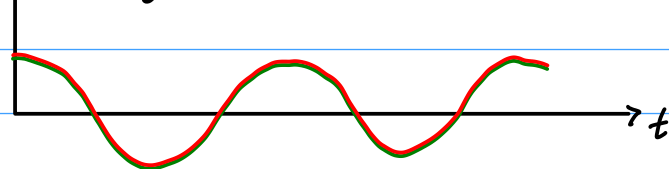
Løsninger: $\omega_1 = \sqrt{\frac{k}{m}}$, $\omega_2 = \sqrt{\frac{3k}{m}}$

Modus ω_1 : $\omega_1 = \sqrt{\frac{k}{m}}$, $K - \omega_1^2 M = \begin{bmatrix} k & -k \\ -k & k \end{bmatrix}$.

$$(K - \omega_1^2 M)a = 0 \Rightarrow k \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix} \begin{bmatrix} a_1 \\ a_2 \end{bmatrix} = 0 \Rightarrow a_1 = a_2 = A e^{-i\delta}$$

$$z(t) = \begin{bmatrix} A \\ A \end{bmatrix} e^{i(\omega_1 t - \delta)}, \quad x(t) = \begin{bmatrix} A \\ A \end{bmatrix} \cos(\omega_1 t - \delta), \quad \begin{aligned} x_1(t) &= A \cos(\omega_1 t - \delta) \\ x_2(t) &= A \cos(\omega_1 t - \delta) \end{aligned}$$

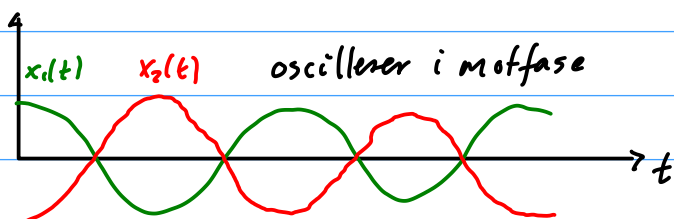
$x_1(t)$ og $x_2(t)$ oscillerer likt, i fase.



Modus ω_2 : $\omega_2 = \sqrt{\frac{3k}{m}}$, $K - \omega_2^2 M = \begin{bmatrix} -k & -k \\ -k & -k \end{bmatrix}$, $(K - \omega_2^2 M)a = 0 \Rightarrow -k \begin{bmatrix} 1 & 1 \\ 1 & 1 \end{bmatrix} \begin{bmatrix} a_1 \\ a_2 \end{bmatrix} = 0$

$$a_1 = -a_2 = A e^{i\delta}$$

$$z(t) = \begin{bmatrix} A \\ -A \end{bmatrix} e^{i(\omega_2 t - \delta)}, \quad x(t) = \begin{bmatrix} A \\ -A \end{bmatrix} \cos(\omega_2 t - \delta) = \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}$$



Generell løsning: $x(t) = A_1 \begin{bmatrix} 1 \\ 1 \end{bmatrix} \cos(\omega_1 t - \delta_1) + A_2 \begin{bmatrix} 1 \\ -1 \end{bmatrix} \cos(\omega_2 t - \delta_2)$, $\omega_2 = \sqrt{3}\omega_1$

Spesialtilfelle svalt koblete oscillatorer:

$$k_1 = k_3 = k, \quad k_2 \ll k. \quad K = \begin{bmatrix} k+k_2 & -k_2 \\ -k_2 & k+k_2 \end{bmatrix}, \quad K - \omega^2 M = \begin{bmatrix} k+k_2 - m\omega^2 & -k_2 \\ -k_2 & k+k_2 - m\omega^2 \end{bmatrix}$$

$$0 = \det(K - \omega^2 M) \Rightarrow \dots \Rightarrow \omega_1 = \sqrt{\frac{k}{m}}, \quad \omega_2 = \sqrt{\frac{k+2k_2}{m}}$$

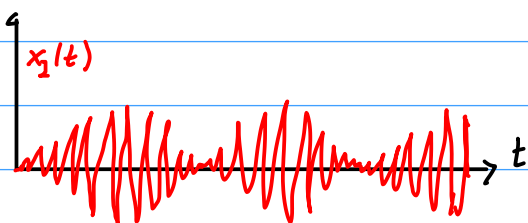
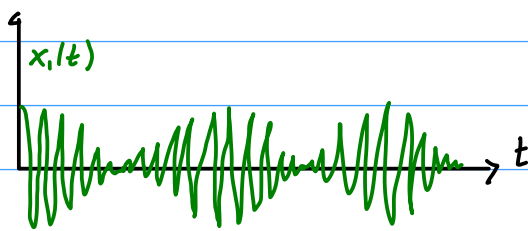
$$\text{Modus } \omega_1: \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix} \begin{bmatrix} a_1 \\ a_2 \end{bmatrix} = 0 \Rightarrow a_1 = a_2 = C_1, \quad z(t) = C_1 \begin{bmatrix} 1 \\ 1 \end{bmatrix} e^{i\omega_1 t}$$

$$\text{Modus } \omega_2: \begin{bmatrix} 1 & 1 \\ 1 & 1 \end{bmatrix} \begin{bmatrix} a_1 \\ a_2 \end{bmatrix} = 0 \Rightarrow a_1 = -a_2 = C_2, \quad z(t) = C_2 \begin{bmatrix} 1 \\ -1 \end{bmatrix} e^{i\omega_2 t}$$

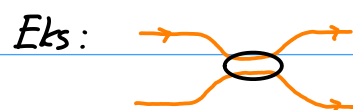
$$\text{Def } \omega_0 = \frac{\omega_1 + \omega_2}{2}, \quad \epsilon = \frac{\omega_2 - \omega_1}{2} : \quad \omega_1 = \omega_0 - \epsilon, \quad \text{der } \epsilon \ll \omega_0. \\ \omega_2 = \omega_0 + \epsilon$$

$$\text{Generell losn.: } z(t) = C_1 \begin{bmatrix} 1 \\ 1 \end{bmatrix} e^{i(\omega_0 - \epsilon)t} + C_2 \begin{bmatrix} 1 \\ -1 \end{bmatrix} e^{i(\omega_0 + \epsilon)t} \\ = \left(C_1 \begin{bmatrix} 1 \\ 1 \end{bmatrix} e^{-i\epsilon t} + C_2 \begin{bmatrix} 1 \\ -1 \end{bmatrix} e^{i\epsilon t} \right) e^{i\omega_0 t}$$

$$\text{Anta } C_1 = C_2 = \frac{A}{2}. \quad z(t) = A \begin{bmatrix} \cos \epsilon t \\ -i \sin \epsilon t \end{bmatrix} e^{i\omega_0 t}, \quad x_1(t) = A \cos \epsilon t \cos \omega_0 t \\ x_2(t) = A \cos \epsilon t \sin \omega_0 t \\ \hat{=} \text{Re}(-ie^{i\omega_0 t}) = \text{Re} e^{i\omega_0 t - i\frac{\pi}{2}} \\ = \cos(\omega_0 t - \frac{\pi}{2}) = \sin \omega_0 t$$



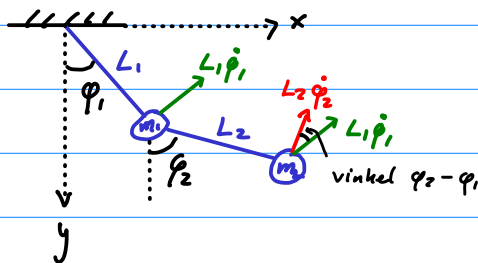
Stor oscillator bil 1,
kobler svalt til bil 2,
kobler tilbake osv.



—
—
—
piano

Eks: Dobbelpendel

Bruke φ_1 og φ_2 som koordinater.



$$U_1 = m_1 g L_1 (1 - \cos \varphi_1)$$

$$U_2 = m_2 g [L_1 (1 - \cos \varphi_1) + L_2 (1 - \cos \varphi_2)]$$

$$T_1 = \frac{1}{2} m_1 (L_1 \dot{\varphi}_1)^2, \quad T_2 = \frac{1}{2} m_2 [(L_1 \dot{\varphi}_1)^2 + (L_2 \dot{\varphi}_2)^2 + 2L_1 L_2 \dot{\varphi}_1 \dot{\varphi}_2 \cos(\varphi_2 - \varphi_1)]$$

(fordi $\vec{v} = \vec{v}_1 + \vec{v}_2 \Rightarrow \vec{v}^2 = \vec{v}_1^2 + \vec{v}_2^2 + 2\vec{v}_1 \cdot \vec{v}_2$)

$\mathcal{L} = T - U = \dots$ gir Lagrange's ligninger...

Tilnærmetse: Små utslag $|\varphi_1|$ og $|\varphi_2|$, og små vinkelhast. $|\dot{\varphi}_1|$ og $|\dot{\varphi}_2|$.

$$\cos \varphi = 1 - \frac{\varphi^2}{2} + \dots$$

$$\Rightarrow T = \frac{1}{2} (m_1 + m_2) L_1^2 \dot{\varphi}_1^2 + \frac{1}{2} m_2 L_2^2 \dot{\varphi}_2^2 + m_2 L_1 L_2 \dot{\varphi}_1 \dot{\varphi}_2 \leftarrow \text{her kan vi sette } \cos(\varphi_1 - \varphi_2) \approx 1$$

> siden allerede multiplisert med små størrelser $\dot{\varphi}_1, \dot{\varphi}_2$

$$U = \frac{1}{2} (m_1 + m_2) g L_1 \varphi_1^2 + \frac{1}{2} m_2 g L_2 \varphi_2^2 \leftarrow \cos \varphi_{1,2} = 1 - \frac{\varphi_{1,2}^2}{2}$$

$$\frac{d}{dt} \left(\frac{\partial \mathcal{L}}{\partial \dot{\varphi}_1} \right) = \frac{\partial \mathcal{L}}{\partial \varphi_1} : (m_1 + m_2) L_1^2 \ddot{\varphi}_1 + m_2 L_1 L_2 \ddot{\varphi}_2 = -(m_1 + m_2) g L_1 \varphi_1$$

$$\frac{d}{dt} \left(\frac{\partial \mathcal{L}}{\partial \dot{\varphi}_2} \right) = \frac{\partial \mathcal{L}}{\partial \varphi_2} : m_2 L_1 L_2 \ddot{\varphi}_1 + m_2 L_2^2 \ddot{\varphi}_2 = -m_2 g L_2 \varphi_2$$

$$M \ddot{\varphi} = -K \varphi, \quad M = \begin{bmatrix} (m_1 + m_2) L_1^2 & m_2 L_1 L_2 \\ m_2 L_1 L_2 & m_2 L_2^2 \end{bmatrix}, \quad K = \begin{bmatrix} (m_1 + m_2) g L_1 & 0 \\ 0 & m_2 g L_2 \end{bmatrix}$$

Spesielttilfelle: $m_1 = m_2 = m, L_1 = L_2 = L$:

$$M = mL^2 \begin{bmatrix} 2 & 1 \\ 1 & 1 \end{bmatrix}, \quad K = mL^2 \begin{bmatrix} 2\omega_0^2 & 0 \\ 0 & \omega_0^2 \end{bmatrix}, \quad \omega_0 = \sqrt{\frac{g}{L}}$$

Kompleks representasjon: $\varphi(t) = \operatorname{Re} z(t)$, $z(t) = a e^{i\omega t} = \begin{bmatrix} a_1 \\ a_2 \end{bmatrix} e^{i\omega t}$,
 der ω er en ukjent "eigenfrekvens".

$$K - \omega^2 M = mL^2 \begin{bmatrix} 2(\omega_0^2 - \omega^2) & -\omega^2 \\ -\omega^2 & \omega_0^2 - \omega^2 \end{bmatrix}, \quad (K - \omega^2 M) a = 0$$

$$\det(K - \omega^2 M) = 0 \Rightarrow 2(\omega_0^2 - \omega^2)^2 - \omega^4 = 0 \Rightarrow \omega^4 - 4\omega_0^2 \omega^2 + 2\omega_0^4 = 0$$

$$\Rightarrow \omega^2 = (2 \pm \sqrt{2}) \omega_0^2$$

$$\text{Løsning } \omega_1^2 = (2 - \sqrt{2}) \omega_0^2: \quad K - \omega_1^2 M = mL^2 \begin{bmatrix} 2(1 - (2 - \sqrt{2}))\omega_0^2 & -(2 - \sqrt{2})\omega_0^2 \\ -(2 - \sqrt{2})\omega_0^2 & (1 - (2 - \sqrt{2}))\omega_0^2 \end{bmatrix}$$

$$= mL^2 \omega_0^2 (\sqrt{2} - 1) \begin{bmatrix} 2 & -\sqrt{2} \\ -\sqrt{2} & 1 \end{bmatrix}$$

$$\begin{bmatrix} 2 & -\sqrt{2} \\ -\sqrt{2} & 1 \end{bmatrix} \begin{bmatrix} a_1 \\ a_2 \end{bmatrix} = 0 \Rightarrow a_2 = \sqrt{2} a_1, \quad a_1 = A_1 e^{-i\delta_1}$$

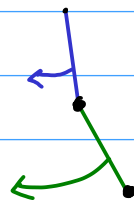
$$\varphi(t) = \begin{bmatrix} \varphi_1(t) \\ \varphi_2(t) \end{bmatrix} = \operatorname{Re} \left\{ \begin{bmatrix} 1 \\ \sqrt{2} \end{bmatrix} A_1 e^{-i\delta_1} e^{i\omega_1 t} \right\}$$

$$= \underline{A_1 \begin{bmatrix} 1 \\ \sqrt{2} \end{bmatrix} \cos(\omega_1 t - \delta_1)}$$

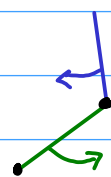
$$\text{Løsning } \omega_2^2 = (2 + \sqrt{2}) \omega_0^2: \quad \varphi(t) = \underline{A_2 \begin{bmatrix} 1 \\ -\sqrt{2} \end{bmatrix} \cos(\omega_2 t - \delta_2)}$$

Generell løsning, gyldig for små utslag og hastigheter:

$$\varphi(t) = A_1 \begin{bmatrix} 1 \\ \sqrt{2} \end{bmatrix} \cos(\omega_1 t - \delta_1) + A_2 \begin{bmatrix} 1 \\ -\sqrt{2} \end{bmatrix} \cos(\omega_2 t - \delta_2)$$



ω_1



ω_2