

Lecture 9

Canonical ensemble

Average values, equipartition, speed
distribution, free energy

Recap+ plan

Statistical mechanics

Models	Counting states @constant (NVU)	Counting states @constant (NVT)	Counting states @constant (μVT)
<ul style="list-style-type: none"> Two-state: Einstein crystal Ideal gas MD 	$\Omega(N, n) = \frac{N!}{(N-n)!n!}$ $\Omega(N, q) = \frac{(N-1+q)!}{(N-1)!q!}$ $\Omega(U, V) = f(N)V^N U^{3N/2}$ measure & average	Partition function $Z = \sum_i e^{-\beta \varepsilon_i}$ Maxwell-Boltzmann speed distribution	Gibbs sum $Z_G = \sum_i e^{-\beta(\varepsilon_i - \mu N_i)}$
Interacting systems	-----> Equilibrium irreversible $\Omega \rightarrow \max \Omega$		

$$S = k \ln \Omega$$

$$F = -kT \ln Z$$

$$U - TS = -kT \ln Z_G$$

Thermodynamics

$$\Delta S_{tot} \geq 0$$

$$T_1 = T_2$$

$$\Delta F \leq 0$$

$$\Delta G \leq 0$$

$$P_1 = P_2$$

$$\mu_1 = \mu_2$$

Implications:

- Thermodynamic identity $dU = TdS - PdV$

$$F \equiv U - TS$$

$$dF = -PdV - SdT$$

$$G \equiv F + PV$$

$$dG = VdP - SdT$$

- 1st law $dU = Q + W$

- Ideal gas: $PV = NkT$, equipartition, heat capacity, entropy of mixing, heat, work, state variables, path independence

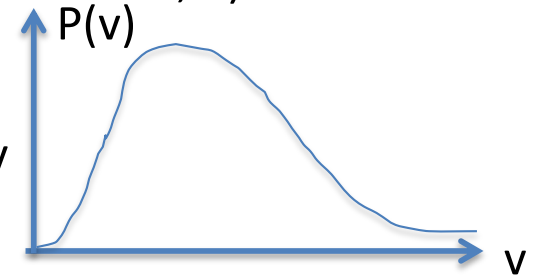
Experiments: measure and model

Probability of jumping people



To states: 1) On the floor, 2) on red box

MB velocity distr.



Mean kinetic energy

$$m\langle v^2 \rangle / 2 = kT$$

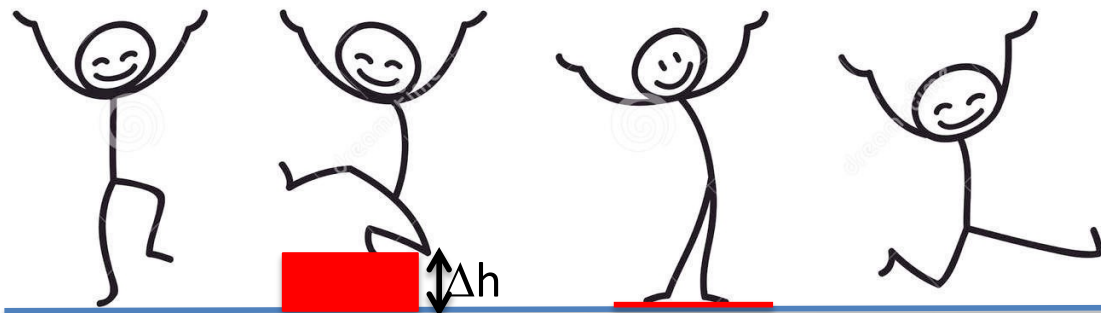
Mean potential energy

$$mg\langle h \rangle = m\langle v^2 \rangle / 2$$



Relative probability

$$\begin{aligned} P(2)/P(1) &= \exp(-E_2/kT) / \exp(-E_1/kT) \\ &= \exp(-(E_2 - E_1)/kT) \\ &= \exp(-\Delta E/kT) \\ &= \exp(-mg\Delta h/kT) \end{aligned}$$



Microcanonical & canonical ensemble

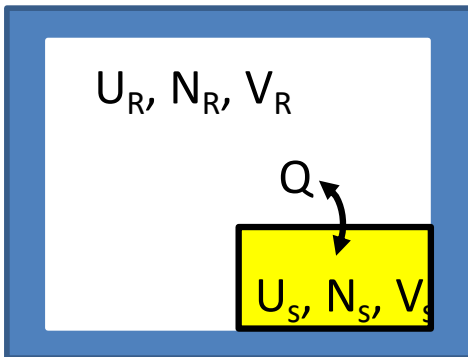


Microcanonical, (NVU) kept constant

Number of microstates in a macrostate: Multiplicity Ω

Probability of a macrostate: $P = \Omega / \sum \Omega$

Entropy: $S = k \ln \Omega$



Canonical, (NVT) kept constant

Exchanges Q with (NVU) reservoir to keep T constant

Boltzmann factor: $e^{-\beta \epsilon_i}, \beta = \frac{1}{kT}$

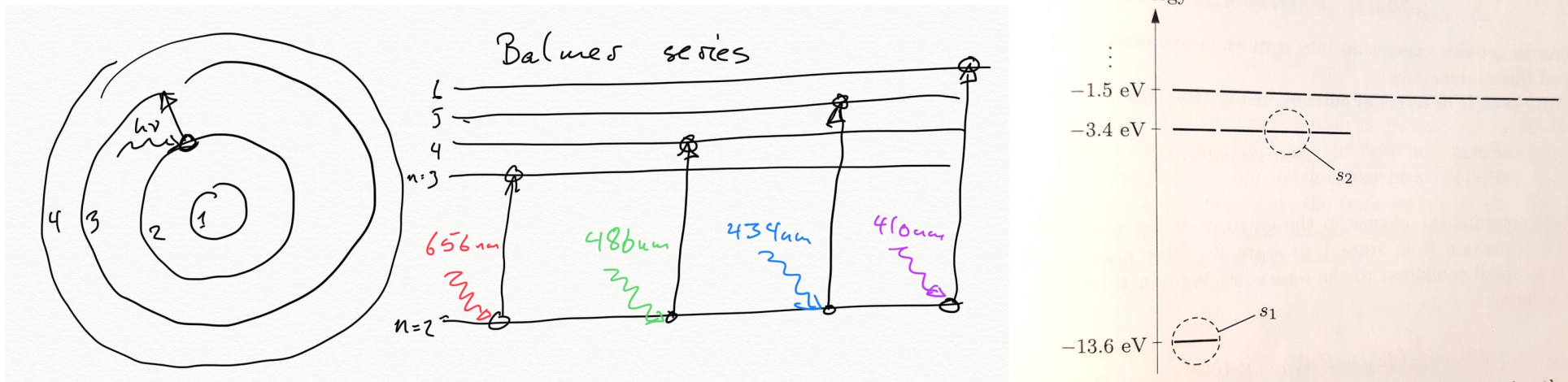
Partition function: sum over all possible microstates:

$$Z = \sum_i e^{-\beta \epsilon_i}$$

Probability of a microstate: $P_i = \frac{e^{-\beta \epsilon_i}}{Z}$

Relative probability: $\frac{P_2}{P_1} = e^{-\beta(\epsilon_2 - \epsilon_1)}$

Example: Hydrogen in the stars

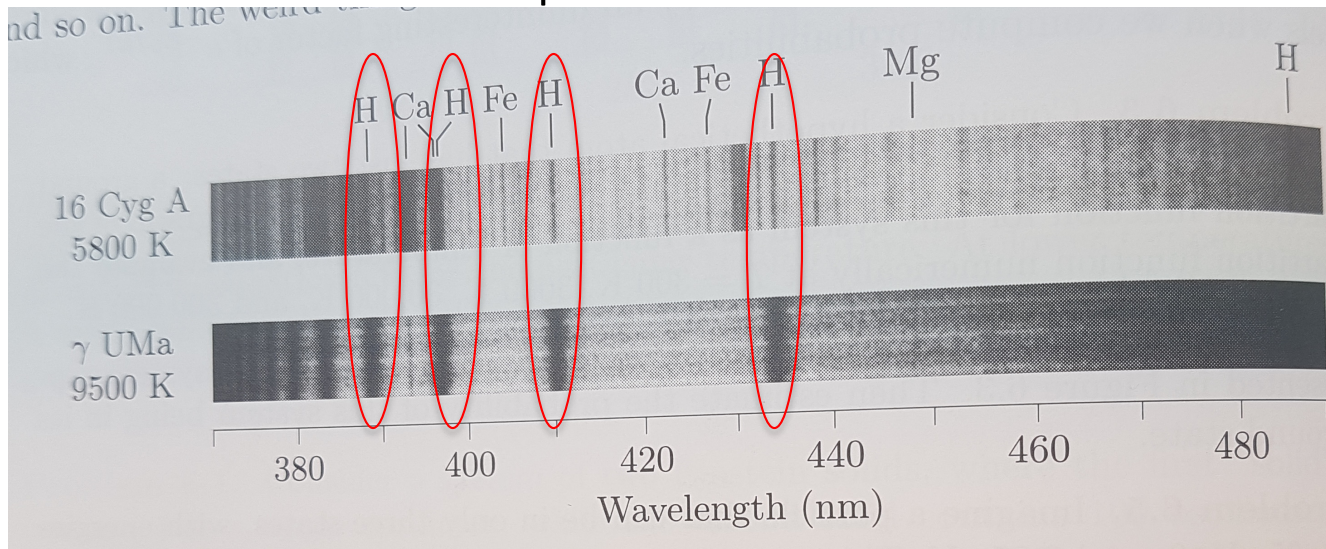


Probability of excited vs. ground state: $\frac{P(s_2)}{P(s_1)} = e^{-\beta(E_2 - E_1)}$ on sun ($T = 5800$ K)

$$\frac{(E_2 - E_1)}{kT} = \frac{13.6 - 3.4}{8.6 \cdot 10^{-5} \cdot 5800} = 20.4 \Rightarrow \frac{P(s_2)}{P(s_1)} = 1.4 \cdot 10^{-9}$$

Excited atoms absorb visible photons when $n = 2 \rightarrow i, \quad i > 2$

Absorption lines



5800 K, $P(s_2)/P(s_1) = 10^{-9}$

9500 K, $P(s_2)/P(s_1) = 10^{-6}$

Average values, Energy

Probability of a microstate: $P(s) = \frac{e^{-\beta E(s)}}{Z}$

Partition function $Z = \sum_s e^{-\beta E(s)}$

The average value of a fluctuating quantity is the first moment of its distribution:

$$\bar{X} = \langle X \rangle = \sum_s X(s)P(s)$$

For the Boltzmann distribution

$$\bar{X} = \frac{1}{Z(T)} \sum_s X(s)e^{-\beta E(s)}$$

Average energy $U = \bar{E} = \frac{1}{Z(T)} \sum_s E(s)e^{-\beta E(s)}$

$$\frac{\partial Z}{\partial \beta} = \sum_s \frac{\partial}{\partial \beta} e^{-\beta E(s)} = \sum_s -E(s)e^{-\beta E(s)}$$

=> $U = \bar{E} = -\frac{1}{Z} \frac{\partial Z}{\partial \beta}$

Thermodynamics: Helmholtz free energy

Internal energy

- ▶ $dU = TdS - PdV + \mu dN$

Helmholtz free energy

- ▶ $F \equiv U - TS$

- ▶ $dF = dU - TdS - SdT$

- ▶ Thermodynamic identity for F: $dF = -PdV - SdT + \mu dN$

- ▶ $dF = \left(\frac{\partial F}{\partial V}\right)_{N,T} dV + \left(\frac{\partial F}{\partial T}\right)_{N,V} dT + \left(\frac{\partial F}{\partial N}\right)_{V,T} dN$

- ▶ $P = -\left(\frac{\partial F}{\partial V}\right)_{N,T}$, $S = -\left(\frac{\partial F}{\partial T}\right)_{N,V}$, $\mu = \left(\frac{\partial F}{\partial N}\right)_{V,T}$

Natural variables: N, V, T, like the Canonical ensemble!

Average values: Helmholtz free energy

Microcanonical ensemble (NVU):

Fundamental variables: Ω, S

Micro to macro: $S = k \ln \Omega$

Second law: $\Delta S \geq 0$

Derived definitions: $T^{-1} \equiv \left(\frac{\partial S}{\partial U}\right)_{N,V}, P \equiv T \left(\frac{\partial S}{\partial V}\right)_{N,U}, \mu \equiv -T \left(\frac{\partial S}{\partial N}\right)_{U,V}$

Canonical ensemble (NVT):

Fundamental variables: $Z, F = U - TS$

Second law: $\Delta F \leq 0$

Derived definitions: $P = -\left(\frac{\partial F}{\partial V}\right)_{N,T}, S = -\left(\frac{\partial F}{\partial T}\right)_{N,V}, \mu = \left(\frac{\partial F}{\partial N}\right)_{T,V}$

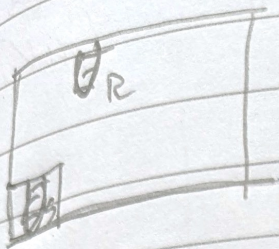
Micro to macro: $F = U + T \left(\frac{\partial F}{\partial T}\right)_{N,V}$

$$F = -kT \ln Z$$

$$\frac{\partial \ln Z}{\partial T} = \frac{\partial \beta}{\partial T} \frac{\partial \ln Z}{\partial \beta} = \frac{-1}{kT^2} \frac{1}{Z} \frac{\partial Z}{\partial \beta} = \frac{U}{kT^2}$$

$$\Rightarrow \frac{\partial F}{\partial T} = -\ln Z \frac{\partial kT}{\partial T} - kT \frac{\partial \ln Z}{\partial T} = \frac{F}{T} - \frac{U}{T}$$

F is minimum in equilibrium



Total energy $U = U_R + U_S$

entropy $S = S_R + S_S$

$$S = S_R(U - U_S) + S_S(U_S)$$

$$\approx S_R(U) + U_S \left(\frac{\partial S_R}{\partial U_R} \right)_{U, N} + S_S(U_S)$$

$$\Rightarrow S \approx S_R(U) - \frac{U_S}{T} + S_S(U_S)$$

$$TS = TS_R(U) - U_S + TS_S(U_S)$$

- F_S

Equilibrium U is constant $\Rightarrow S_R(U)$ is constant
 S is max $\Leftrightarrow F$ is minimum.

Paramagnetism

One fluctuating spin

– magnet. moment: μ

– spin: $s = \pm \frac{1}{2}$

– energy: $\epsilon_s = -2s\mu B$

Probability up:

$$P_{\uparrow} = \frac{1}{Z} e^{-\beta\epsilon_s} = \frac{1}{Z} e^{\beta\mu B}$$

Probability down:

$$P_{\downarrow} = \frac{1}{Z} e^{-\beta\epsilon_s} = \frac{1}{Z} e^{-\beta\mu B}$$

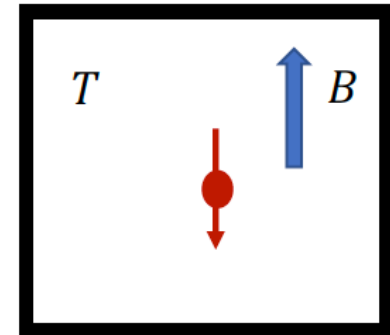
Partition function:

$$Z = e^{-\beta\mu B} + e^{\beta\mu B} = 2\cosh(\beta\mu B)$$

Average energy: $\bar{\epsilon} = -\frac{1}{Z} \frac{\partial Z}{\partial \beta} = -\frac{2\mu B \sinh \beta\mu B}{2 \cosh(\beta\mu B)} = -\mu B \tanh(\beta\mu B)$

Average mag. mom.: $\bar{\mu} = (+\mu)P_{\uparrow} + (-\mu)P_{\downarrow} = \frac{\mu}{Z} (e^{\beta\mu B} - e^{-\beta\mu B})$
 $= \mu \tanh(\beta\mu B)$

N spins, magnetization: $M = N\bar{\mu}$, energy: $U = -N\mu B \tanh(\beta\mu B)$



Equipartition theorem

Systems with quadratic degrees of freedom

$$E(q) = cq^2$$

Partition function $Z = \sum_q e^{-\beta cq^2} = \frac{1}{\Delta q} \sum_q \Delta q e^{-\beta cq^2}$

$\Delta q \ll kT,$ $= \frac{1}{\Delta q} \int_{-\infty}^{\infty} dq e^{-\beta cq^2} = \frac{1}{\Delta q} \sqrt{\frac{\pi}{\beta c}} = C\beta^{-\frac{1}{2}}$

Mean energy $\bar{E} = U = -\frac{1}{Z} \frac{\partial Z}{\partial \beta} = -\frac{\beta^{\frac{1}{2}}}{C} \left(-\frac{1}{2} C \beta^{-\frac{3}{2}} \right) = \frac{1}{2} kT$

Kinetic energy $\bar{E}_k = \frac{1}{2} m \overline{v^2} = \frac{1}{2} m (\overline{v_x^2} + \overline{v_y^2} + \overline{v_z^2}) = \frac{3}{2} kT$

RMS speed $v_{rms} = \sqrt{\overline{v^2}} = \sqrt{\frac{3kT}{m}}$

Partition function of many particles

Getting the counting right

1 particle, N states $Z_1 = \sum_s^N e^{-\beta \epsilon_s}$

Identical particles

Noninteracting / independent particles

no extra states are created due to having more particles

Distinguishable particles

2 particles, N states: total energy $\sum_i^2 \epsilon_{i,s} = \epsilon_{1,s} + \epsilon_{2,s}$

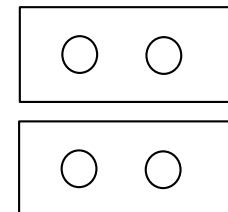
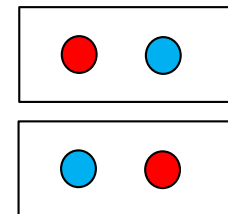
$$\begin{aligned} Z_2 &= \sum_s e^{-\beta(\epsilon_{1,s} + \epsilon_{2,s})} \\ &= \sum_s e^{-\beta \epsilon_{1,s}} e^{-\beta \epsilon_{2,s}}, \quad s = (s_1, s_2) \\ &= \sum_{s_1} \sum_{s_2} e^{-\beta \epsilon_{1,s_1}} e^{-\beta \epsilon_{2,s_2}} \\ &= Z_1 Z_1 = Z_1^2 \end{aligned}$$

$$Z_N = Z_1^N$$

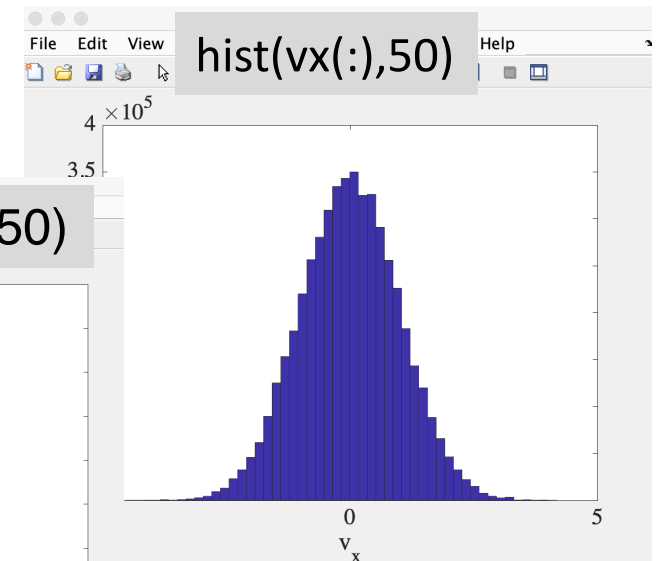
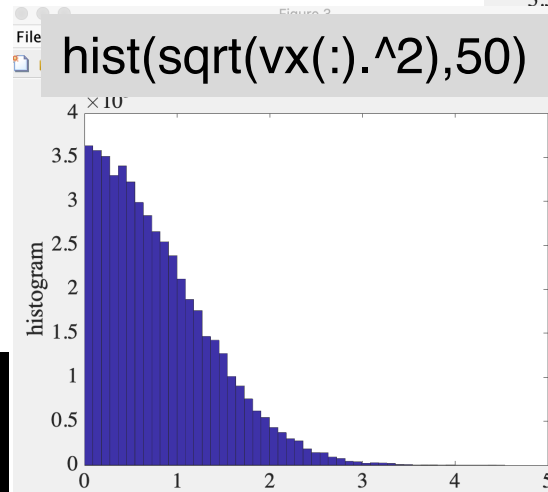
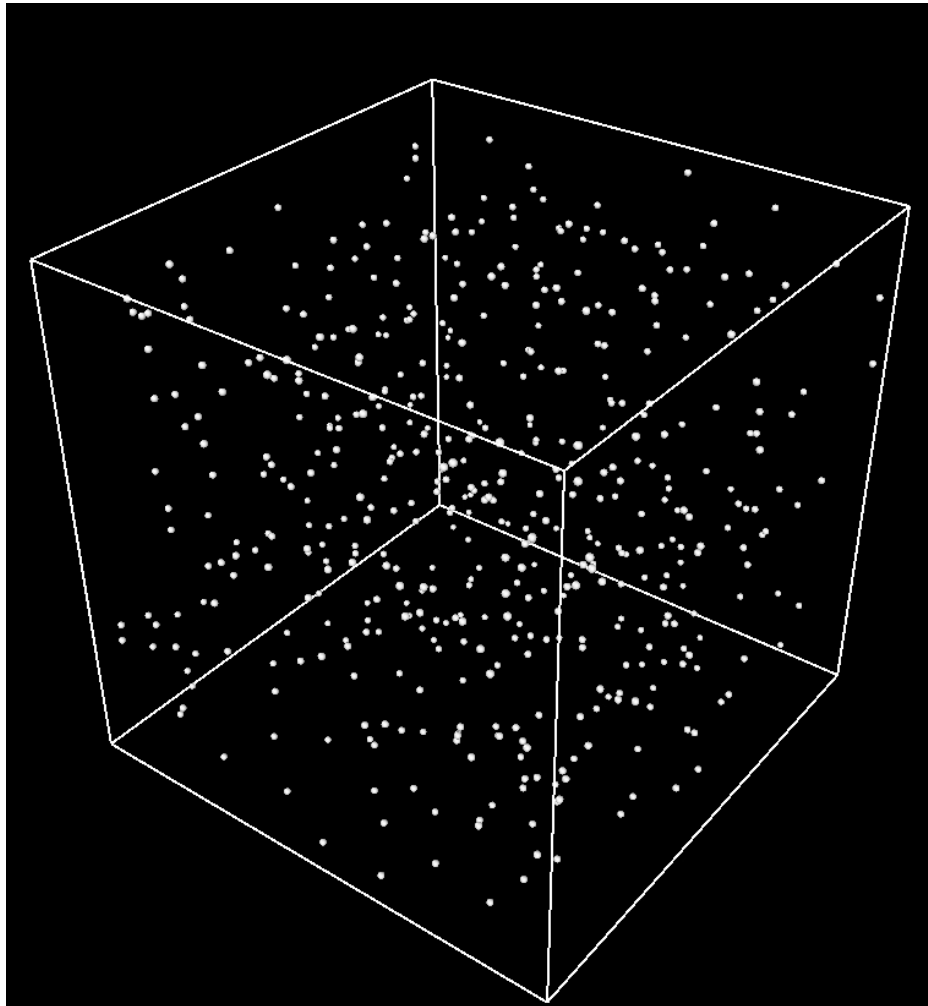
Indistinguishable particles

$$Z_2 = \frac{1}{2} Z_1^2$$

$$Z_N = \frac{1}{N!} Z_1^N$$

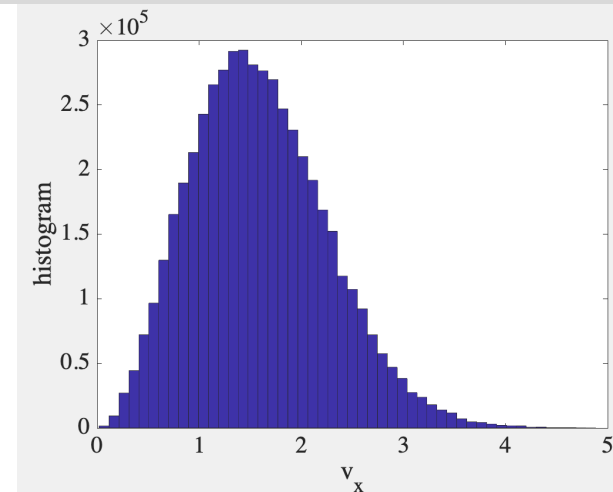


Maxwell-Boltzmann speed distribution



- Average velocity $\overline{v_x} = 0$
- Velocity is a vector, speed is a scalar
- Velocity is in 3D
- Probability of speed v , $P(v) = 0$
- Probability distribution: $P(v: v + dv) = D(v)dv$

hist(sqrt(vx(:).^2+vy(:).^2+vz(:).^2),50)




```
# Define variables
# Length of sides of cubic box
variable L equal 5
# Temperature:
variable T equal 1
# Density:
variable rho equal 0.001
# Pressure:
variable P equal 0.01
# Set Pconst=1 for pressure constant, Pconst=0 for constant volume
variable Pconst equal 0
```

```
units lj
dimension 3
lattice fcc ${rho}
region simbox block 0 $L 0 $L 0 $L
create_box 1 simbox
create_atoms 1 box
```

```
mass 1 1.0
velocity all create $T 87287
pair_style lj/cut 2.5
pair_coeff 1 1 1.0 1.0 2.5
```

```
if "${Pconst} == 0" then &
"fix nvt all nvt temp $T $T 1" &
else &
"fix npt all npt temp $T $T 1 iso $P $P 1"
neigh_modify every 1 delay 0 check yes
```

```
#Thermalization run
run 1000
```

```
if "${Pconst} == 0" then &
"unfix nvt" &
"fix nve all nve" &
else &
"unfix npt" &
"fix nph all nph iso $P $P 1"
```

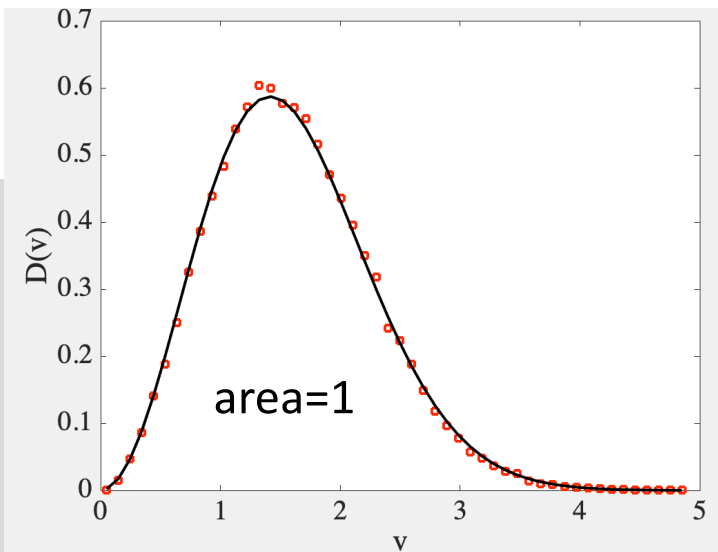
```
#Start heat addition and run to stabilize
variable eFlux equal 1
#fix heat all heat 1 ${eFlux} region simbox
run 1000
```

```
#Define new thermodynamic output and trajectory dump and start production run
thermo_style custom step temp epair etotal press
thermo 100
dump 1 all custom 100 dump.lammpstrj id x y z vx vy vz
run 1000000
```

```
$ Imp_serial < heatcapLJ.in
```

```
vsq=vx.^2+vy.^2+vz.^2;
speed=sqrt(vsq);
h = histogram(speed(:),50);
hv=h.Values;
s=sum(hv);
hbe=h.BinEdges;
dv=h.BinWidth;
hbc=hbe+dv/2;
v=hbc(1:end-1);
Dofv=hv/s/dv;
norm=sum(Dofv)*dv;
meanv=sum(v.*Dofv*dv)/norm;
vrms=sqrt(mean(vsq(:)))
plot(v,Dofv,'.',v,2*v.^2.*exp(-v.^2./2)./sqrt(2*pi),'-')
```

$$kT = 1, m = 1 \Rightarrow$$



v

$D(v)$

$$\sum D(v)dv = 1$$

$$\bar{v} = \sum vD(v)dv = 1.5966$$

$$v_{rms} = \sqrt{\langle v_x^2 + v_y^2 + v_z^2 \rangle} = 1.7287$$

$$D(v) = \frac{2}{\sqrt{2\pi}} v^2 e^{-\frac{v^2}{2}}$$

Distribution function: $[D(v)] = [v^{-1}]$

$$[P] = \left[\int D(v)dv \right] = 1$$

Maxwell-Boltzmann speed distribution

Probability distribution:

$$P(v: v + dv) = D(v)dv$$

Boltzmann probability:

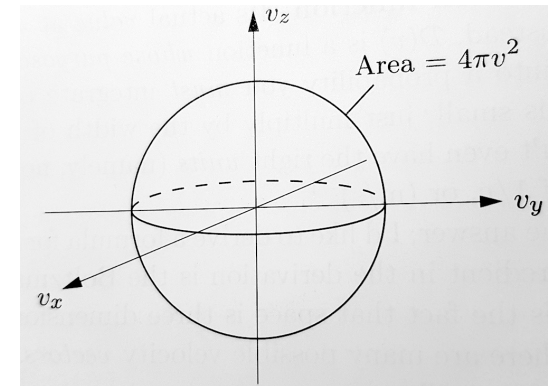
$$D(v)dv \propto e^{-\frac{E_k(v)}{kT}} dn = e^{-\frac{mv^2}{2kT}} dn$$

Number of vectors of speed v :

$$dn \propto 4\pi v^2 dv$$

Normalization

$$\int_0^\infty D(v)dv = 1$$



$$C \int_0^\infty e^{-\frac{mv^2}{2kT}} 4\pi v^2 dv = 1 \implies C = \left(\frac{m}{2\pi kT}\right)^{3/2}$$

M-B speed distribution

$$D(v) = \left(\frac{m}{2\pi kT}\right)^{3/2} 4\pi v^2 e^{-\frac{mv^2}{2kT}}$$

Maxwell-Boltzmann speed distribution

Average speed:

$$\bar{v} = \int_0^{\infty} vD(v)dv = \sqrt{\frac{8kT}{\pi m}} (= 1.5958)$$

$$\text{meanv}=\text{sum}(v.*Dofv*dv)= 1.5966$$

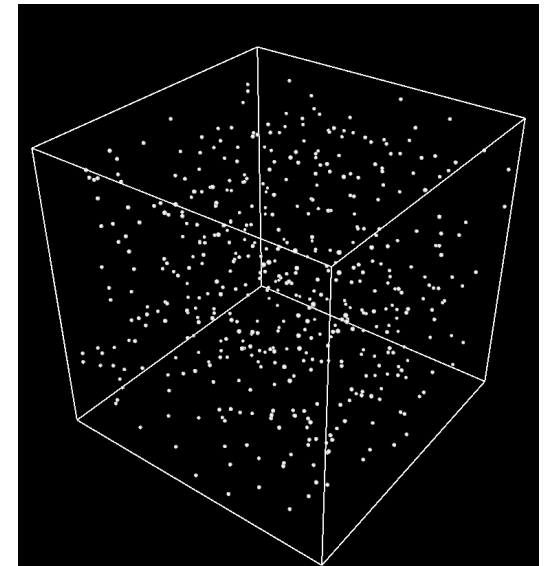
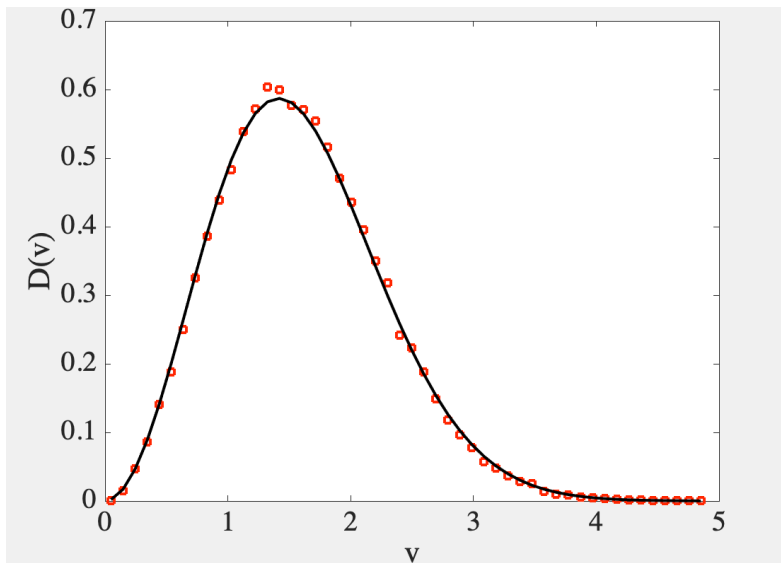
RMS speed:

$$v_{rms} = \sqrt{\int_0^{\infty} v^2 D(v)dv} = \sqrt{\frac{3kT}{m}} (= 1.7321)$$

$$v_{rms}=\text{sqrt}(\text{mean}(v\text{sq}(:)))=1.7287$$

Most probable speed

$$\frac{\partial D(v)}{\partial v} = 0 \Rightarrow v_{mp} = \sqrt{\frac{2kT}{m}} (= 1.4142)$$



Ensembles, counting and probabilities



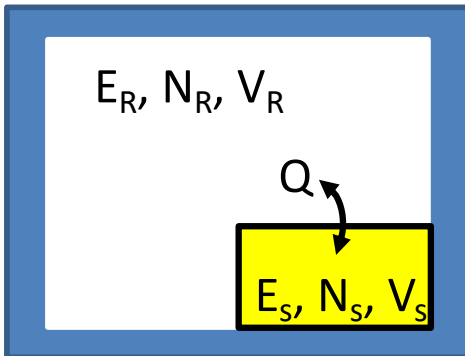
Microcanonical, (N, V, U) kept constant

Number of microstates, s , in a macrostate: Multiplicity Ω

Probability of macrostate: $P = \Omega / \sum \Omega$

Probability of a microstate for a given macrostate: $P(s) = 1/\Omega$

Entropy: $S = k \ln \Omega$



Canonical, (N, V, T) kept constant

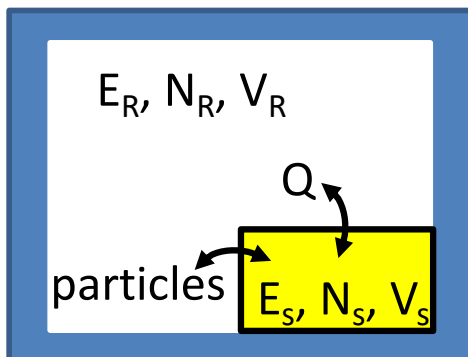
Exchanges Q with (N, V, U) reservoir to keep T constant

Boltzmann factor: $e^{-\beta \epsilon_i}$

Partition function: sum over all possible microstates: $Z = \sum_i e^{-\beta \epsilon_i}$

Probability of a microstate: $P_i = e^{-\beta \epsilon_i} / Z$

Free energy: Helmholtz, $F = -kT \ln Z$



Grand canonical, (μ, V, U) kept constant

Exchanges Q & particles with (N, V, U) reservoir to keep T & μ constant

Gibbs factor: $e^{-\beta(\epsilon_i - \mu N_i)}$

Gibbs sum: sum over all possible microstates: $Z_G = \sum_i e^{-\beta(\epsilon_i - \mu N_i)}$

Probability of a microstate: $P_i = e^{-\beta(\epsilon_i - \mu N_i)} / Z_G$

Free energy: Grand potential, $\Phi = -kT \ln Z_G = U - TS - \mu N$