

Proof that the force from an external magnetic field \vec{B} on a ^{static} current density $\vec{j}(\vec{r})$ is given by

$$\vec{F}_m = (\vec{m} \cdot \vec{\nabla}') \vec{B}(\vec{r}')|_{\vec{r}'=0} + \dots$$

We have shown in lecture that

$$\vec{F}_m = \int \vec{j}(\vec{r}) \times (\vec{r} \cdot \vec{\nabla}') \vec{B}(\vec{r}') d^3\vec{r} + \dots$$

Now we use that

$$(\vec{r} \cdot \vec{\nabla}') \vec{B}(\vec{r}') = \vec{\nabla}' (\vec{r} \cdot \vec{B}(\vec{r}')) - \vec{r} \times [\vec{\nabla}' \times \vec{B}(\vec{r}')]]$$

this can be seen from Rottmann p. 64

$$\vec{\nabla}(\vec{a} \cdot \vec{b}) = (\vec{b} \cdot \vec{\nabla}) \vec{a} + (\vec{a} \cdot \vec{\nabla}) \vec{b} + \vec{a} \times (\vec{\nabla} \times \vec{b}) + \vec{b} \times (\vec{\nabla} \times \vec{a})$$

with $\vec{a} = \vec{r}$ and $\vec{b} = \vec{B}(\vec{r}')$ and using that $\vec{\nabla}' \times \vec{r} = 0$ and $(\vec{B}(\vec{r}') \cdot \vec{\nabla}') \vec{r} = 0$ since the differentiation is wrt. \vec{r}' .

Further $\vec{\nabla}' \times \vec{B}(\vec{r}') = 0$ since \vec{B} is external and has no source near $\vec{j}(\vec{r})$ (Ampère's law)

$$\begin{aligned} \text{Thus } \int \vec{j}(\vec{r}) \times (\vec{r} \cdot \vec{\nabla}') \vec{B}(\vec{r}') d^3\vec{r} &= \int \vec{j}(\vec{r}) \times \vec{\nabla}' (\vec{r} \cdot \vec{B}(\vec{r}')) d^3\vec{r} \\ &= - \int \vec{\nabla}' \times (\vec{r} \cdot \vec{B}(\vec{r}')) \vec{j}(\vec{r}) d^3\vec{r} = - \vec{\nabla}' \times \int (\vec{r} \cdot \vec{B}(\vec{r}')) \vec{j}(\vec{r}) d^3\vec{r} \quad (*) \end{aligned}$$

We can now write (triple cross product rule, p. 62 Rottmann)

$$- \int (\vec{r} \cdot \vec{B}(\vec{r}')) \vec{j}(\vec{r}) d^3\vec{r} = \int \vec{B}(\vec{r}') \times (\vec{r} \times \vec{j}(\vec{r})) d^3\vec{r} - \int \vec{r} (\vec{B}(\vec{r}') \cdot \vec{j}(\vec{r})) d^3\vec{r} \quad (**)$$

$$\text{Furthermore } \int \vec{r} (\vec{B}(\vec{r}') \cdot \vec{j}(\vec{r})) d^3\vec{r} = - \int (\vec{r} \cdot \vec{B}(\vec{r}')) \vec{j}(\vec{r}) d^3\vec{r} \quad (***)$$

This follows from the same argument used to find the magnetic dipole contribution to the vector potential, namely that

$$\int \vec{r}' (\vec{j}(\vec{r}') \cdot \vec{r}') d^3\vec{r}' = - \int \vec{j}(\vec{r}') (\vec{r}' \cdot \vec{r}') d^3\vec{r}'$$

From (*), (**) and (***)

$$\vec{F}_m = \frac{1}{2} \vec{\nabla}' \times \int \vec{B}(\vec{r}') \times (\vec{r} \times \vec{j}(\vec{r})) d^3\vec{r} = \frac{1}{2} \vec{\nabla}' \times \vec{B}(\vec{r}') \times \int \vec{r} \times \vec{j}(\vec{r}) d^3\vec{r}$$

$$\begin{aligned}
&= \vec{\nabla}' \times \vec{B}(\vec{r}') \times \vec{m} = (\vec{m} \cdot \vec{\nabla}') \vec{B}(\vec{r}') - (\vec{B}(\vec{r}') \cdot \vec{\nabla}') \vec{m} + \vec{B}(\vec{r}') (\vec{\nabla}' \cdot \vec{m}) - \vec{m} (\vec{\nabla}' \cdot \vec{B}(\vec{r}')) \\
&= (\vec{m} \cdot \vec{\nabla}') \vec{B}(\vec{r}') - \vec{m} (\vec{\nabla}' \cdot \vec{B}(\vec{r}'))
\end{aligned}$$

Here we have used (Pottmann p. 64) that

$$\vec{\nabla} \times \vec{a} \times \vec{b} = (\vec{b} \cdot \vec{\nabla}) \vec{a} - (\vec{a} \cdot \vec{\nabla}) \vec{b} + \vec{a} (\vec{\nabla} \cdot \vec{b}) - \vec{b} (\vec{\nabla} \cdot \vec{a})$$

with $\vec{a} = \vec{B}(\vec{r}')$ and $\vec{b} = \vec{m}$ and the fact that \vec{m} is a constant.

Since $\vec{\nabla}' \cdot \vec{B}(\vec{r}') = 0$ (no monopoles)

$$\vec{F}_m = (\vec{m} \cdot \vec{\nabla}') \vec{B}(\vec{r}')$$

q.e.d.