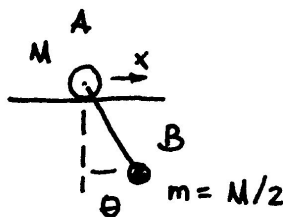


Exam FYS 3120/4120 Spring semester 2006

Solutions

Problem 1



$$I = \frac{1}{2} M R^2$$

$$\dot{x} = \omega R$$

(rolling cylinder)

a) Kinetic energy

$$T_A = \frac{1}{2} M \dot{x}^2 + \frac{1}{2} I \omega^2 = \frac{3}{4} M \dot{x}^2$$

$$T_B = \frac{1}{2} m \left( \underbrace{(\dot{x}^2 + L \cos \theta \dot{\theta})^2}_{\dot{x}_B^2} + \underbrace{L^2 \sin^2 \theta \dot{\theta}^2}_{\dot{y}_B^2} \right) = \frac{1}{4} M (\dot{x}^2 + L^2 \dot{\theta}^2 + 2 L \cos \theta \dot{\theta} \dot{x})$$

Potential energy

$$V_A = 0, \quad V_B = -mgL \cos \theta = -\frac{M}{2} g L \cos \theta$$

Lagrangian

$$L = T - V = \underline{M \left( \dot{x}^2 + \frac{1}{4} L^2 \dot{\theta}^2 + \frac{1}{2} L \cos \theta \dot{\theta} \dot{x} + \frac{1}{2} g L \cos \theta \right)}$$

b)  $\frac{\partial L}{\partial x} = 0 \Rightarrow x$  cyclic  $\Rightarrow \frac{d}{dt} \left( \frac{\partial L}{\partial \dot{x}} \right) = 0$

$$K \equiv \frac{\partial L}{\partial \dot{x}} = \underline{M \left( 2\dot{x} + \frac{1}{2} L \cos \theta \dot{\theta} \right)} \text{ is constant of motion}$$

$$\frac{\partial L}{\partial t} = 0 \Rightarrow \frac{dH}{dt} = 0$$

$$\Rightarrow H = T + V = \underline{M \left( \dot{x}^2 + \frac{1}{4} L^2 \dot{\theta}^2 + \frac{1}{2} L \cos \theta \dot{\theta} \dot{x} - \frac{1}{2} g L \cos \theta \right)} \equiv E \text{ (energy)}$$

is constant of motion

c)  $\frac{\partial L}{\partial \theta} = M \left( \frac{1}{2} L^2 \dot{\theta} + \frac{1}{2} L \cos \theta \dot{x} \right) \quad \frac{\partial L}{\partial \theta} = M \left( -\frac{1}{2} L \sin \theta \dot{\theta} \dot{x} + \frac{1}{2} g L \sin \theta \right)$

$$\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{\theta}} \right) - \frac{\partial L}{\partial \theta} = 0 \Rightarrow \underline{\ddot{\theta} + \frac{1}{L} \cos \theta \dot{x} + \frac{g}{L} \sin \theta = 0}$$

\*equation \*discussed under b)

$$\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{x}} \right) = M (2 \ddot{x} + \frac{1}{2} L \cos \theta \ddot{\theta} - \frac{1}{2} L \sin \theta \dot{\theta}^2) = 0$$

$$\Rightarrow \ddot{x} = -\frac{1}{4} L \cos \theta \ddot{\theta} + \frac{1}{4} L \sin \theta \dot{\theta}^2$$

Inserted in the  $\theta$ -equation:

$$(1 - \frac{1}{4} \cos^2 \theta) \ddot{\theta} + \frac{1}{4} \sin \theta \cos \theta \dot{\theta}^2 + \frac{g}{l} \sin \theta = 0$$

Small angle approximation  $\theta \ll 1 \Rightarrow \cos \theta \approx 1, \sin \theta \approx \theta$

Approximation to first order in  $\theta, \dot{\theta}, \ddot{\theta}$ :

$$(1 - \frac{1}{4}) \ddot{\theta} + \frac{g}{l} \theta = 0 \Rightarrow \underline{\ddot{\theta} + \frac{4}{3} \frac{g}{l} \theta = 0}$$

Harmonic osc. equation with frequency  $\omega = \sqrt{\frac{4}{3} \frac{g}{l}}$

General solution  $\theta(t) = A \cos(\omega t - \varphi)$

Initial conditions  $\theta(0) = \theta_0, \dot{\theta}(0) = 0 \Rightarrow \underline{\theta(t) = \theta_0 \cos \omega t}$

Solution for the  $x$ -coordinate

$$\dot{x}(0) = \dot{\theta}(0) = 0 \Rightarrow K = 0$$

$$\Rightarrow \dot{x} = -\frac{1}{4} L \cos \theta \dot{\theta} \approx -\frac{1}{4} L \dot{\theta}$$

$$\Rightarrow x = -\frac{1}{4} L \theta + x_0 \quad \text{const}$$

$$t=0 \quad x_0 = \frac{1}{4} L \theta_0 \Rightarrow x(t) = +\frac{1}{4} L (\theta_0 - \theta(t))$$

$$= \frac{1}{4} L \theta_0 (1 - \cos \omega t) = \underline{\underline{\frac{1}{2} L \theta_0 \sin^2 \frac{\omega t}{2}}}$$

## Problem 2

a) Equation of motion

$$m_e (\gamma \vec{a} + \dot{\gamma} \vec{v}) = e \vec{E}$$

Instantaneous inertial frame moves along  $\vec{E}$ :  
no transformation of  $\vec{E}$

In this frame:  $\vec{v} = 0$ ,  $\gamma = 1$ ,  $\dot{\gamma} = 0$ ,  $\vec{a} \equiv \vec{a}_0$

$$\Rightarrow \underline{m_e \vec{a}_0 = e \vec{E}}$$

b)  $\frac{d}{dt} (\gamma \vec{v}) = \frac{e \vec{E}}{m_e} \Rightarrow \gamma \vec{v} = \frac{e \vec{E}}{m_e} t + \text{const} \quad \swarrow = 0$

$$\Rightarrow \gamma^2 \beta^2 = \left( \frac{eE}{m_e c} \right)^2 t^2 \quad \beta = \frac{v}{c}$$

$$\gamma^2 = \frac{1}{1-\beta^2} \Rightarrow \gamma^2 \beta^2 = \gamma^2 - 1$$

$$\gamma^2 = \gamma^2 \beta^2 + 1 = \left( \frac{eE}{m_e c} \right)^2 t^2 + 1$$

$$\Rightarrow \gamma = \sqrt{1 + \kappa^2 t^2} \quad \underline{\kappa = \frac{eE}{m_e c} = \frac{a_0}{c}}$$

c)

$$\gamma = \cosh \kappa \tau$$

$$\Rightarrow \kappa^2 t^2 = \cosh^2 \kappa \tau - 1 = \sinh^2 \kappa \tau \Rightarrow t = \frac{1}{\kappa} \sinh \kappa \tau$$

$$\frac{dt}{d\tau} = \cosh \kappa \tau = \underline{\gamma}$$

Compare definition of proper time

$$d\tau^2 = dt^2 - \frac{1}{c^2} dx^2 = \left(1 - \frac{v^2}{c^2}\right) dt^2 = \frac{1}{\gamma^2} dt^2$$

$$\text{or } \underline{\frac{dt}{d\tau} = \gamma}$$

### Problem 3

a) Electric dipole radiation (formula coll.)

$$\vec{B}(\vec{r}, t) = \frac{\mu_0}{4\pi cr} \ddot{\vec{p}}(t_{\text{ret}}) \times \vec{n}, \quad t_{\text{ret}} = t - \frac{r}{c}, \quad \vec{n} = \frac{\vec{r}}{r}$$

$$\ddot{\vec{p}} = -\omega^2 \vec{p} = -\omega^2 l q (\cos \omega t \vec{i} + \sin \omega t \vec{j})$$

$$\vec{n} = \sin \theta \cos \varphi \vec{i} + \sin \theta \sin \varphi \vec{j} + \cos \theta \vec{k} \quad \text{polar coord.}$$

$$\Rightarrow \vec{B}(\vec{r}, t) = -\frac{\mu_0 \omega^2 l q}{4\pi cr} (\cos \theta \sin \omega t_{\text{ret}} \vec{i} - \cos \theta \cos \omega t_{\text{ret}} \vec{j} - \sin \theta \sin(\omega t_{\text{ret}} - \varphi) \vec{k})$$

$$\uparrow \quad \underline{B_0} = -\frac{\mu_0 \omega^2 l q}{4\pi cr}$$

For radiation fields (formula coll.)  $\underline{\vec{E}(\vec{r}, t) = c \vec{B}(\vec{r}, t) \times \vec{n}}$

b) Radiation in the x-direction;  $\vec{n} = \vec{i}$  ( $\theta = \frac{\pi}{2}, \varphi = 0$ ):

$$\vec{B}(\vec{r}, t) = -B_0 \sin \omega t_{\text{ret}} \vec{k}, \quad \text{oscillates in the z-direction}$$

$\Rightarrow$  linear polarization

Radiation in the z-direction;  $\vec{n} = \vec{k}$  ( $\theta = 0$ )

$$\vec{B}(\vec{r}, t) = B_0 (\sin \omega t_{\text{ret}} \vec{i} - \cos \omega t_{\text{ret}} \vec{j}), \quad \text{rotates in the x, y-plane}$$

$\Rightarrow$  circular polarization

c) Energy density

$$u = \frac{1}{2} (\epsilon_0 \vec{E}^2 + \frac{1}{\mu_0} \vec{B}^2) = \frac{1}{\mu_0} \vec{B}^2 = \frac{1}{\mu_0} B_0^2 (\cos^2 \theta + \sin^2 \theta \sin^2(\omega t_{\text{ret}} - \varphi))$$

$$\text{time average } \overline{\sin^2 \omega t} = \frac{1}{2} \Rightarrow \bar{u} = \frac{1}{2\mu_0} B_0^2 (1 + \cos^2 \theta)$$

for fixed  $r$  maximal for  $\cos \theta = \pm 1$ ; along the z-axis

This is the direction of maximal radiation, since radiated energy is proportional to  $u$ :

$$\vec{S} = \frac{1}{\mu_0} \vec{E} \times \vec{B} = \frac{c}{\mu_0} B^2 \vec{n} = \underline{c u \vec{n}}$$