

## Orthogonal sets of functions

Recall our standard homogeneous differential equation:

$$y'' + P(x)y' + Q(x)y = 0.$$

Now, there is a special class of such equations that take the form

$$\frac{d}{dx}[p(x)y'] + [q(x) + \lambda r(x)]y = 0, \text{ with } r(x) > 0. \text{ Thus,}$$

$$p(x)y'' + p'(x)y' + [q(x) + \lambda r(x)]y = 0.$$

### Examples:

Legendre:  $(1-x^2)y'' - 2xy' + n(n+1)y = 0$ , ( $\lambda = n(n+1)$ ,  $p(x) = 1-x^2$ ,  $q(x) = 0$ , and  $r(x) = 1$ ).

Fourier:  $y'' + \left(\frac{n\pi}{L}\right)^2 y = 0$ ,  $p(x) = r(x) = 1$ ,  $q(x) = 0$ .

Hermite:  $y'' - 2xy' + 2ny = 0$ , multiplication by  $e^{-x^2}$  gives it the right form:

$$e^{-x^2}y'' - 2xe^{-x^2}y' + 2ne^{-x^2}y = 0, \text{ with } p(x) = r(x) = e^{-x^2}, \text{ and } q(x) = 0, \lambda = 2n.$$

Laguerre:  $xy'' + (1-x)y' + \lambda y = 0$ , multiplied by  $e^{-x}$  to get the correct form:

$$xe^{-x}y'' + (1-x)e^{-x}y' + \lambda e^{-x}y = 0, \text{ with } p(x) = xe^{-x}, r(x) = e^{-x}, \text{ and } q(x) = 0.$$

We are actually investigating so called eigenvalue equations of the general form:

$$Dy + \lambda r(x)y = 0, \text{ where the differential operator } D \text{ has the form}$$

$$D = p(x)\frac{d^2}{dx^2} + p'(x)\frac{d}{dx} + q(x).$$

Boundary conditions are applied at  $x = a$  and  $x = b$ , i.e.  $x \in [a, b]$ .

Legendre:  $x \in [-1, 1]$ ,

Fourier:  $x \in [-L, L]$ ,

Hermite:  $x \in \langle -\infty, \infty \rangle$ ,

Laguerre:  $x \in [0, \infty > .$

Notice that for all the equations listed above we have  $p(a) = p(b)$  (at the boundaries).

The boundary conditions at  $x = a$  and  $x = b$  typically lead to a series of discrete eigenvalues  $\lambda_n$ , and corresponding solutions (eigenfunctions)  $y_n(x)$ .

Without proof (which may be rather demanding!) we state that the eigenfunctions are orthogonal:

$$\int_a^b r(x) y_n(x) y_m(x)^* dx = 0 \text{ for } \lambda_n \neq \lambda_m .$$

### Complete set of eigenfunctions:

If any function  $f(x)$  (without infinite discontinuities) can be expanded in a convergent series in terms of the set of eigenfunctions  $\{y_n(x)\}$  for  $x \in [a, b]$  (or the appropriate open interval):

$$f(x) = \sum_{n=1}^{\infty} a_n y_n(x),$$

then we say that the set  $\{y_n(x)\}$  is complete.

The coefficients  $a_n$  are determined by utilizing the orthogonality of the eigenfunctions:

$$\int_a^b f(x) r(x) y_m(x)^* dx = \sum_n a_n \int_a^b r(x) y_n(x) y_m(x)^* dx = a_m \int_a^b r(x) y_m(x) y_m(x)^* dx .$$

It is usual to normalize the eigenfunctions, i.e. multiply them by a constant  $C(n)$  ( $y_n(x)$  is replaced by  $C(n)y_n(x)$ ), so that

$$\int_a^b r(x) y_m(x) y_m(x)^* dx = 1. \quad \text{Thus,}$$

$$a_m = \int_a^b f(x) r(x) y_m(x)^* dx .$$

### Hermitian operators

For which type of eigenvalue equations and boundary conditions do the eigenfunctions form a complete orthogonal set?

If the differential operator  $D$  with given boundary conditions at  $x = a$  and  $x = b$  obeys the condition

$$\int_a^b y_n(x)^* Dy_m(x) dx = \int_a^b y_m(x) Dy_n(x)^* dx ,$$

then  $D$  is a Hermitian operator.

For Hermitian operators it follows (without proof here):

1. The eigenvalues are real (cf. observable quantities in quantum mechanics)
2. The eigenfunctions are orthogonal (over  $x \in [a, b]$ )
3. The eigenfunctions form a complete set for  $x \in [a, b]$ .

The eigenvalue equations listed as examples above, are all of the Hermitian type.

**Examples:**

Legendre equation:  $(1-x^2)y_n'' - 2xy_n' + n(n+1)y_n = 0$ .

Solutions: Legendre polynomials  $y_n(x) = P_n(x)$  for  $x \in [-1, 1]$

Orthogonality:  $\int_{-1}^1 P_n(x)P_m(x)dx = \frac{2}{2n+1}\delta_{nm}$ .

Completeness:  $f(x) = \sum_{n=0}^{\infty} a_n P_n(x)$ ,  $x \in [-1, 1]$ .  $a_n = \frac{2n+1}{2} \int_{-1}^1 f(x)P_n(x)dx$ .

Hermite equation:  $y_n'' - 2xy_n' + 2ny_n = 0$ ,  $r(x) = e^{-x^2}$ .

Solutions: Hermite polynomials:  $y_n(x) = H_n(x)$  for  $x \in \langle -\infty, \infty \rangle$ .

Orthogonality:  $\int_{-\infty}^{\infty} e^{-x^2} H_n(x)H_m(x)dx = 2^n n! \sqrt{\pi} \delta_{nm}$ .

Completeness:  $f(x) = \sum_{n=0}^{\infty} a_n H_n(x)$ ,  $a_n = \frac{1}{2^n n! \sqrt{\pi}} \int_{-\infty}^{\infty} e^{-x^2} f(x)H_n(x)dx$ .