

Problem set 9

9.1 The canonical commutation relations.

The basic commutation relations of the electromagnetic field can be written as

$$[\hat{A}_{\mathbf{k}a}, \hat{E}_{\mathbf{k}'a'}^\dagger] = -i \frac{\hbar}{\epsilon_0} \delta_{\mathbf{k}\mathbf{k}'} \delta_{aa'} \quad (1)$$

where \mathbf{k} is the wave vector of the electromagnetic field and a is a polarization index. The photon creation and annihilation operators are related to the field operators by

$$\hat{A}_{\mathbf{k}a} = \sqrt{\frac{\hbar}{2\omega_k \epsilon_0}} (\hat{a}_{\mathbf{k}a} + \hat{a}_{-\mathbf{k}\bar{a}}^\dagger), \quad \hat{E}_{\mathbf{k}a} = i \sqrt{\frac{\hbar\omega_k}{2\epsilon_0}} (\hat{a}_{\mathbf{k}a} - \hat{a}_{-\mathbf{k}\bar{a}}^\dagger) \quad (2)$$

where \bar{a} is related to a by a permutation of the two values of the polarization index. Use the above expressions to show that the creation and annihilation operators satisfy the standard harmonic oscillator commutation relations,

$$\begin{aligned} [\hat{a}_{\mathbf{k}a}, \hat{a}_{\mathbf{k}'a'}^\dagger] &= \delta_{\mathbf{k}\mathbf{k}'} \delta_{aa'} \\ [\hat{a}_{\mathbf{k}a}, \hat{a}_{\mathbf{k}'a'}] &= [\hat{a}_{\mathbf{k}a}^\dagger, \hat{a}_{\mathbf{k}'a'}^\dagger] = 0 \end{aligned} \quad (3)$$

9.2 The electromagnetic field energy and momentum.

The classical expressions for the electromagnetic field energy \mathcal{E} and field momentum \mathcal{P} are

$$\begin{aligned} \mathcal{E} &= \frac{1}{2} \int d^3r (\epsilon_0 \mathbf{E}^2 + \frac{1}{\mu_0} \mathbf{B}^2) \\ \mathcal{P} &= \epsilon_0 \int d^3r \mathbf{E} \times \mathbf{B} \end{aligned} \quad (4)$$

The same expressions are valid in the quantum description, when the classical fields are replaced by the operator fields, $\mathbf{E} \rightarrow \hat{\mathbf{E}}$, $\mathbf{B} \rightarrow \hat{\mathbf{B}}$. The energy \mathcal{E} is then replaced by the hamiltonian \hat{H} and the momentum \mathcal{P} by the momentum operator $\hat{\mathcal{P}}$.

Use expressions for the field operators from the lecture notes to show that \hat{H} and $\hat{\mathcal{P}}$ has the following form, when written in terms of the photon number $\hat{N}_{\mathbf{k}a}$ operator

$$\begin{aligned} \hat{H} &= \sum_{\mathbf{k}a} \hbar\omega_k \hat{N}_{\mathbf{k}a} + E_0 \mathbb{1} \\ \hat{\mathcal{P}} &= \sum_{\mathbf{k}a} \hbar\mathbf{k} \hat{N}_{\mathbf{k}a} \end{aligned} \quad (5)$$

where $\hat{N}_{\mathbf{k}a}$ is defined by

$$\hat{N}_{\mathbf{k}a} = \hat{a}_{\mathbf{k}a}^\dagger \hat{a}_{\mathbf{k}a} \quad (6)$$

Comment on the interpretation of the constant E_0 and explain the meaning of removing this from the definition of the Hamiltonian, as one usually does.

9.2 The electromagnetic field equation

Show that the field operator $\hat{\mathbf{A}}(\mathbf{r}, t)$, in the Heisenberg picture, satisfies the same field equation as the classical vector potential

$$\left(\frac{\partial^2}{\partial t^2} - c^2 \nabla^2\right)\hat{\mathbf{A}}(\mathbf{r}, t) = 0 \quad (7)$$

9.3 Polarization vectors

Show that the polarization vectors $\epsilon_{\mathbf{k}a}$ satisfy, together with the wave vector \mathbf{k} , the following completeness relation,

$$\sum_a (\epsilon_{\mathbf{k}a})_i (\epsilon_{\mathbf{k}a}^*)_j + \frac{k_i k_j}{\mathbf{k}^2} = \delta_{ij} \quad (8)$$

9.4 Field commutators

The electric and magnetic fields do not commute, but rather satisfy a commutation relation that can be written as

$$[\hat{E}_i(\mathbf{r}), \hat{B}_j(\mathbf{r}')] = i \frac{\hbar}{\epsilon_0} \sum_{\mathbf{k}} \epsilon_{ijk} \frac{\partial}{\partial x_k} \delta(\mathbf{r} - \mathbf{r}') \quad (9)$$

with ϵ_{ijk} as the three-dimensional Levi-Civita symbol. Show this by considering the Fourier transform of the equation.

To make precise what is meant by derivatives of delta-functions, one usually consider the corresponding Fourier transformed functions. As a reminder, the Fourier transform of the three-dimensional delta function in a finite volume V is

$$\delta(\mathbf{r}) = \frac{1}{V} \sum_{\mathbf{k}} e^{i\mathbf{k}\cdot\mathbf{r}} \quad (10)$$