### FYS 4110 Non-relativistic Quantum Mechanics, Fall Semester 2011

# Problem set 5

#### 5.1 Displacement operators in phase space

For a particle moving in one dimension the position coordinate x and the momentum p define the coordinates of the two-dimensional classical *phase space*.

In the quantum description of the one-dimensional harmonic oscillator non-hermitian lowering operator is defined as

$$\hat{a} = \frac{1}{\sqrt{2m\hbar\omega}} (m\omega\,\hat{x} + i\hat{p}) \tag{1}$$

We may consider this as the operator of a complex phase space variable, with position as the real part and momentum as the imaginary part. It has a dimensionless form due to the constants introduced in the expression.

A coherent state, in a similar way is characterized by a complex number z, the eigenvalue of  $\hat{a}$ , which we may interpret as a complex phase space coordinate,

$$z = \frac{1}{\sqrt{2m\hbar\omega}} (m\omega x_c + ip_c) \tag{2}$$

The following operator

$$\hat{\mathcal{D}}(z) = e^{(z\hat{a}^{\dagger} - z^*\hat{a})} \tag{3}$$

acts as a displacement operator in phase space, in the sense

$$\hat{\mathcal{D}}(z)\hat{x}\hat{\mathcal{D}}(z)^{\dagger} = \hat{x} - x_c , \quad \hat{\mathcal{D}}(z)\hat{p}\hat{\mathcal{D}}(z)^{\dagger} = \hat{p} - p_c \tag{4}$$

Show that displacements in two different directions in general will not commute but rather satisfy a relation of the form

$$\hat{\mathcal{D}}(z_a)\hat{\mathcal{D}}(z_b) = e^{i\alpha(z_a, z_b)}\hat{\mathcal{D}}(z_b)\hat{\mathcal{D}}(z_a)$$
(5)

with  $\alpha(z_a, z_b)$  as a complex phase. Determine the phase as a function of  $z_a$  and  $z_b$ . What is the condition for the two operators to commute?

### 5.2 Energy eigenstates in the coherent state representation

The energy eigenstates of a harmonic oscillator are denoted by  $|n\rangle$  and the coherent states by  $|z\rangle$ . We consider states described as wave functions in the coherent state representation,  $\psi(z) = \langle z | \psi \rangle$ .

In the coherent state representation the energy eigenstates have the form of wave functions,  $\psi_n(z) = \langle z|n \rangle = \exp(-|z|^2/2)[z^{*2}/\sqrt{n!}]$ . Show that the modulus squared of the energy eigenfunctions,  $|\psi_n(z)|^2$ , only depends on the absolute value of  $z, r \equiv |z|$ , and find the position  $r_n$  of its maximum. Plot the functions  $|\psi_n(z)|^2$  as function of r for values of n from 0 to 5.

## 5.3 Eigenvectors for $\hat{a}^{\dagger}$ ?

The coherent states  $|z\rangle$  are defined as eigenvectors of the lowering operator  $\hat{a}$ . Assume  $|\bar{z}\rangle$  to be eigenvector of the raising operator  $\hat{a}^{\dagger}$ ,

$$\hat{a}^{\dagger}|\bar{z}\rangle = \bar{z}|\bar{z}\rangle \tag{6}$$

Show that no normalizable vector exists that satisfies this equation by expanding the state  $|\bar{z}\rangle$  in the energy eigenstates  $|n\rangle$ .

## 5.4 Driven harmonic oscillator (Midterm Exam 2008)

The Hamiltonian of a one-dimensional harmonic oscillator is given by the expression

$$\hat{H} = \frac{1}{2m} (\hat{p}^2 + m^2 \omega_0^2 \hat{x}^2) = \hbar \omega_0 (\hat{a}^\dagger \hat{a} + \frac{1}{2})$$
(7)

with the raising and lowering operators defined by

$$\hat{a} = \frac{1}{\sqrt{2m\hbar\omega_0}} (m\omega_0 \hat{x} + i\hat{p}) , \quad \hat{a}^{\dagger} = \frac{1}{\sqrt{2m\hbar\omega_0}} (m\omega_0 \hat{x} - i\hat{p})$$
(8)

The time evolution operator is

$$\hat{\mathcal{U}}_0(t) = e^{-\frac{i}{\hbar}t\hat{H}_0} \tag{9}$$

The coherent states of the oscillator are defined as eigenvectors of the lowering operator

$$\hat{a}|z\rangle = z|z\rangle \tag{10}$$

The general coherent state  $|z\rangle$  is related to the ground state of the oscillator  $|0\rangle$  by

$$|z\rangle = \hat{\mathcal{D}}(z)|0\rangle = e^{-z^* z} e^{z\hat{a}^{\dagger}}|0\rangle \tag{11}$$

where the unitary shift operator is given by

$$\hat{\mathcal{D}}(z) = e^{z\hat{a}^{\dagger} - z^*\hat{a}} \tag{12}$$

a) Show that for a general operator  $\hat{A}$  we have the relation

$$\hat{\mathcal{U}}e^{\hat{A}}\hat{\mathcal{U}}^{-1} = e^{\hat{\mathcal{U}}\hat{A}\hat{\mathcal{U}}^{-1}} \tag{13}$$

and use that to calculate the operator  $\hat{\mathcal{U}}_0(t)\hat{\mathcal{D}}(z)\hat{\mathcal{U}}_0(t)^{\dagger}$ . Make use of the result to determine the time dependent state vector  $|\psi(t)\rangle$ , when this initially is a coherent state  $|\psi(0)\rangle = |z_0\rangle$ . Show that  $|\psi(t)\rangle$  at later times t is also a coherent state.

We next assume the harmonic oscillator to be under influence of a time dependent external potential, so that the hamiltonian now is

$$\hat{H} = \hat{H}_0 + \hat{W}(x, t)$$
 (14)

In the following we assume the external potential to have the specific form

$$\hat{W}(x,t) = A\hat{x}\sin\omega t \tag{15}$$

with A as a constant and  $\omega$  as the oscillation frequency of the external potential.

b) Find the Heisenberg equation of motion for  $\hat{x}$  and  $\hat{p}$  and show that they correspond to the equation of motion of a *driven* harmonic oscillator, that is subject to the periodic force  $f(t) = -A \sin \omega t$ .

c) Give the definition of the time evolution operator  $\hat{\mathcal{U}}_{I}(t)$  in the *interaction picture* and show that it satisfies an equation of the form

$$i\hbar \frac{d}{dt}\hat{\mathcal{U}}_I(t) = \hat{H}_I(t)\hat{\mathcal{U}}_I(t)$$
(16)

Assume  $\hat{W}$  is treated as the interaction. Show that  $\hat{H}_I(t)$  then is a linear function of  $\hat{a}$  and  $\hat{a}^{\dagger}$ ,

$$\dot{H}_{I}(t) = \theta(t)^{*} \,\hat{a} + \theta(t) \,\hat{a}^{\dagger} \tag{17}$$

and determine the function  $\theta(t)$ .

d) Show that the equation (16) has a solution of the form

$$\hat{\mathcal{U}}_{I}(t) = e^{\xi(t)\hat{a}^{\dagger} - \xi^{*}(t)\hat{a}} e^{i\phi(t)}$$
(18)

with  $\xi(t)$  as a complex and  $\phi(t)$  as a real function of time. What are the equations that these two functions should satisfy?

e) Use the expressions for  $\hat{\mathcal{U}}_0(t)$  and  $\hat{\mathcal{U}}_I(t)$  to find the time dependent state vector  $|\psi(t)\rangle$  in the Schrödinger picture, with the same initial condition as in a). Show that also in this case it describes a time dependent coherent state, of the form  $|\psi(t)\rangle = e^{i\gamma(t)}|z(t)\rangle$ . Find z(t) expressed in terms of  $z_0$ ,  $\xi(t)$  and  $\omega_0$ .

f) Determine the function  $\xi(t)$  and find an explicit expression for z(t). The corresponding real coordinate is  $x(t) = \sqrt{2\hbar/m\omega_0} \operatorname{Re} z(t)$ . Does this coordinate satisfy the classical equation of motion of the driven harmonic oscillator?