## FYS 4110 Non-relativistic Quantum Mechanics, Fall Semester 2015

## Problem set 2

### 2.1 Operator identities

Assume $\hat{A}$ and $\hat{B}$ to be two operators, generally not commuting.
We define the following to composite operators:

$$
\begin{align*}
\hat{F}(\lambda) & =e^{\lambda \hat{A}} \hat{B} e^{-\lambda \hat{A}} \\
\hat{G}(\lambda) & =e^{\lambda \hat{A}} e^{\lambda \hat{B}} \tag{1}
\end{align*}
$$

a) Show the following relation

$$
\begin{equation*}
\frac{d \hat{F}}{d \lambda}=[\hat{A}, \hat{F}] \tag{2}
\end{equation*}
$$

and use it to derive the expansion

$$
\begin{equation*}
\hat{F}(\lambda)=\hat{B}+\lambda[\hat{A}, \hat{B}]+\frac{\lambda^{2}}{2}[\hat{A},[\hat{A}, \hat{B}]] \ldots \tag{3}
\end{equation*}
$$

b) Show the following relation between $\hat{G}(\lambda)$ and $\hat{F}(\lambda)$,

$$
\begin{equation*}
\frac{d \hat{G}}{d \lambda}=(\hat{A}+\hat{F}) \hat{G} \tag{4}
\end{equation*}
$$

and use this to demonstrate the following expansion (Campbell-Baker-Hausdorff)

$$
\begin{equation*}
\hat{G}(\lambda)=e^{\lambda \hat{A}+\lambda \hat{B}+\frac{\lambda^{2}}{2}[\hat{A}, \hat{B}]+\ldots} \tag{5}
\end{equation*}
$$

by calculating the exponent on the right-hand side to second order in $\lambda$.
c) When $[\hat{A}, \hat{B}]$ commutes with both $\hat{A}$ and $\hat{B}$ the expression (5) is exact without the higher order terms indicated by ... in (5).
Verify this by use of (3) and (4), and by noting that the eigenvalues of $\hat{G}$ satisfy a differential equation that can be integrated.

### 2.2 Gaussian integral

The following formula gives the integral of a gaussian function

$$
\begin{equation*}
I \equiv \int_{-\infty}^{\infty} d x e^{-\lambda x^{2}}=\sqrt{\frac{\pi}{\lambda}} \tag{6}
\end{equation*}
$$

This is correct for complex $\lambda$ provided the real part of $\lambda$ is positive. Verify this by writing the $I^{2}$ as a two-dimensional integral

$$
\begin{equation*}
I^{2}=\int_{-\infty}^{\infty} \int_{-\infty}^{\infty} d x d y e^{-\lambda\left(x^{2}+y^{2}\right)} \tag{7}
\end{equation*}
$$

and by changing to to polar coordinates in the evaluation.

### 2.3 Gaussian wave function

The quantum state of a free particle with mass $m$ in one dimension is, at time $t=0$, described by the following momentum wave function of Gaussian form,

$$
\begin{equation*}
\psi(p)=\sqrt[4]{\frac{\lambda}{\pi}} e^{-\frac{\lambda}{2}\left(p-p_{0}\right)^{2}} \tag{8}
\end{equation*}
$$

where $\lambda$ is a real, positive parameter that determines the width of the Gaussian. What is the corresponding time dependent wave function $\psi(x, t)$ in the coordinate representation? Show that this function is also a Gaussian, of the form

$$
\begin{equation*}
\psi(x, t)=N^{\prime} e^{-\frac{\lambda^{\prime}}{2}\left(x-x_{0}\right)^{2}} \tag{9}
\end{equation*}
$$

with $\lambda^{\prime}$ and $x_{0}$ as a time-dependent parameters and $N^{\prime}$ as a normalization factor. Determine $\lambda^{\prime}$ and $x_{0}$ and examine how the maximum of the wave packet moves with time and how the width changes. Compare with the classical motion of the particle.

### 2.4 Time dependent transformations

Two unitarily equivalent descriptions of a quantum system are related by a time dependent unitary transformation $\hat{U}(t)$, which acts on state vectors as

$$
\begin{equation*}
|\psi(t)\rangle \quad \rightarrow \quad\left|\psi^{\prime}(t)\right\rangle=\hat{U}(t)|\psi(t)\rangle \tag{10}
\end{equation*}
$$

and on the observables as

$$
\begin{equation*}
\hat{A} \quad \rightarrow \quad \hat{A}^{\prime}(t)=\hat{U}(t) \hat{A} \hat{U}(t)^{-1} . \tag{11}
\end{equation*}
$$

Show that the Hamiltonian $\hat{H}^{\prime}$, which determines the Schrödinger equation of the transformed state vector $\left|\psi^{\prime}(t)\right\rangle$, includes an additional term which depends on the time derivative of $\hat{U}(t)$,

$$
\begin{equation*}
\hat{H} \rightarrow \hat{H}^{\prime}(t)=\hat{U}(t) \hat{H} \hat{U}(t)^{-1}+i \hbar \frac{d \hat{U}}{d t} \hat{U}^{-1} \tag{12}
\end{equation*}
$$

Discuss the meaning of the difference between the equations (11) and (12).
A unitary operator can generally be expressed as

$$
\begin{equation*}
\hat{U}=e^{i \hat{B}} \tag{13}
\end{equation*}
$$

where $\hat{B}$ is a hermitian operator. Assume $\hat{B}$, and therefore also $\hat{U}$, is time dependent. Assume also that the following commutator vanishes, $\left[\left[\hat{B}, \frac{d \hat{B}}{d t}\right], \hat{B}\right]=0$. Find the time derivative of $\hat{U}$ expressed in terms of $\hat{B}$ and its time derivative, and show that when introduced in Eq.(12), this equation gets the following form

$$
\begin{equation*}
\hat{H}^{\prime}=e^{i \hat{B}} \hat{H} e^{-i \hat{B}}-\hbar\left(\frac{d \hat{B}}{d t}+\frac{i}{2}\left[\hat{B}, \frac{d \hat{B}}{d t}\right]\right) \tag{14}
\end{equation*}
$$

Note in particular the contribution from the commutator between $\hat{B}$ and its time derivative.

