

# Problem set 8

## 8.1 Photon emission

A particle with mass  $m$  and charge  $e$  is trapped in a one-dimensional harmonic oscillator potential, with the motion restricted to the  $z$ -axis. The frequency of the oscillator is  $\omega$ . At time  $t = 0$  the particle is excited to energy level  $n$  and it then performs a transition to level  $n - 1$  by emitting one photon of energy  $\hbar\omega$ . We write the energy eigenstates of the composite system of charged particle and photons as  $|n, n_{\mathbf{k}a}\rangle$ . With initially no photon present the state is  $|i\rangle = |n, 0\rangle$ , while the final state with one photon present is  $|f\rangle = |n - 1, 1_{\mathbf{k}a}\rangle$ . To first order in perturbation theory the angular probability distribution  $p(\theta, \phi)$  of the emitted photon is

$$p(\theta, \phi) = \kappa \sum_a |\langle n - 1, 1_{\mathbf{k}a} | \hat{H}_{emis} | n, 0 \rangle|^2 \quad (1)$$

with  $(\theta, \phi)$  as the polar angle of the photon quantum number  $\mathbf{k}$  and  $\kappa$  as a proportionality factor.  $\hat{H}_{emis}$  is the emission part of the interaction Hamiltonian, which in the dipole approximation is

$$\hat{H}_{emis} = -\frac{e}{m} \sum_{\mathbf{k}a} \sqrt{\frac{\hbar}{2V\epsilon_0\omega}} \hat{\mathbf{p}} \cdot \boldsymbol{\epsilon}_{\mathbf{k}a} \hat{a}_{\mathbf{k}a}^\dagger \quad (2)$$

a) Show that for an arbitrary (real) vector  $\mathbf{a}$  we have the identity

$$\sum_a (\mathbf{a} \cdot \boldsymbol{\epsilon}_{\mathbf{k}a})^2 = \mathbf{a}^2 - \left(\mathbf{a} \cdot \frac{\mathbf{k}}{k}\right)^2 \quad (3)$$

b) Determine the particle matrix element  $\langle n - 1 | \hat{\mathbf{p}} | n \rangle$ .

c) Find the probability distribution  $p(\theta, \phi)$ .

The relation between the momentum operator and the ladder operators of the harmonic oscillator is found in Sect. 1.4.4 of the lecture notes.

## 8.2 Electric dipole transition in hydrogen (Exam 2008)

We consider the transition in hydrogen from the excited  $2p$  level to the ground state  $1s$ , where a single photon is emitted. The initial atomic state (A) we assume to have  $m = 0$  for the  $z$ -component of the orbital angular momentum, so that the quantum numbers of this state are  $(n, l, m) = (2, 1, 0)$ , with  $n$  as the principle quantum number and  $l$  as the orbital angular momentum quantum number. Similarly the ground state (B) has quantum numbers  $(n, l, m) = (1, 0, 0)$ . When expressed in polar coordinates the wave functions of the two states (with intrinsic spin of the electron not included) are given by

$$\begin{aligned} \psi_A(r, \phi, \theta) &= \frac{1}{\sqrt{32\pi a_0^3}} \cos \theta \frac{r}{a_0} e^{-\frac{r}{2a_0}} \\ \psi_B(r, \phi, \theta) &= \frac{1}{\sqrt{\pi a_0^3}} e^{-\frac{r}{a_0}} \end{aligned} \quad (4)$$

where  $a_0$  is the Bohr radius.

We remind you about the form of the interaction matrix element in the dipole approximation,

$$\langle B, 1_{\mathbf{k}a} | \hat{H}_{emis} | A, 0 \rangle = ie \sqrt{\frac{\hbar\omega}{2V\epsilon_0}} \boldsymbol{\epsilon}_{\mathbf{k}a} \cdot \mathbf{r}_{BA} \quad (5)$$

where  $e$  is the electron charge,  $\mathbf{k}$  is the wave vector of the photon,  $a$  is the polarization quantum number,  $\omega$  is the photon frequency and  $\boldsymbol{\epsilon}_{\mathbf{k}a}$  is a polarization vector.  $V$  is a normalization volume for the electromagnetic wave functions,  $\epsilon_0$  is the permittivity of vacuum and  $\mathbf{r}_{BA}$  is the matrix element of the electron position operator between the initial and final atomic states.

- Explain why the x- and y-components of  $\mathbf{r}_{BA}$  vanish while the z-component has the form  $z_{BA} = \nu a_0$ , with  $\nu$  as a numerical factor. Determine the value of  $\nu$ . (A useful integration formula is  $\int_0^\infty dx x^n e^{-x} = n!$ .)
- To first order in perturbation theory the interaction matrix element (5) determines the direction of the emitted photon, in the form of a probability distribution  $p(\phi, \theta)$ , where  $(\phi, \theta)$  are the polar angles of the wave vector  $\mathbf{k}$ . Determine  $p(\phi, \theta)$  from the above expressions.
- The life time of the 2p state is  $\tau_{2p} = 1.6 \cdot 10^{-9} s$  while the excited 2s state (with angular momentum  $l = 0$ ) has a much longer life time,  $\tau_{2s} = 0.12 s$ . Do you have a (qualitative) explanation for the large difference?

### 8.3 Spinflip radiation

We will study the transition between the two spin states of an electron in an external magnetic field directed along the  $z$ -axis,  $\mathbf{B} = b\mathbf{e}_z$ . The Hamiltonian can be expressed as  $H = H_0 + H_1$ , where  $H_0$  describes the energy of a magnetic dipole in the external field, while  $H_1$  describes the interaction with the radiation field. Then we have

$$H_0 = \frac{\hbar}{2} \omega_B \sigma_z, \quad \omega_B = -\frac{eB}{m}$$

where  $m$  is the electron mass and  $e$  the electron charge (which is negative so that  $\omega_B > 0$ ). The matrix element of the interaction part  $H_1$  for the case of emission of a single photon is in the dipole approximation given by

$$\langle B, 1_{\mathbf{k}a} | H_1 | A, 0 \rangle = i \frac{e\hbar}{2m} \sqrt{\frac{\hbar}{2\omega V \epsilon_0}} (\mathbf{k} \times \boldsymbol{\epsilon}_{\mathbf{k}a}) \cdot \boldsymbol{\sigma}_{BA}$$

where  $|A\rangle$  is the excited spin state (spin up) and  $|B\rangle$  is the ground state (spin down).  $\boldsymbol{\epsilon}_{\mathbf{k}a}$  is the polarization vector and  $\omega = ck$  is the angular frequency of the emitted photon.  $V$  is the normalization volume for the electromagnetic radiation and  $\boldsymbol{\sigma}_{BA} = \langle B | \boldsymbol{\sigma} | A \rangle$  is the matrix element of the vector  $\boldsymbol{\sigma}$  of the Pauli matrices.

- To first order in perturbation theory, the angular dependency of the squared matrix element  $|\langle B, 1_{\mathbf{k}a} | H_1 | A, 0 \rangle|^2$  will determine the probability distribution for the direction of the emitted photon,  $p(\theta, \phi)$ , where  $(\theta, \phi)$  are the polar coordinates for the wavevector  $\mathbf{k}$ . Determine  $p(\theta, \phi)$  using the above expression for the matrix element. It may be useful to know that when summing over the polarization states we have  $\sum_a |\boldsymbol{\epsilon}_{\mathbf{k}a} \cdot \mathbf{b}|^2 = |\mathbf{b}|^2 - |\mathbf{b} \cdot \frac{\mathbf{k}}{k}|^2$  for an arbitrary vector  $\mathbf{b}$ . The probability distribution should be normalized as  $\int d\phi \int d\theta \sin\theta p(\theta, \phi) = 1$ .

- b) The squared matrix element also determines, for a given  $\mathbf{k}$ , the probability distribution for the polarization direction of the photon. Assume that a photon detector is set to detect photons emitted in the  $x$ -direction and with polarization vector  $\epsilon(\alpha) = \cos \alpha \mathbf{e}_y + \sin \alpha \mathbf{e}_z$  (here  $\mathbf{e}_y$  and  $\mathbf{e}_z$  are unit vectors in the  $y$ - and  $z$ - directions). Find the probability distribution  $p(\alpha)$  to detect the emitted photon as a function of the angle  $\alpha$ .
- c) To a good approximation, the probability to find the spin in the excited state decays exponentially with time

$$P_A(t) = e^{-t/\tau}$$

where the lifetime  $\tau$  is, to first order in the interaction, determined by the time independent transition rate

$$w_{BA} = \frac{V}{(2\pi\hbar)^2} \int d^3k \sum_a |\langle B, 1_{\mathbf{k}a} | H_1 | A, 0 \rangle|^2 \delta(\omega - \omega_B)$$

Use this to find an expression for the lifetime  $\tau$ .