## FYS4160 - General Relativity Problem Set 0 Solutions Spring 2024

These solutions are credited to Jake Gordin, who wrote them in the years 2020-23.

If you spot any typos, mistakes, don't hesitate to contact me at halvor.melkild@fys.uio.no. For any physics related question please use the forum at astro-discourse.uio.no.

The idea of these solutions is to give you a sense of what a 'model' answer should be, and they also elaborate on some discussions from the help sessions. I try to make them "pedagogical": i.e. hopefully comprehensive and most steps should be explained.

## Problem 1. Index notation

a) • 
$$(\nabla S)_i = \partial_i S$$

- $\nabla \cdot \mathbf{A} = \partial_i A^i$ 
  - $(\nabla \times \mathbf{A})_i = \epsilon_{ijk} \partial^j A^k$
  - $\operatorname{Tr}(M) = M_{ij}\delta^{ij} = M^{i}_{\ i}$
  - $(M_{ij})^T = M_{ji}$

b) •  $\nabla \cdot (\nabla \times \mathbf{A}) = 0$ , where **A** is a 3D vector:

$$\partial_i (\epsilon^{ijk} \partial_j A_k) = \epsilon^{ijk} \partial_i \partial_j A_k = \epsilon^{ijk} \partial_{(i} \partial_{j)} A_k = \epsilon^{(ij)k} \partial_{(i} \partial_{j)} A_k = 0.$$

Step three to four follows from equation (1.83) in Carroll. The answer is then zero because of the antisymmetry of  $\epsilon^{ijk}$  – it contributes no symmetric part when summed over.

•  $\nabla \times (\nabla S) = 0$ , where S is a scalar:

 $\epsilon^{ijk}\partial_j(\partial_k S) = \epsilon^{ijk}\partial_{(j}\partial_{k)}S = \epsilon^{(ij)k}\partial_{(j}\partial_{k)}S = 0 \text{ (Same argument as above.)}$ 

## Problem 2. Index notation part 2, electric boogaloo

- a) •  $\partial_{\mu}x^{\nu} = \delta^{\nu}_{\mu}$ Right:  $\partial_{\mu}x^{\nu} = \frac{\partial x^{\nu}}{\partial x^{\mu}} = \delta^{\nu}_{\mu}$ •  $\partial_{\mu}x^{\mu} = 1$ 
  - Wrong:  $\partial_{\mu}x^{\mu} = \delta^{\mu}_{\mu} = 4$
  - $\partial^{\mu}x^{\nu} = g^{\mu\nu}$ Right:  $\partial^{\mu}x^{\nu} = g^{\mu\rho}\partial_{\rho}x^{\nu} = g^{\mu\rho}\delta^{\nu}_{\rho} = g^{\mu\nu}$
  - $T_{\alpha}{}^{\beta}{}_{\gamma} = g^{\beta\mu}T_{\alpha\mu\gamma} = g^{\mu\beta}T_{\alpha\mu\gamma}$ Right: the metric is sym-metric (I'll show myself out).
  - $T_{\alpha}{}^{\beta}{}_{\beta} = g_{\alpha\mu}g^{\beta\alpha}T^{\mu}{}_{\alpha\beta}$

Wrong: this is a case of nonsensical notation, there are three  $\alpha$ 's - the sum over them is meaningless. Replace with a different dummy index to get  $T_{\alpha}{}^{\beta}{}_{\beta} = g_{\alpha\mu}g^{\beta\nu}T^{\mu}{}_{\nu\beta} = g^{\beta\nu}T_{\alpha\nu\beta}$ 

b) • All independent Lorentz scalars from two four-vectors A and B:

 $A^{\mu}A_{\mu}, B^{\mu}B_{\mu}, A^{\mu}B_{\mu}$  (Lorentz scalar here simply means no free indices.)

- All independent Lorentz scalars from a rank-2 tensor T:  $T^{\mu\nu}T_{\mu\nu}, T^{\mu\nu}T_{\nu\mu}, T^{\mu}{}_{\mu}, T^{\mu}{}_{\nu}T^{\nu}{}_{\rho}T^{\rho}{}_{\mu}, T^{\mu}{}_{\nu}T^{\nu}{}_{\rho}T^{\rho}{}_{\lambda}T^{\lambda}{}_{\mu}, \dots$  (again, no free indices) The list can go on forever, but the list of <u>independent</u> scalars is still finite, as the number of components of T is finite.
- All independent Lorentz scalars from a rank-2 tensor T and two four-vectors A and B:  $T^{\mu\nu}A_{\mu}A_{\nu}, T^{\mu\nu}A_{\mu}B_{\nu}, T^{\mu\nu}B_{\mu}A_{\nu}, T^{\mu\nu}B_{\mu}B_{\nu}, T^{\mu}{}_{\mu}A^{\nu}B_{\nu}, T^{\mu}{}_{\nu}T^{\nu}{}_{\rho}A^{\rho}A_{\mu}, T^{\mu\nu}T_{\mu\nu}A^{\rho}A_{\rho}, \dots$  (no free indices again - double check though!)
- All independent Lorentz four-vectors from a scalar S and two four-vectors A and B:  $SA^{\mu}$ ,  $SB^{\mu}$  (Four-vector means one free index, so these are the only combinations.)