## FYS4160 - General Relativity Problem Set 7 Solutions Spring 2024

These solutions are credited to Jake Gordin, who wrote them in the years 2020-23.

If you spot any typos, mistakes, don't hesitate to contact me at halvor.melkild@fys.uio.no. For any physics related question please use the forum at astro-discourse.uio.no.

The idea of these solutions is to give you a sense of what a 'model' answer should be, and they also elaborate on some discussions from the help sessions. I try to make them "pedagogical": i.e. hopefully comprehensive and most steps should be explained.

## Problem 23. Curved d'Alembertians.

We start with the definition of  $\Box$  in an arbitrary curved spacetime:

$$\Box = g_{\mu\nu} \nabla_{\mu} \nabla_{\nu}.$$

Here's a cute little trick: the Christoffel symbol can be written as<sup>1</sup>

$$\Gamma^{\mu}_{\mu\nu} = \frac{1}{\sqrt{g}} \partial_{\nu}(\sqrt{g}).$$

The divergence of a vector field can therefore be written as

$$\nabla_{\mu}V^{\mu} = \frac{1}{\sqrt{g}}\partial_{\mu}(\sqrt{g} V^{\mu}).$$

Back to our question: note that f is a *scalar* test function, and therefore the quantity  $\nabla^{\mu} f$  is a vector field. Now, we need to put them together,

$$\Box f = g_{\mu\nu} \nabla_{\mu} \nabla_{\nu} f$$
$$\Box f = \nabla_{\mu} \nabla^{\mu} f$$
$$\Box f = \frac{1}{\sqrt{g}} \partial_{\mu} (\sqrt{g} \nabla^{\mu} f)$$

## Problem 24. Fluid equations.

- (a) Done step-by-step in Carroll, pg. 35-37, so I won't repeat the details. (P.S. I know I say this occasionally as a solution, but it isn't a valid answer in an exam lol).
- (b) Here we consider the derivation of the Navier-Stokes equations<sup>2</sup>. This can be done by adding a term  $\epsilon^{\mu\nu}$  to the energy-momentum tensor:

<sup>&</sup>lt;sup>1</sup>I derive this at the end of the solutions, for those who are interested.

<sup>&</sup>lt;sup>2</sup>This is entirely unrelated to physics, but I first heard of the Navier-Stokes equations in the book *Fashionable Nonsense* by Alan Sokal and Jean Bricmont. The Navier-Stokes equations come up because of a postmodern philosopher, Luce Irigaray, who claimed that they remain unsolved because they concern fluids and flow: these concepts are "feminine", in Irigaray's view, and the mostly male community of physicists are thus ill-equipped to understand such physics. Does the fact that the physics community is mostly male - now, in the 2020s, let alone the 90s - present significant challenges for women and impede the progress of physics as a whole? Absolutely. Does it mean that certain aspects of physics itself are gendered? I'll leave that as an exercise to the reader.

$$T^{\mu\nu} = (\rho + p)u^{\mu}u^{\nu} + p\eta^{\mu\nu} + \epsilon^{\mu\nu}$$

Let's think carefully about what the components of the energy-momentum tensor actually represent. Looking at our handy diagram below, we can see what terms actually correspond to shear stress: the off-diagonal, spatial components. Therefore  $\epsilon^{\mu\nu} \neq 0$  if  $\epsilon^{\mu\nu} = \epsilon^{jk}$  and  $j \neq k$ .

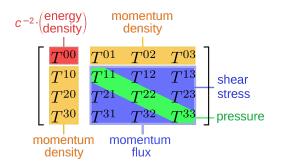


Figure 1: The breakdown of  $T^{\mu\nu}$ . (Taken once again from Wikipedia; if it ain't broke, don't fix it).

Recall that to derive the Euler equation, we considered the divergence of the energy-momentum tensor projected orthogonally to the fluid velocity.

$$P^{\sigma}{}_{\nu}\partial_{\mu}T^{\mu\nu} = 0.$$

The derivation follows the same procedure as for part (a), with the exception of an extra shear term:

$$\begin{aligned} (\delta^{\sigma}{}_{\nu} + u^{\sigma}u_{\nu})\partial_{\mu}\epsilon^{\mu\nu} &= \delta^{\sigma}{}_{\nu}\partial_{\mu}\epsilon^{\mu\nu} + u^{\sigma}u_{\nu}\partial_{\mu}\epsilon^{\mu\nu} \\ &= \partial_{\mu}\epsilon^{\mu\sigma} + u^{\sigma}u_{\nu}\partial_{\mu}\epsilon^{\mu\nu}. \end{aligned}$$

As part of the Euler derivation, the choice would be made to set  $\sigma = j$ , and so

$$(\delta^{\sigma}{}_{\nu} + u^{\sigma}u_{\nu})\partial_{\mu}\epsilon^{\mu\nu} = \partial_{\mu}\epsilon^{\mu j} + u^{j}u_{\nu}\partial_{\mu}\epsilon^{\mu\nu}.$$

We now take the non-relativistic limit (eq. 1.119 in Carroll), and sum:

$$\partial_{\mu}\epsilon^{\mu j} + u^{j}u_{\nu}\partial_{\mu}\epsilon^{\mu\nu} = \partial_{0}\epsilon^{0j} + \partial_{i}\epsilon^{ij} + v^{j}u_{0}\partial_{\mu}\epsilon^{\mu0} + v^{j}v_{i}\partial_{\mu}\epsilon^{\mu i}$$
$$= \partial_{i}\epsilon^{ij} + v^{j}v_{i}\partial_{k}\epsilon^{ki}$$
$$= \partial_{i}\epsilon^{ij} + \mathcal{O}(v^{2})$$

where any shear terms with a 0 component vanish because those components of  $\epsilon^{\mu\nu}$  are zero. In the non-relativistic limit, also, velocities are small, which is why we neglect  $v^2$  terms. If we define  $\epsilon = -\tau$ , the shear contribution to the now modified Euler equation is therefore

$$\rho(\partial_t \mathbf{v} + (\nabla \cdot \mathbf{v})) = -\nabla p + \nabla \cdot \tau$$

(in vector notation). The continuity equation does not get modified because this kind of shear, *viscous* stress, does not act parallel (or along) the fluid velocity.

## Appendix: derivation of the Christoffel in terms of the metric determinant.

For some generic matrix A, we have Jacobi's formula,:

$$\frac{d}{dt}\log\det(A(t)) = \operatorname{tr}\left(A^{-1}\frac{d}{dt}A(t)\right)$$

Differentiation of the LHS gives

$$\frac{1}{\det A(t)}\frac{d}{dt}(\det(A(t))) = \operatorname{tr}\left(A^{-1}\frac{d}{dt}A(t)\right)$$

Let's consider a 4x4 matric  $g_{\mu\nu}$  and use the shorthand for its determinant g; we differentiate

$$\partial_v g = (g) \operatorname{tr} \left( g^{-1} \partial_v g \right) = (g) g^{\mu \lambda} \partial_\nu g_{\mu \lambda}$$

Hence, the derivative of  $\sqrt{g}$  would then be

$$\partial_v \sqrt{g} = \frac{1}{2} g^{-1/2} \partial_\nu g$$
$$= \frac{1}{2} g^{-1/2} g g^{\mu\lambda} \partial_\nu g_{\mu\lambda}$$
$$= \frac{1}{2} \sqrt{g} g^{\mu\lambda} \partial_\nu g_{\mu\lambda}$$

Dividing by  $\sqrt{g}$  gives

$$\Gamma^{\mu}_{\mu\nu} = \frac{1}{2} g^{\mu\lambda} \partial_{\nu} g_{\mu\lambda},$$

the canonical Christoffel. (Compare to a contraction of equation 3.1 in Carroll, noting here that two indices are contracted).

Hence we can say that

$$\Gamma^{\mu}_{\mu\nu} = \frac{1}{\sqrt{\det g}} \partial_{\nu} (\sqrt{\det g}).$$