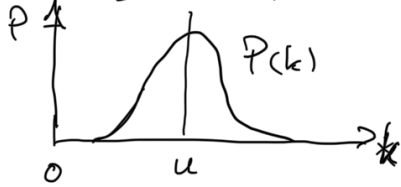


Diffusion and friction in fluids

Diffusion ~ random motion
 - independent of details

4.1.3 1D random walk



step lengths kL

$$u = \langle k_j \rangle = \sum_k k P_k \quad - \text{drift}$$

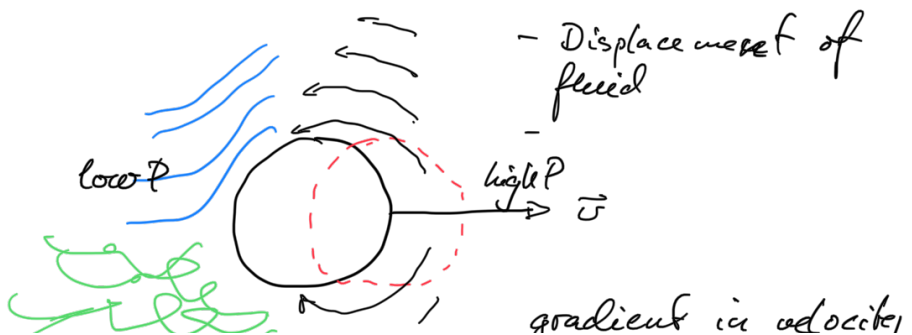
$$\langle x_N \rangle = NuL$$

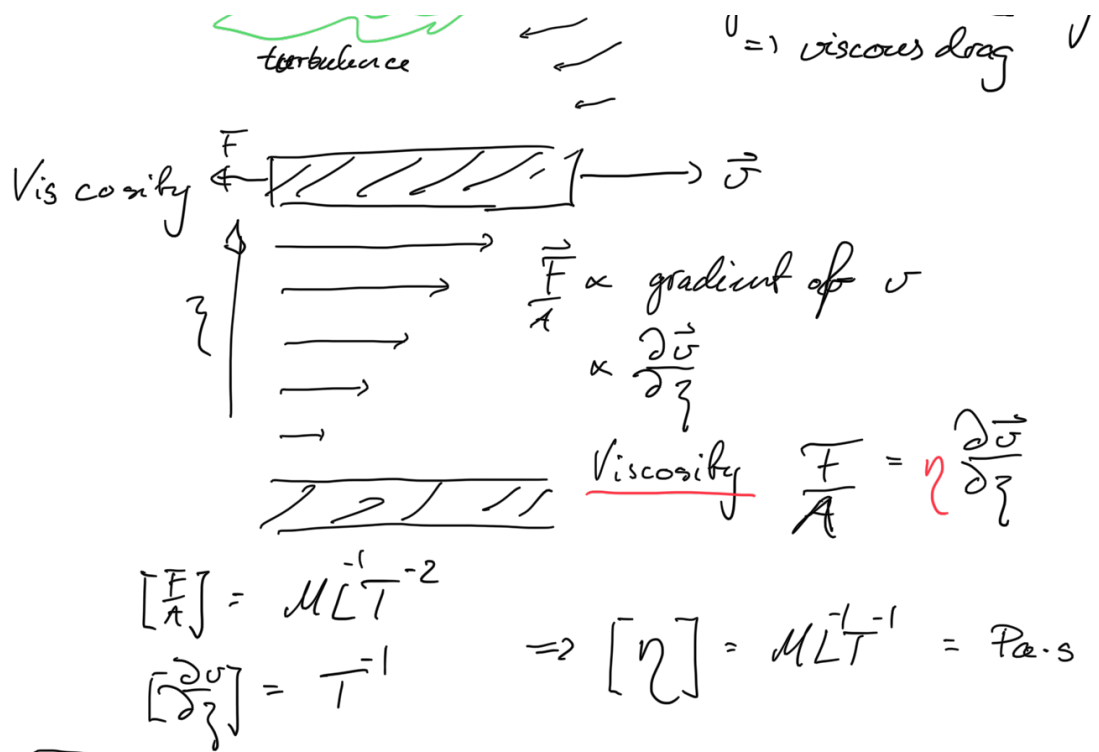
$$\sigma_N^2 = \langle (x_N - \langle x_N \rangle)^2 \rangle = 2Dt$$

This is general. Calculating σ_N^2 in detail is model dependent.

Dimensional analysis → drag

Length	L	m
time	T	s
mass	M	kg
area		
volume		
speed	LT^{-1}	
acceler.	LT^{-2}	
force	MLT^{-2}	N
work	ML^2T^{-2}	J
stress / pressure	$MT^{-2}L^{-1}$	Pa





Important quantities

- Velocity	\vec{v}	LT^{-1}
- Size	r	L
- Density	ρ	ML^{-3}
- Viscosity	η	MLT^{-1}
- Drag	F_d	MLT^{-2}

Buckingham's Π -theorem: The number of independent, dimensionless parameters necessary to describe the system is no. of physical quantities - no. of indep. dimensions

$$5 - 3 = 2$$

Dimensionless param: $Re = \frac{\rho v r}{\eta}$ $\frac{\rho v - \text{inertial}}{\eta - \text{viscous}}$

$Re \gg 1$ turbulence
 $Re \ll 1$ only viscous damping

F_d not used - make dim. less

$[F_d] = MLT^{-2}$ in terms of η and ρ

T T F, η MLT^{-2} L^{-2} L^2 L^{-2}

$$\left[\frac{a}{g} \right] = \frac{ML^{-3}}{ML^{-3}} = L^1 \quad \text{---} \quad \frac{1}{r^2} \quad L^2$$

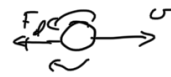
$$\Rightarrow C_R = \frac{F_d}{\rho v^2 r^2}$$

$$\text{II} \quad \frac{F_d}{\eta} = \frac{MLT^{-2}}{ML^{-1}T^{-1}} = L^2 T^{-1} \quad \begin{matrix} / v \\ / r \end{matrix}$$

$$C_s = \frac{F_d}{\eta v r}$$

Stokes drag $F_d = C_s \eta v r$ Sphere $C_s = 6\pi$
allow Re

Viscous friction $f = \frac{F_d}{v}$

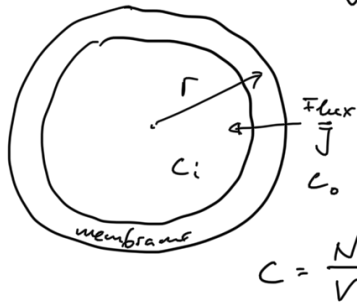


$$f = 6\pi \eta r$$

Exercise: Your book 4C, p120 \Rightarrow

$$fD = k_B T \quad \text{Einstein relation}$$

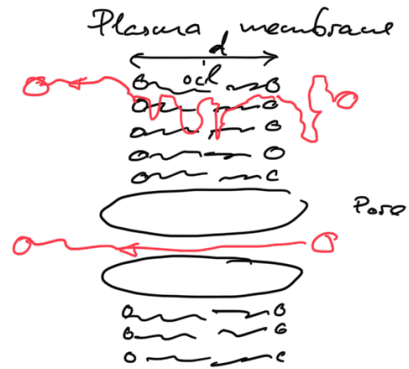
Permeation through a membrane



$$C = \frac{N}{V}$$

Conc. difference $\Delta C = c_o - c_i$

Flux $J = \frac{1}{A} \frac{\Delta N}{\Delta t}$ flow rate per unit area



How does c_i change with time? (c_o - const)

$$\frac{\partial c_i}{\partial t} = - \frac{\partial \Delta C}{\partial t} = \frac{1}{V} \frac{\partial N}{\partial t} = \frac{3}{r} J$$

$$A = 4\pi r^2$$

$$V = \frac{4}{3}\pi r^3$$

Permeability P : $J = P \Delta C$

$$\frac{V}{A} = \frac{r}{3}$$

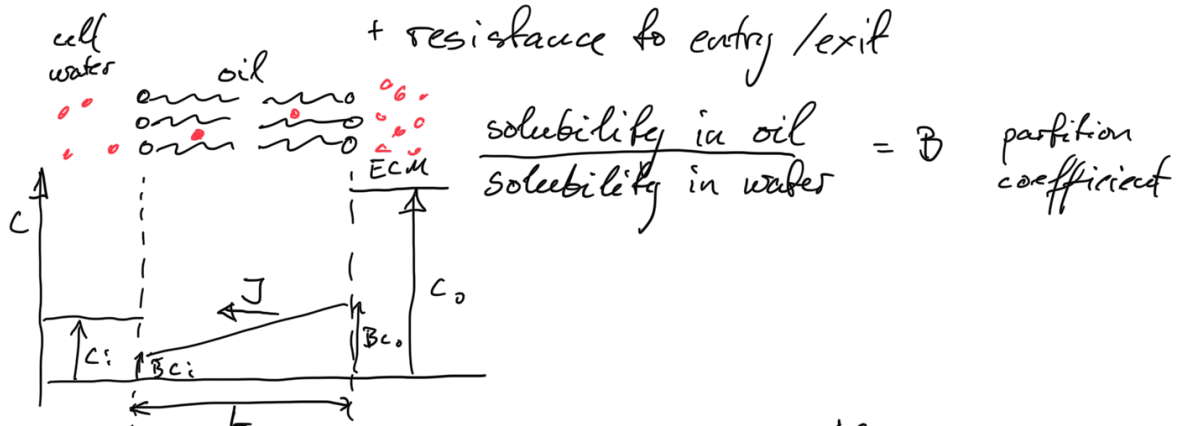
$$\Rightarrow - \frac{\partial \Delta C}{\partial t} = \frac{3}{r} P \Delta C$$

$$\int_{AC_0}^{AC} \frac{\partial AC}{AC} = -\frac{3P}{r} \int_0^t dt$$

$$\ln AC - \ln AC_0 = -\frac{3P}{r} t$$

$$\frac{AC}{AC_0} = e^{-t/\tau}, \quad \tau = \frac{r}{3P}$$

Permeability is + diffusion through membrane



Diffusion in membrane

$$J = -D_m B \frac{\Delta C}{L} = P \Delta C$$

$$\Rightarrow |P| = \frac{D_m B}{L}$$

B is small \Rightarrow passive permeation is small
 \Rightarrow "active" permeation biologically more important

For next lecture: Life at low Reynolds number

- p 3-4 Intro, swimming
- p 5-6 Scallop + other swimmers
- p 7-8 Corkscrew \rightarrow Datum
- p 9-10 Swimming & diffusion

lvat \leq simple implementation of Visbeck model?

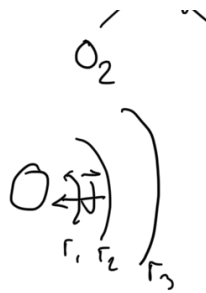
Dag: Interest of swimming of *E. coli*, active matter

4.6.2 p 138

Bacterial metabolism: O_2 limit



O_2 - consumed by bacteria



- transport from outside
- x Ocean: turbulence
- x ECM, mud, pond: diffusion

Assume instant uptake at bacteria surface

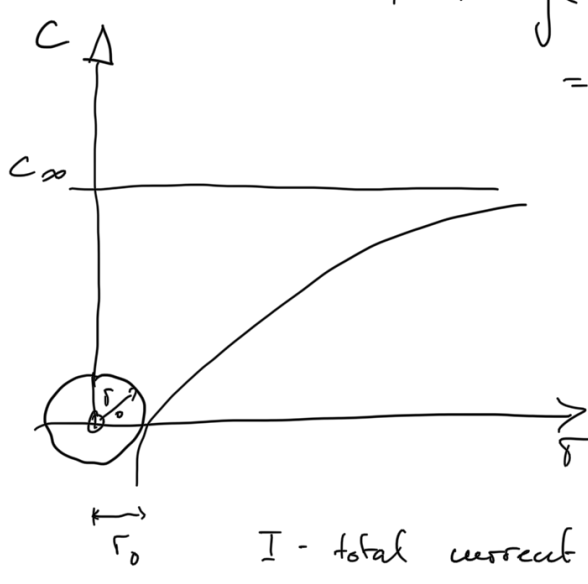
$$\text{Flux } j(r) = \frac{I}{4\pi r^2} = -D \frac{\partial c}{\partial r}$$

$$\Rightarrow \frac{-I}{4\pi D} \frac{\partial r}{r^2} = \partial c$$

$$\frac{-I}{4\pi D} \frac{1}{r} \Big|_{r_0}^{\infty} = c \Big|_0^{c_{\infty}}$$

$$\frac{-I}{4\pi D} \left[-\frac{1}{r_0} \right] = c_{\infty}$$

$$\Rightarrow \underline{I = 4\pi D c_{\infty} r_0}$$



I - total current of O₂

Ions at membranes - a first look

Electrolyte in electric field



Electric force on ions: $F = qE$

Drift velocity $v = \frac{F}{\zeta}$

$$\Rightarrow \text{Ion flux } j = c \cdot v = \frac{c q E}{\zeta}$$

Charge buildup at walls



Stationary $j = \frac{c q E}{\zeta} - D \frac{\partial c}{\partial x} = 0 \quad \left| \quad j = D \left(-\frac{\partial c}{\partial x} + \frac{qE}{k_B T} \right) \right.$

Einstein relation: $\zeta D = k_B T$ Nernst-Planck

$$\Rightarrow c q E - k_B T \frac{\partial c}{\partial x} = 0$$

$$\frac{q E}{k_B T} \partial x = \frac{\partial c}{c}$$

"B"

$$\frac{qE}{k_B T} L = \ln \frac{C_{right}}{C_{left}}$$

