Problem sheet 5

FYS5120-Advanced Quantum Field Theory

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Problem 5.2: Renormalized perturbation theory

In this exercise you will renormalize QED to one-loop order. To this end, start with the bare QED Lagrangian 1

$$\mathcal{L} = -\frac{1}{4} F_{0\,\mu\nu} F_0^{\mu\nu} + \bar{\psi}_0 \Big(i \partial \!\!\!/ - e_0 A_0 - m_0 \Big) \psi_0,.$$
(1)

By realizing that the *bare* Lagrangian does not describe the physical parameters, we can rescale the fields and coupling in the following way

$$\psi_0 = \sqrt{\mathcal{Z}_2} \, \psi \tag{2}$$

$$A_{0\,\mu} = \sqrt{\mathcal{Z}_3} \, A_\mu \tag{3}$$

$$m_0 = \mathcal{Z}_m m \tag{4}$$

$$e_0 = \mathcal{Z}_e e \,. \tag{5}$$

The idea from here is that we want to expand around some tree level values for these parameters, which naturally takes the from

$$\mathcal{Z}_i = 1 + \delta_i \,, \tag{6}$$

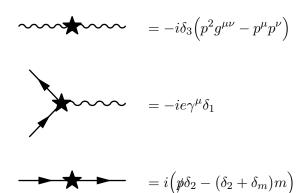
where δ_i starts at $\mathcal{O}(e^2)$. First, define

$$\mathcal{Z}_1 = \mathcal{Z}_e \mathcal{Z}_2 \sqrt{\mathcal{Z}_3} \,, \tag{7}$$

and show that

$$\delta_e = \delta_1 - \delta_2 - \frac{1}{2}\delta_3 + \mathcal{O}(e^4).$$
(8)

Finally, show that the Feynman rules for the counterterms can be written as



¹We ignore the gauge fixing term here.

Problem 5.2: Electron self energy

By using the Feynman rules for counterterms, show that the renormalized two point function can be written as (momentum space)

$$i\mathcal{G}(p) = \frac{i}{\not p - m + \Sigma(\not p)} \tag{9}$$

where

$$\Sigma(p) = \Sigma_2(p) - (\delta_m + \delta_2)m + \mathcal{O}(e^4).$$
⁽¹⁰⁾

Calculate the one-loop electron self energy, i.e. $\Sigma_2(p)$, and show that²

$$\delta_2 = -\frac{\alpha}{4\pi} \frac{1}{\bar{\epsilon}} \tag{11}$$

$$\delta_m = -\frac{3\alpha}{4\pi} \frac{1}{\bar{\epsilon}} \,, \tag{12}$$

in the \overline{MS} scheme. With these counterterms, you should find the UV-finite expression

$$\Sigma(p) = \frac{\alpha}{4\pi} \left(2m - p + 2\int_0^1 dx \left[(1-x)p - 2m \right] \ln \frac{\mu^2}{\Delta} \right) + \mathcal{O}(e^4)$$
(13)

where

$$\Delta = xm^2 - x(1-x)p^2.$$
(14)

Now, there is an IR-divergence in this expression, but we will not worry about these now. The easy way to cure this divergence is to add a photon mass to the photon propagator, but these IR-divergences will be treated more correctly in chapter 20.

Problem 5.3: Vertex correction

In Problem sheet 4 you found that the vertex correction could be written as

$$\Gamma_{1-\text{loop}}^{\mu} = F_1^{1-\text{loop}}(q^2)\gamma^{\mu} + \frac{i}{m}q_{\nu}S^{\mu\nu}F_2^{1-\text{loop}}(q^2)$$
(15)

where the one-loop form factors took the form

$$F_1^{1-\text{loop}} = -2ie^2\mu^{4-d} \int_0^1 dx dy dz \delta(x+y+z-1) \int \frac{d^d k}{(2\pi)^d} \frac{\mathcal{N}_1}{(k^2 - \Delta + i\epsilon)^3}$$
(16)

$$F_2^{1-\text{loop}} = 2ie^2 \mu^{4-d} \int_0^1 dx dy dz \delta(x+y+z-1) \int \frac{d^d k}{(2\pi)^d} \frac{\mathcal{N}_2}{(k^2 - \Delta + i\epsilon)^3}$$
(17)

where

$$\Delta = (1 - z)^2 m^2 - xyq^2$$
(18)

and

$$\mathcal{N}_1 = \frac{(d-2)^2}{d}k^2 - (d-2)\left((1-z)^2m^2 + xyq^2\right) + 2z(2m^2 - q^2) \tag{19}$$

$$\mathcal{N}_2 = 2m^2 (2z(1-z) + (4-d)(1-z)^2)$$
⁽²⁰⁾

²Remember that $\bar{\epsilon}$ includes $\ln 4\pi$ and γ_E .

By calculating $F_2^{1-\text{loop}}(0)$ you found the correction to the magnetic moment of the electron. In this exercise you will calculate $F_1^{1-\text{loop}}(0)$ and find the counterterm δ_1 . First, to leading order we have that

$$F_1(p^2) = 1. (21)$$

Hence, it is natural to define the renormalization condition (why?)

$$\Gamma^{\mu}(0) = \gamma^{\mu} \,. \tag{22}$$

Use renormalized perturbation theory and show that

$$\delta_1 = -\frac{\alpha}{4\pi} \frac{1}{\bar{\epsilon}},\tag{23}$$

in the \overline{MS} scheme. The last counterterm we need to remove all divergences to one-loop in QED, is δ_3 . From the renormalized Lagrangian, you should conclude that δ_3 cancels the divergence in the photon propagator. The one-loop correction to the photon propagator was calculated in *Problem sheet 3*, and you should find that

$$\delta_3 = -\frac{\alpha}{3\pi} \frac{1}{\bar{\epsilon}} \,. \tag{24}$$