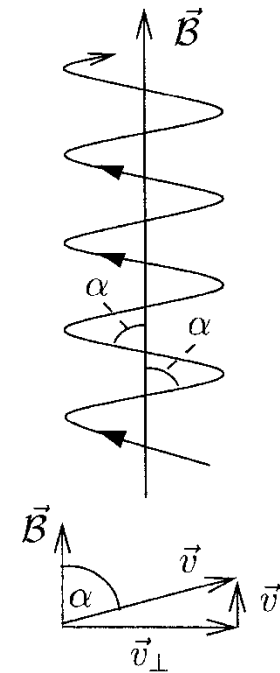
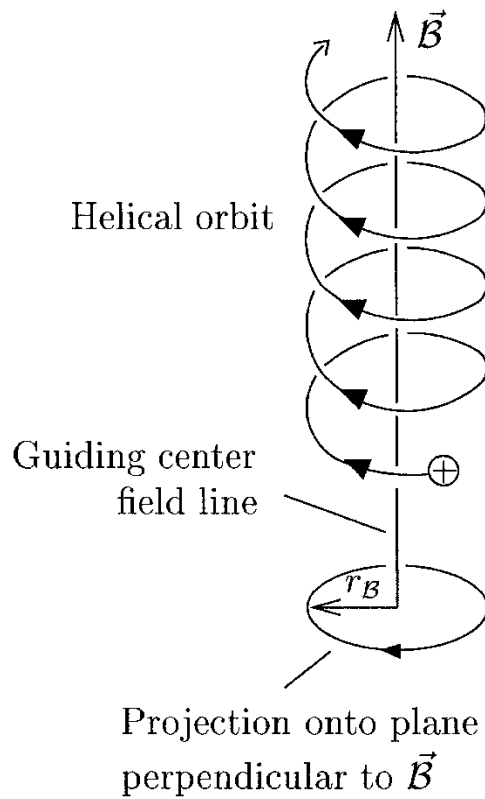


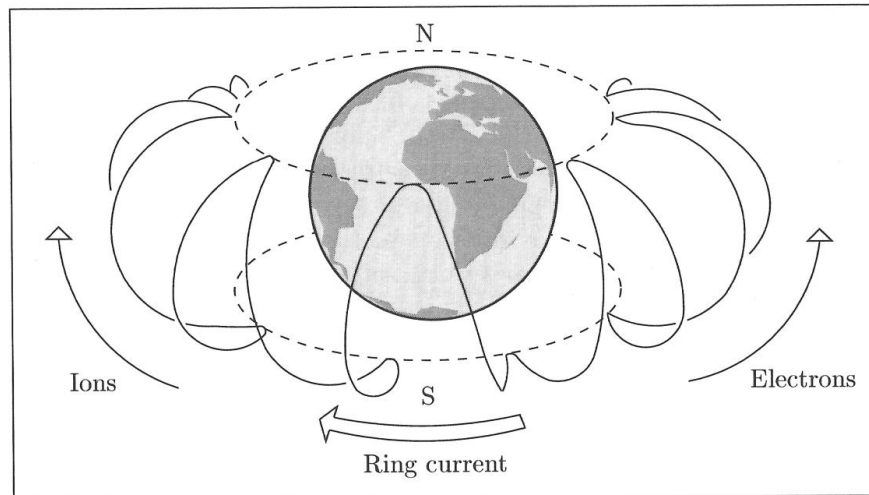
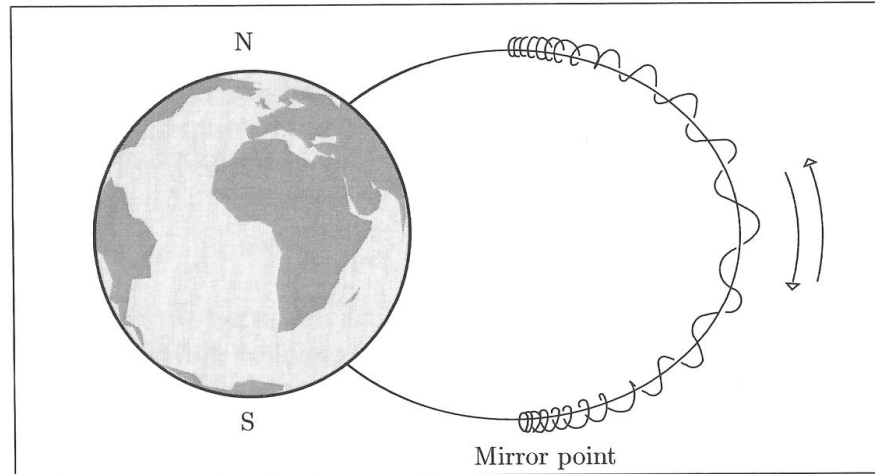
# Gyration

$$v_{\parallel} \neq 0$$

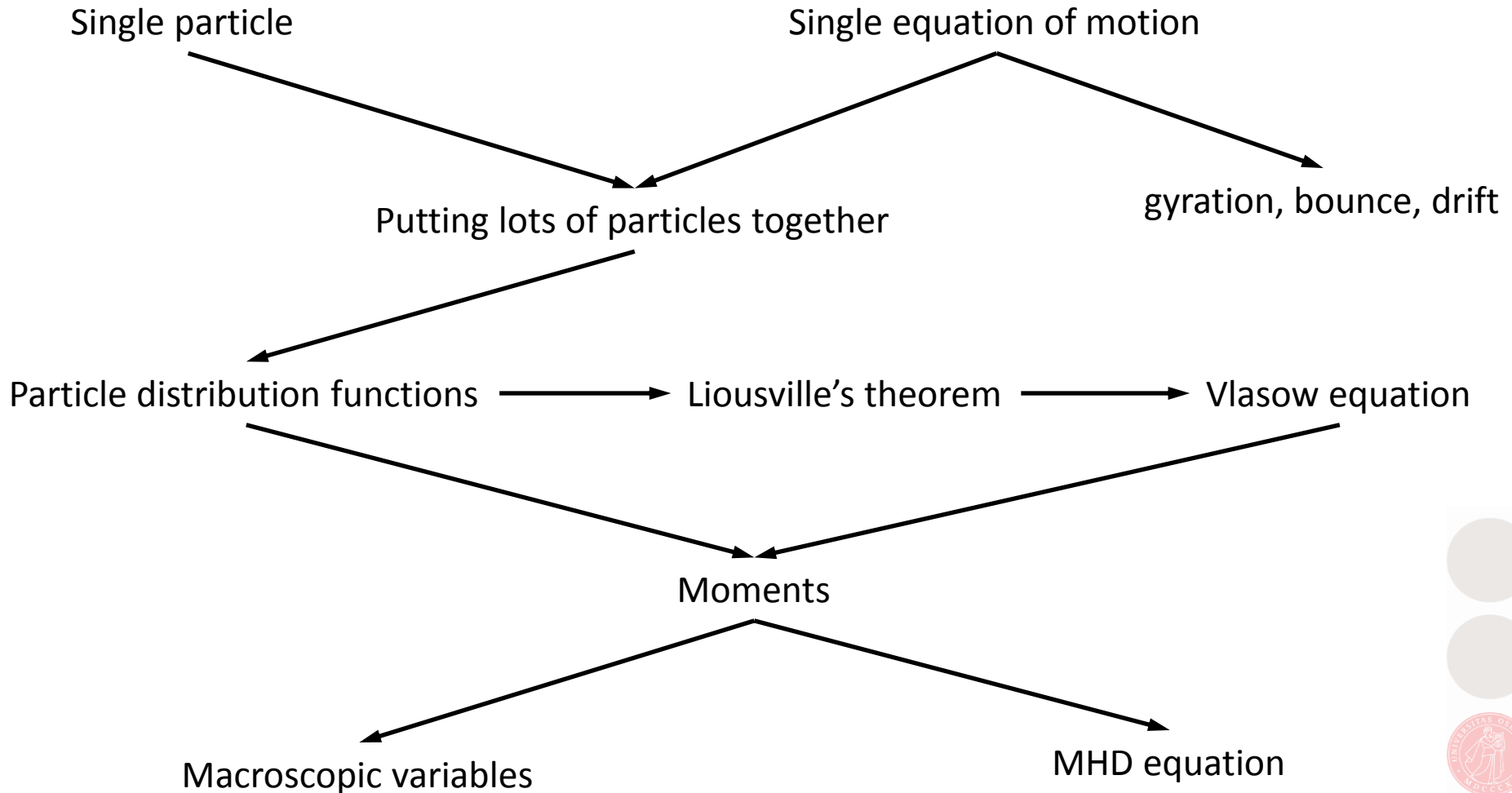
Projection onto plane  
parallel to  $\vec{B}$



# Bounce and drift motion



# SPM to MHD



# The road to MHD

	State	Fields	Dynamics
Single particle motion	$q, m, \vec{r}, \vec{v}$	they just are	$m \frac{d\vec{v}}{dt} = q(\vec{E} + \vec{v} \times \vec{B})$
kinetic theory	$q_s, m_s, f_s(\vec{v}, \vec{r}, t)$	$\rho = \sum_s q_s \int f_s(\vec{v}, \vec{r}, t) d^3v$ $\vec{j} = \sum_s q_s \int \vec{v} f_s(\vec{v}, \vec{r}, t) d^3v$ $\nabla \cdot \vec{E} = \frac{\rho}{\epsilon_0} \quad \nabla \cdot \vec{B} = 0$ $\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t} \quad \nabla \times \vec{B} = \mu_0 \vec{j} + \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t}$	$\frac{df_s}{dt} = 0$
Magnetohydrodynamics (MHD)	$m, n, \vec{v}, p$	$\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t} \quad \nabla \times \vec{B} = \mu_0 \vec{j}$ $\vec{j} = \sigma(\vec{E} + \vec{v} \times \vec{B}) \quad p\rho^{-\gamma} = \text{const.}$	$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \vec{v}) = 0$ $\rho \frac{d\vec{v}}{dt} = -\nabla p + \vec{j} \times \vec{B}$



# MHD equations

$$\frac{\partial \rho}{\partial t} + \nabla(\rho \vec{u}) = 0$$

$$\rho \frac{D\vec{u}}{Dt} = -\nabla p + \rho \vec{g} + \vec{j} \times \vec{B}$$

$$p = \alpha \rho^{\gamma^*}$$

$$\nabla \times \vec{B} = \mu_0 \vec{j}$$

$$\frac{\partial \vec{B}}{\partial t} = \nabla \times (\vec{u} \times \vec{B})$$



# Diffusion vs. frozen-in

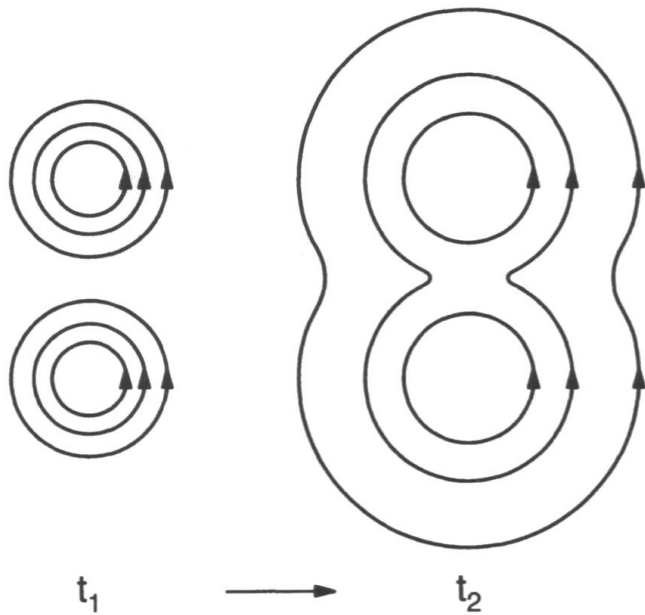
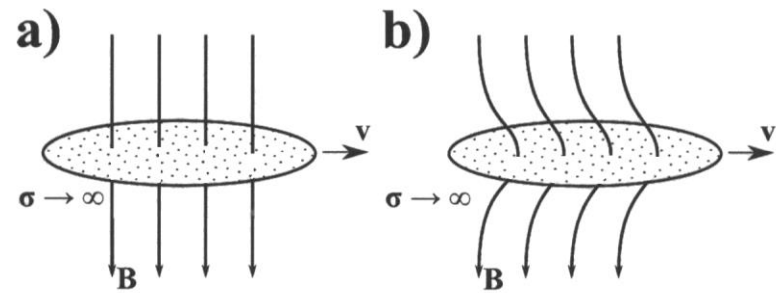


Fig. 5.1. Diffusion of magnetic field lines.



**Figure 1.2:** Illustration of the "frozen-in" theorem. **a)** A magnetic field penetrates a highly conducting plasma. **b)** As the plasma moves, the magnetic field is "frozen-in" and follows the motion of the plasma.



# Magnetic reconnection

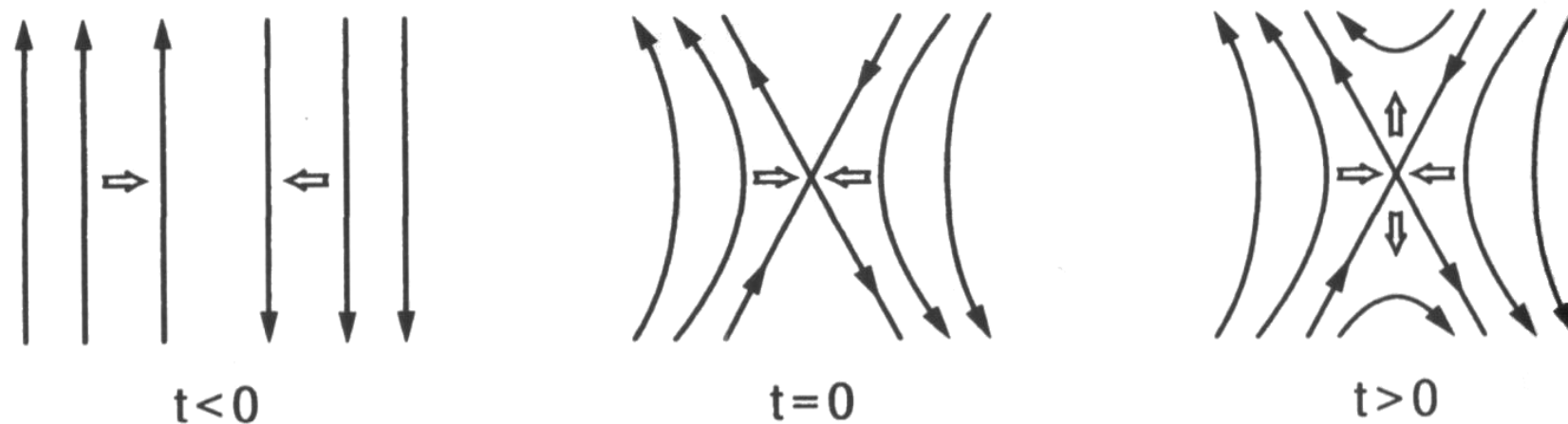


Fig. 5.3. Evolution of field line merging.

