

**GEF4610 – OBLIGATORY EXERCISE 1**  
**All symbols are defined in the GEF4610 Lecture Notes**

**Problem 1**

The equations governing the Ekman currents are:

$$\begin{aligned} 0 &= f v + A^{(z)} \frac{d^2 u}{dz^2}, \\ 0 &= -f u + A^{(z)} \frac{d^2 v}{dz^2}. \end{aligned} \tag{1.1}$$

a) State the assumptions behind this set of equations.

b) We apply a constant wind stress  $\tau$  along the y-axis at the surface  $z = 0$ , and a no-slip bottom ( $u = v = 0$ ) at  $z = -H$ . Solve (1.1) and show that the complex solution  $W = u + iv$  may be written

$$W(z) = \frac{i\tau}{\rho_r a A^{(z)}} \frac{\sinh((z+H)a)}{\cosh(Ha)}, \tag{1.2}$$

where  $a^2 = if / A^{(z)}$ .

c) We now consider the angle  $\alpha$  between the wind direction and the surface velocity. Express (1.2) in terms of the Ekman depth  $D_E = \pi(2A^{(z)} / f)^{1/2}$ , and apply a numerical software, e.g. matlab, to plot  $\alpha$  for various values of the ratio  $D_E / H$ . Discuss your findings.

**Problem 2**

a) State the assumptions leading to the equations

$$\begin{aligned} u_t + uu_x + vv_y - fv &= -g\eta_x - P_{Sx} / \rho, \\ v_t + uv_x + vv_y + fu &= -g\eta_y - P_{Sy} / \rho, \\ \eta_t + ((H + \eta)u)_x + ((H + \eta)v)_y &= 0, \end{aligned} \tag{2.1}$$

and derive the conservation equation for the potential vorticity:

$$\frac{DQ}{dt} \equiv \frac{D}{dt} \frac{f + \zeta}{H + \eta} = 0. \tag{2.2}$$

This means that  $Q$  is constant for a given fluid column. We will use these equations on an  $f$ -plane, and consider steady flow in this problem.

b) We have constant geostrophic flow  $U_0$  in the  $x$ -direction in a layer of constant depth  $H_0$ . What are the surface pressure gradients in this case?

c) We introduce a small sub-sea ridge  $q(x)$  of width  $2L$  which is infinitely long in the  $y$ -direction. The fluid depth then becomes  $H = H_0 - q$ ; see Fig. 1.

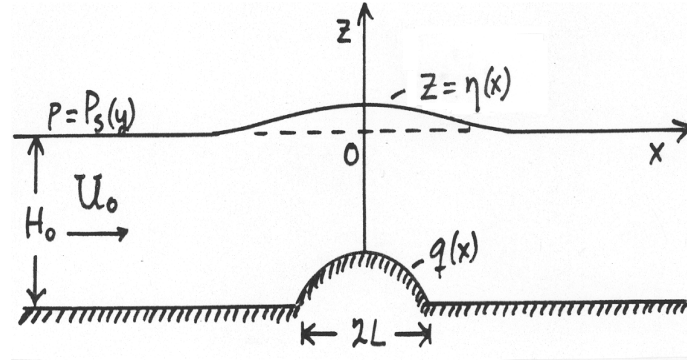


Figure 1. Figure sketch.

The flow will now be modified. The velocities in the  $x$ - and  $y$ -directions can be written  $U_0 + u(x)$  and  $v(x)$ , respectively, where  $u$  and  $v$  are small compared to  $U_0$ . Assume that the velocity far upstream of the ridge is  $U_0$ , the depth is  $H_0$ , and the relative vorticity is zero. Show from (2.2) that

$$v_x = \frac{f}{H_0}(\eta - q). \quad (2.3)$$

d) Show that the linearized momentum equations in (2.1) becomes

$$\begin{aligned} U_0 u_x - f v &= -g \eta_x, \\ U_0 v_x + f u &= 0. \end{aligned} \quad (2.4)$$

e) Combine (2.3) and (2.4), and show that

$$(1 - F_0^2) \eta_{xx} - \frac{1}{a^2} \eta = -F_0^2 q_{xx} - \frac{1}{a^2} q, \quad (2.5)$$

where

$$F_0 = \frac{U_0}{(gH_0)^{1/2}}, \quad a = \frac{(gH_0)^{1/2}}{f}. \quad (2.6)$$

$F_0$  and  $a$  are important fluid parameters. What is their significance?

f) Assume subcritical flow, and take that  $\eta \rightarrow 0$ ,  $x \rightarrow \pm\infty$ . Use that  $d/dx^2$  scales as  $1/L^2$  in order to simplify (2.5). Discuss the surface elevation over the ridge in the extreme cases  $L^2 \ll a^2$ , and  $L^2 \gg a^2$ .

g) Sometimes the rigid lid approximation is used for flow over a ridge (the surface is horizontal everywhere). Now (2.3) reduces to

$$v_x = -\frac{f}{H_0} q, \quad (2.7)$$

i.e. anti-cyclonic relative vorticity is generated over the entire ridge. Show from (2.7) that  $v$  downstream of the ridge is constant, and dependent on the cross-sectional area  $A$  of the ridge, given by

$$A = \int_{-L}^L q(x) dx. \quad (2.8)$$

h) Use the continuity equation in (2.1), and find  $u$  in this case.

i) Show that the angle between the current upstream of the ridge and the current downstream of the ridge in the case of a rigid lid is given by

$$\theta = \arctan\left(\frac{fA}{U_0 H_0}\right). \quad (2.9)$$