

INF-GEO 4310 Imaging

Lecture 16.09.2010

Fritz Albregtsen
Geometrical Optics part II

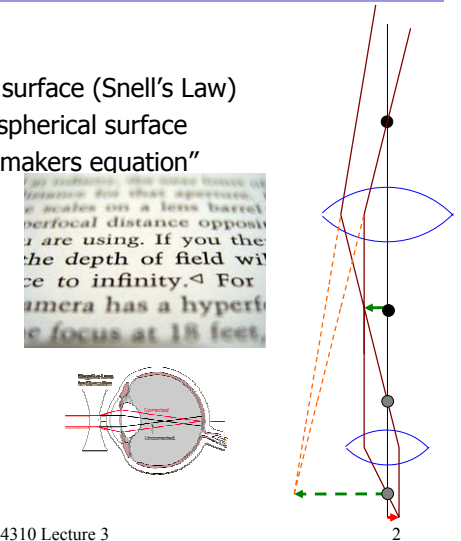


- Themes today:
 - Imaging by Refraction
 - Geometrical Optics: Diffraction
 - Geometrical Optics: Scattering

- Literature:
 - F. Albregtsen: "2. Reflection, refraction, diffraction, and scattering" (pages 37 - 82)

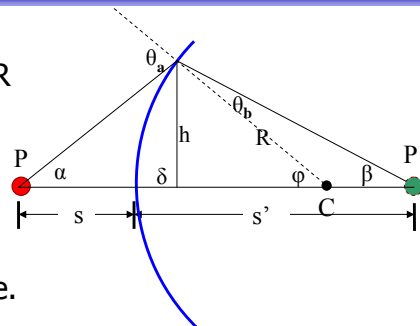
Geometrical Optics: Refraction

- Imaging by refraction
 - Refraction at a planar surface (Snell's Law)
 - Refraction at a single spherical surface
 - Thin lenses; "The lensmakers equation"
 - The camera;
 - "Depth-of-field"
 - The eye
 - The magnifier
 - The eyepiece
 - Microscopes
 - Telescopes
 - Multiple lens systems



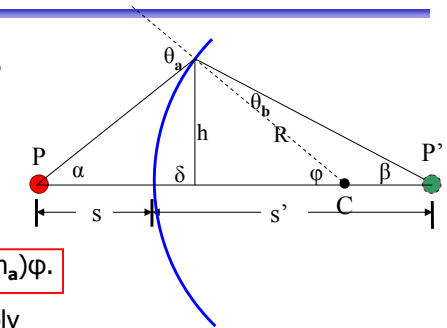
Refraction by spherical surface - I

- Convex surface with radius of curvature R facing incident light originating from P on the optical axis.
- A ray from P at an angle α to the axis is refracted at the surface.
- Angle of refraction is given by Snell's law, refracted ray crosses the optical axis at an angle β .
- All rays from P will intersect axis at the same point P' , provided that the angle α is small.



Refraction by spherical surface - II

- Snell's law : $n_a \sin \theta_a = n_b \sin \theta_b$.
- Paraxial approximation
 $\Rightarrow n_a \theta_a = n_b \theta_b$.
- Combining this with $\theta_a = \alpha + \phi$ gives $\theta_b = (\alpha + \phi) n_a / n_b$.
- Substituting this into $\phi = \beta + \theta_b$ we get $(n_a \alpha + n_b \beta) = (n_b - n_a) \phi$.
- Tangents of α , β , and ϕ are simply
 $\text{tg}(\alpha) = h/(s+\delta)$, $\text{tg}(\beta) = h/(s'-\delta)$, $\text{tg}(\phi) = h/(R-\delta)$.
- If the angle α is small, so are β and ϕ .
- Under the paraxial approximation, δ may be neglected compared to s , s' , and R .
- $\Rightarrow \alpha = h/s$, $\beta = h/s'$, $\phi = h/R$.



Refraction by spherical surface - III

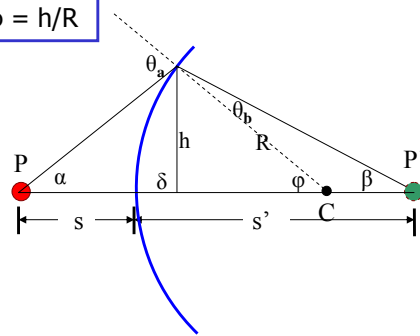
- Substituting $\alpha = h/s, \beta = h/s', \varphi = h/R$

into

$$(n_a \alpha + n_b \beta) = (n_b - n_a) \varphi$$

we get general **object-image relation**, for single spherical surface:

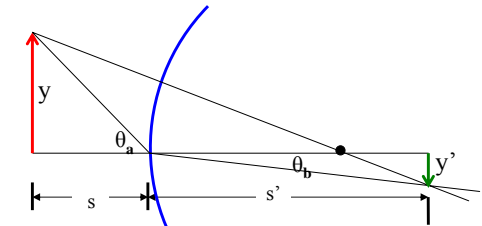
$$\frac{n_a}{s} + \frac{n_b}{s'} = \frac{n_b - n_a}{R}$$



- Same expression, no matter if n_a is greater or less than n_b .
- Does not contain the angle α :
=> all light rays coming from P will intersect at P'.

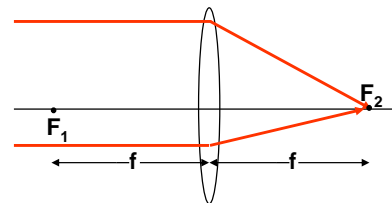
Imaging by spherical surface

- Given
(y, s, R)
- What is
 - the size y'
 - position s'
 - of the real image ?
- We see that
 $\text{tg}(\theta_a) = y/s$ and $\text{tg}(\theta_b) = -y'/s'$.
- Snell's law
 $n_a \sin(\theta_a) = n_b \sin(\theta_b)$
- Small angles => $n_a y/s = -n_b y'/s'$.
- Magnification: $m = y'/y = -(n_a s')/(n_b s)$.
- (negative sign => inverted image)

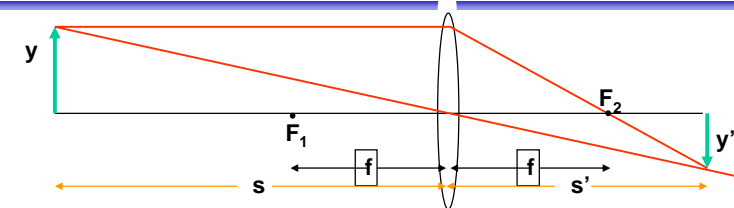


Thin lenses

- Thin lens: two refracting surfaces close enough to neglect distance between them.
- Centers of curvature of spherical surfaces lie on and define optical axis.
- The first and second focal points are on either side of the lens.
- Focal length: distance from focal point to middle of lens.
 - Focal length of convex lens is positive
 - Focal length of concave lens is negative.



Object-image relation

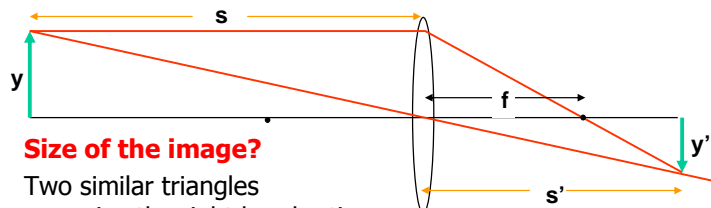


- We have two pairs of similar triangles, giving:

$$\frac{y'}{y} = \frac{s'}{s} \quad \text{and} \quad \frac{y'}{y} = \frac{s'-f}{f} \quad \text{implying} \quad \frac{1}{s} + \frac{1}{s'} = \frac{1}{f}$$

- This is known as the "*object-image relation*".

Image size and magnification



- **Size of the image?**
- Two similar triangles give the right hand ratios.
- Then use "object-image relation".

- The **magnification** is

$$m = y'/y$$

- If object is far away:

$$s \gg f \Rightarrow m \ll 1$$

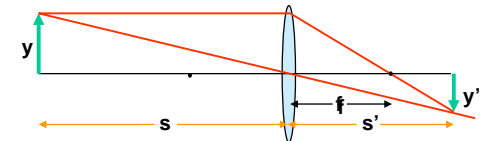
$$\frac{y'}{y} = \frac{s'}{s} \Rightarrow y' = \frac{y s'}{s}$$

$$\frac{1}{s} + \frac{1}{s'} = \frac{1}{f} \Rightarrow s' = \frac{s f}{(s - f)}$$

$$\Rightarrow y' = \frac{y f}{(s - f)}$$

Size of image of the Moon ?

- What is the size the Moon imaged by lens having $f = 50 \text{ mm}$?
 - Moon has a diameter of 3476 km
 - Distance to Moon is 384 405 km.
- $s = 384\,405 \text{ km}$, $f = 50 \text{ mm}$, $y = 3\,476 \text{ km}$



- Then

$$y' = y f / (s - f) = 3476 \text{ km} \times 50 \text{ mm} / (384405 \text{ km} - 50 \text{ mm}) = 0.45 \text{ mm}.$$

- This modest image fills less than 0.02 % of the 24 x 36 mm film area!

"The lensmakers equation"

- Lens has two surfaces => **object-image relation** applied twice:

$$\frac{n_a}{s_1} + \frac{n_b}{s_1'} = \frac{n_b - n_a}{R_1}, \quad \frac{n_b}{s_2} + \frac{n_a}{s_2'} = \frac{n_a - n_b}{R_2}$$

s_1 = distance to object, s_2' = distance to final image. $s_2 = -s_1'$.

Set $n_a = 1$, $n_b = n$. We get

$$\frac{1}{s_1} + \frac{n}{s_1'} = \frac{n-1}{R_1}, \quad -\frac{n}{s_1'} + \frac{1}{s_2'} = \frac{1-n}{R_2}$$

- Adding equations:

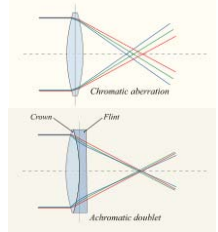
$$\frac{1}{f} = \frac{1}{s_1} + \frac{1}{s_2'} = (n-1) \left(\frac{1}{R_1} - \frac{1}{R_2} \right)$$

- Note that:

- The two focal lengths are always equal, despite different curvatures.

Chromatic aberration

- Index of refraction depends on wavelength => images at different λ focus at different distances.

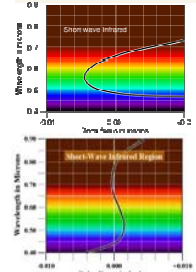


- **Achromatic lens**

- two materials (e.g. crown and flint) bonded
- focus two wavelengths into same focal plane
- Reduces chromatic aberration

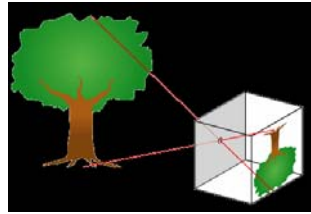
- **Apochromatic lens**

- more than two lenses of different materials
- focus three λ (e.g. R,G,B) into same plane
- order of magnitude better than achromat



The camera (*lat.: small room*)

- A camera consists of:
 - light-tight box
 - lens (several elements)
 - adjustable aperture
 - controllable shutter
 - film or electronic detectors.



pinhole camera

- Fixed focal plane =>
 - lens closer to detector for distant object
 - lens farther away from detector for nearby object.
- This is very different from imaging in the eye.

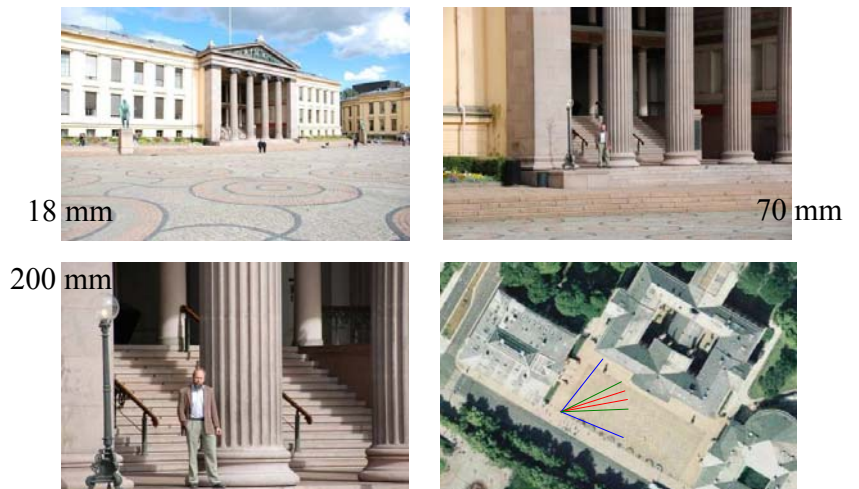
f-number and exposure

- Energy per unit area in the focal plane is proportional to
 - aperture area and length of the exposure time interval.
 - f-number of camera lens $N = f/D$.
 - Intensity in focal plane is proportional to $(D/f)^2$.
 - Changing D by $\sqrt{2}$ changes intensity by a factor of 2.
 - f-numbers are often related by $\sqrt{2}$
 - 1/500 s at $f/4$
 - 1/250 s at $f/5.6$
 - 1/125 s at $f/8$
- } all correspond to the same exposure.
- Shorter exposure times
 - minimize motion-blurring
 - allows larger effective lens aperture, giving
 - better resolution of details in the image
 - reduced **depth of field** and **depth of focus**.

Field of view (FOW)

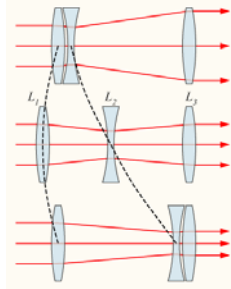
- Using 24 x 36 mm film, FOW measured along the diagonal will be
 - 75° for $f = 28$ mm (wide angle, landscapes)
 - 47° for $f = 50$ mm ("normal")
 - 25° for $f = 105$ mm (ideal for portraits)
 - 8° for $f = 300$ mm (full moon is $1/2^\circ$).
- Replace 24 x 36 mm film with digital detector, => smaller registered field of view.
 - correspond to approximately 1.5 times longer focal lengths in cameras using 24 x 36 mm film.

Same scene, different focal lengths



Zoom lenses

- **Digital zoom:** Cropping and enlarging an image
 - always lower quality than optical zoom, no resolution gained.
- **Optical zoom:** Several lens elements are used
 - focus and focal length can be varied, maintaining focal plane.
- Described by ratio of focal lengths:
 - a 20 to 200 mm zoom is a 10:1 or "10x" zoom.
- Two parts:
 - a fixed-focal-length lens (L3)
 - an **afocal** zoom lens system (L1 + L2)
 - does not focus the light
 - alters the size of a beam
 - alters overall magnification.



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17

The Hitchcock zoom

- Zooming can manipulate perspective in time sequences.
- If the camera is pulled away from the object
 - while lens zooms in to maintain Field Of View, or vice versa,
 - the size of the foreground objects will be constant,
 - but background details will change size relative to foreground.
- Continuous perspective distortion is counter-intuitive:
 - Perspective change without a size change is highly unsettling.
- Special effect used in Hitchcock's Vertigo
 - hence called "Vertigo-" or "Hitchcock-zoom".
 - Also used in "Jaws", "ET", ...



Camera is close, zoomed out



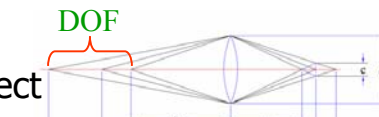
Camera is further away, zoomed in

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Depth of field

- DOF = distance in front of and beyond the focused object that appears to be in focus.
- A large DOF brings both foreground and background into focus.
- A small DOF will focus on part of an interesting object, and defocus a distracting background.



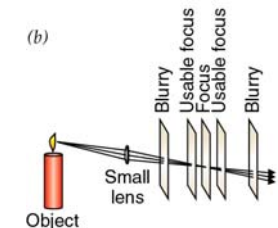
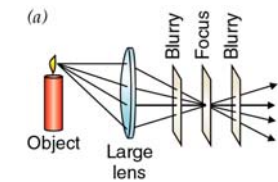
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19

Lens aperture and Depth Of Field

- Large lens aperture D , low f/D - value:
 - Shorter exposure time
 - Focus more critical
 - Better angular resolution
- Small lens aperture D , high f/D -value:
 - Longer exposure time
 - Focus less critical
 - Poorer angular resolution



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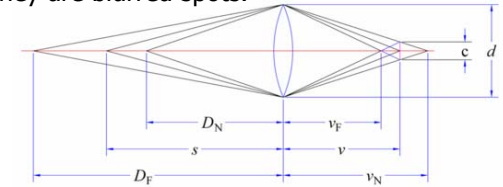
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What determines DOF ?

- DOF is determined by three factors:
 - the focal length of the lens
 - the f-number of the lens aperture
 - the camera-to-object distance.
- Increasing the f-number (smaller aperture) increases the DOF.
 - reduces the amount of light transmitted
 - increases diffraction
 - reduces angular resolution
 => There is a practical limit to the reduction of aperture.

Near and far limits of DOF

- An object at distance s from the lens is in focus at image distance v .
- Objects at D_F and D_N are in focus at image distances v_F and v_N .
- At the image distance v , they are blurred spots.



- When blur spot diameter is equal to the acceptable circle of confusion c (COC), the near and far limits of DOF are at D_N and D_F .
- From similar triangles we see that

$$\frac{v_N - v}{v_N} = \frac{c}{d} \quad \frac{v - v_F}{v_F} = \frac{c}{d}$$

Limits of DOF from $N=f/D$

- "f-number" is given by focal length f and aperture diameter d :

$$N = f/d.$$
- Substituting for d we get the focus limits on the image side of the lens

$$v_N = \frac{fv}{f - Nc} \quad v_F = \frac{fv}{f + Nc}$$

- Thin lens equation:

$$\frac{1}{u} + \frac{1}{v} = \frac{1}{f}$$

- Substituting this give limits of DOF in terms of
 - focal length f
 - "f-number" N
 - object distance s
 - circle of confusion c

$$D_N = \frac{sf^2}{f^2 + Nc(s-f)} \quad D_F = \frac{sf^2}{f^2 - Nc(s-f)}$$

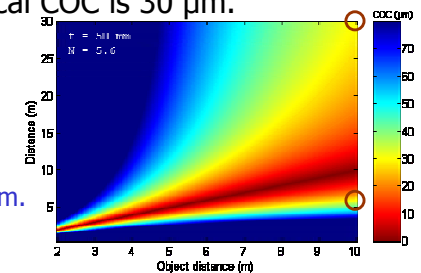
Practical limits of DOF

- DOF beyond the object is always greater than DOF in front of the object.
- For longer focal lengths the ratio tends towards unity.
- For the 35-mm format, a typical COC is 30 μ m.

- A practical example:

$f = 50$ mm
 $N = 5.6$
 $s = 10$ m
 $c = 30$ μ m

- DOF ranges from 6 m to 30 m.
- => DOF extends
 - 4 meters in front of and
 - 20 meters beyond the focus distance.
- A smaller/larger COC gives a larger/smaller DOF.



DOF of a camera zoom lens

- On an old fashioned zoom lens, **far and near limits of DOF** indicated for the chosen **f-number** and **focal length**.



- Hyperfocal distance** is the nearest focus distance at which **the far limit of the DOF extends to infinity**.

The hyperfocal distance

- The hyperfocal distance is the nearest focus distance H at which the far limit of the DOF extends to infinity.**

$$D_F = \frac{s f^2}{f^2 - N c (s - f)}$$

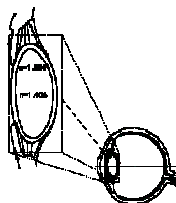
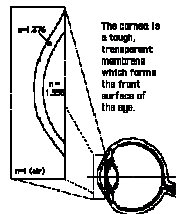
- Setting the far limit D_F to infinity and solving for s gives us H:

$$s = H = \frac{f^2}{N c} + f \approx \frac{f^2}{N c}$$

- Focusing the camera at the hyperfocal distance gives the largest possible DOF for a given f -number.
- You will see the hyperfocal distance again later!

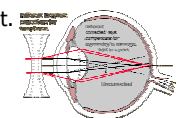
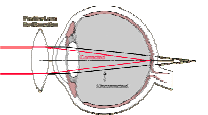
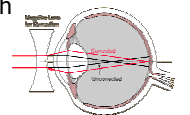
The eye

- Eye is nearly spherical and about 2.5 cm in diameter.
 - cornea* protects the eye, performs much of the focusing
 - iris* and *pupil* controls how much light will be let through
 - lens* produces a sharp image
 - retina* contains photo detectors
 - optic nerve* transmits signals to brain.
- Lens
 - responsible for ca 20 % of the refraction.
 - focal length, $f \approx 1.5$ cm.
- Focusing
 - Radial ligaments around the lens stretch it to a flattened disc
 - focus on far-away objects.
 - If ring-shaped muscle around the radial fibers is relaxed, lens becomes more spherical and its focal length is shortened
 - Focus on closer objects.
 - Lens - retina distance is constant, unlike fixed-focal-length lenses.**
 - Focusing power in "dioptres", d , where $d = 1/f$ (f in meters)**
 - Eye lens ≈ 67 d, cornea responsible for 45 d.**



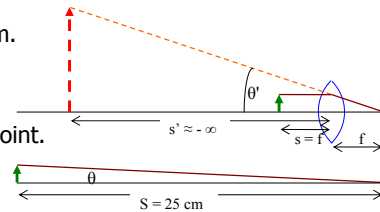
Accommodation and defects

- Accommodation** – an automatic ability to alter the focal length
 - is affected by ageing.
 - young individuals may alter focal power by up to 4 dioptres.
- Presbyopia** - near point recedes as one grows older
- Myopia** (near-sighted) - infinity focused in front of retina.
 - Corrected by diverging lens ($f = 1/d < 0$).
 - moves virtual image of distant object at or inside far point.
- Hyperopia** (far-sighted) - infinity focused behind retina.
 - Corrected by converging lens ($f = 1/d > 0$).
 - forms virtual image of nearby object at or beyond near point.
- Astigmatism** – different focus in horizontal and vertical plane.
 - Remedied by lenses that are not rotationally symmetric.



The magnifier

- Max viewing angle is obtained at ≈ 25 cm.
- Converging lens forms virtual image
 - More distant => larger than object.
- Then object may be moved inside near point.
- Angular magnification M given by
 - ratio of the angle θ' (with the magnifier) to the angle θ (without the magnifier).



- Assume that the eye is relaxed => virtual image placed at infinity.
- Assuming that θ (given in radians) = $\sin(\theta) = \tan(\theta)$:

Magnification:

$$\theta = \frac{y}{25 \text{ cm}} \quad \theta' = \frac{y}{f}$$

$$M = \frac{\theta'}{\theta} = \frac{y/f}{y/25 \text{ cm}} = \frac{25 \text{ cm}}{f}$$

- Simple magnifiers give $M \approx 5$.
- For more magnification, microscope is needed.

The eyepiece

- Concave mirrors and lenses form real images.
- To inspect real image, we may use a second lens.
- This magnifier produce an enlarged virtual image.
 - Found in telescopes and microscopes.
- Telescope and microscope oculars are compound lenses
 - corrected for chromatic and geometric aberrations.



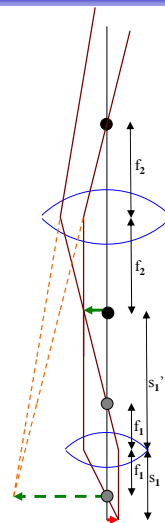
Optical microscope

- Object **just outside** first focal point of the objective => real, inverted and enlarged image.
- This image **just inside** first focal point of eyepiece => a final virtual image.
- Lateral magnification of objective: $m_1 = -s_1'/s_1$
 - If object close to focal point: $m_1 = -s_1'/f_1$.
- Angular magnification of eyepiece: $M_2 = (25 \text{ cm})/f_2$ (if real image is close to focal plane).

- Total angular magnification M:

$$M = m_1 M_2 = -\frac{(25 \text{ cm})s_1'}{f_1 f_2}$$

- Negative sign indicates that image is inverted.
- Use objectives of different f_1 to vary magnification.



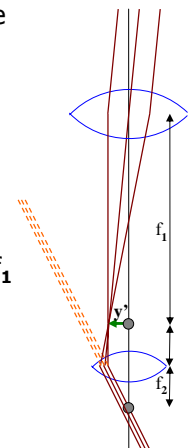
Refracting telescope

- Objective lens forms real, inverted and reduced image
 - distant object => real image at second focal point.
- Real image at the first focal point of eyepiece
 - => final virtual image at infinity
 - => objective to eyepiece distance = sum of focal lengths.
- Angle subtended at the eye by the final virtual image $\theta' = y'/f_2$ (As with the magnifier)
- Angle of object when viewed by unaided eye $\theta = -y'/f_1$

- Magnification:

$$M = \frac{\theta'}{\theta} = -\frac{y'/f_2}{y'/f_1} = -\frac{f_1}{f_2}$$

- Negative sign indicates that image is inverted.
- Use eyepieces of different f_2 to vary magnification.



Multiple lenses

- Focal length for two thin lenses in contact (achromat):

$$\frac{1}{f} = \frac{1}{f_1} + \frac{1}{f_2} \Rightarrow f = \frac{f_1 f_2}{f_1 + f_2}$$

- Back focal length (BFL) of two lenses separated by d :

$$BFL = \frac{f_2(d - f_1)}{d - (f_1 + f_2)} \quad (d \rightarrow 0 \Rightarrow BFL \rightarrow f)$$

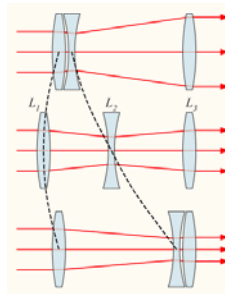
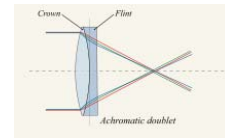
- $d = f_1 + f_2 \Rightarrow$ BFL is infinite. (collimating afocal)

- Alters width of beam
- Found in zoom lenses & simple telescopes

- Magnification given by: $M = -\frac{f_1}{f_2}$

- $f_1 > 0, f_2 > 0 \Rightarrow M < 0$ & inverted image
(two convex lenses).

- $f_1 > f_2 > 0 \Rightarrow M > 0$ & upright image
(convex plus concave lens).



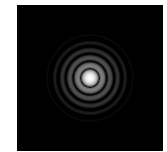
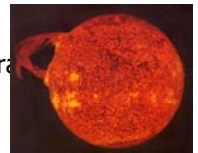
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33

Geometrical Optics: Diffraction

- Fraunhofer diffraction pattern
 - Single slit, twin slit, multiple slits
 - Diffraction grating, spectrograph, spectroheliograph, spectrograph
- Diffraction profile of circular aperture
 - Airy disc and Rayleigh criterion
 - Effect of central obstruction
- The smallest visible detail
 - In a high quality camera
 - In a compact digital camera
 - In a mobile-phone camera
- Depth of focus
- Convolving PSF and sampling aperture



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34

Diffraction by an edge

- Light can bend around corners.
- When a point source casts shadow of a straight edge, the edge of the shadow is not a step edge.
 - some light in the area expected to be in shadow
 - alternating bright and dark fringes in illuminated area.
- Result of interference between many light waves
 - (Huygens' Principle).



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35

Near- and far-field

- Fresnel (near-field)** diffraction
 - Both light source and observation plane are close to the aperture.
 - Curvature of the wave fronts must be taken into account.
 - Fresnel diffraction effects later in the course.
- Fraunhofer (far-field)** diffraction
 - Wave fronts at aperture and observation plane considered planar.
 - usually \Rightarrow light source and observation plane far from slit.
 - We may use collimating lenses
 - lens having light source in its primary focal point will collimate the beam before it reaches the aperture;
 - lens behind the aperture may collimate the beam traveling towards the observation plane.
- The near/far-field limit is at the hyperfocal distance !!!

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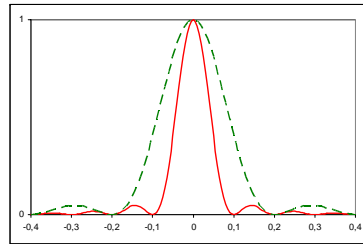
36

Single slit Fraunhofer diffraction

- Diffraction pattern of long, narrow slit:
 - a central bright band,
 - may be much wider than width of slit,
 - bordered by parallel dark and bright bands
 - with rapidly decreasing intensity.

$$I(\theta|a, \lambda) = \left(\frac{\sin[\pi a(\sin \theta) / \lambda]}{\pi a(\sin \theta) / \lambda} \right)^2 I_0$$

a = the width of the slit
 λ = the wavelength
 I_0 = intensity at $\theta = 0$.



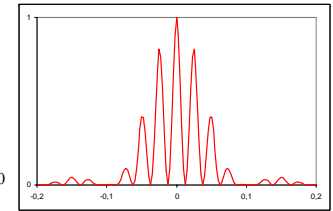
- Visible light: $\lambda < a$, and $\sin(\theta) \approx \theta$.
 \Rightarrow first minimum at $\theta = \lambda/a$.
 \Rightarrow width of central bright band proportional to λ / a .

single slit
 $-\pi/8 < \theta < \pi/8$
 $a = 5 \lambda$ (dashed line)
 $a = 10 \lambda$ (solid line)

Twin slit Fraunhofer diffraction

- Diffraction pattern from two slits
 - a = slit width
 - d = distance between slits
 - λ = monochromatic wavelength
 - I_0 = intensity at $\theta = 0$.

$$I(\theta|a, d, \lambda) = \cos^2(\pi d(\sin \theta) / \lambda) \left(\frac{\sin[\pi a(\sin \theta) / \lambda]}{\pi a(\sin \theta) / \lambda} \right)^2 I_0$$



- Single-slit pattern acts as envelope on cosine interference pattern of two slits.
 - Local maxima may be suppressed by envelope minima.
- Width of envelope proportional to λ / a .

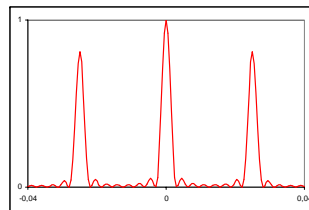
two slits
 $-\pi/16 < \theta < \pi/16$
 $a = 10 \lambda, d = 4a$

Multiple slit diffraction

- Constructive interference for diffracted rays at an angle θ that arrive in observation plane with path length difference = integer number of λ : $d \cdot \sin(\theta) = m \lambda, m=0, \pm 1, \pm 2, \dots$

- Diffraction pattern from N slits
 - a = width of slit
 - d = distance between slits
 - λ = monochromatic wavelength
 - I_0 = intensity at $\theta = 0$.

$$I(\theta|a, d, \lambda, N) = \left[\frac{\sin\left(\frac{N \pi d}{\lambda} \sin \theta\right)}{\sin\left(\frac{\pi d}{\lambda} \sin \theta\right)} \right]^2 \cdot \left[\frac{\sin\left(\frac{\pi a}{\lambda} \sin \theta\right)}{\frac{\pi a}{\lambda} \sin \theta} \right]^2 I_0$$



- Principal maxima
 - same positions as in the two-slit case
 - width proportional to $1/N$.
- N-1 minima between pair of principal maxima
 - secondary maxima get smaller as N increases.

N = 8 slits
 $-\pi/80 < \theta < \pi/80$
 $a/\lambda = 10$
 $d = 4a$.

Diffraction grating

- An assembly of narrow slits - or grooves in planar (or curved) mirror.
- Gratings for $\lambda = 400$ to 700 nm usually have about 1000 lines/mm, corresponding to d on the order of $1/1000$ mm = 1000 nm.
- When a beam is incident on a grating with an angle θ_i (measured from the normal of the grating), it is diffracted into several beams.
 - Specular reflection beam is called zero order ($m = 0$).
 - Other orders given by non-zero integers m in **grating equation**.

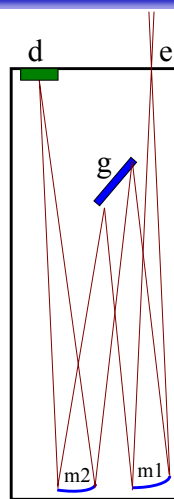
$$d \cdot [\sin \theta_m(\lambda) + \sin \theta_i] = m \lambda, m=0, \pm 1, \pm 2, \dots$$

d = groove period
 λ = wavelength of incident light
 $\theta_m(\lambda)$ = value of the diffracted angle in order m .



A slit spectrograph

- Light is focused onto entrance slit (e).
- Tilted concave collimating mirror (m1) reflects onto a plane reflecting diffraction grating (g).
- Dispersed light of some order m from grating is focused by a second concave mirror (m2) onto detector array (d).
- Special-purpose spectrographs are complex
 - to avoid
 - internal reflections
 - unwanted straylight.



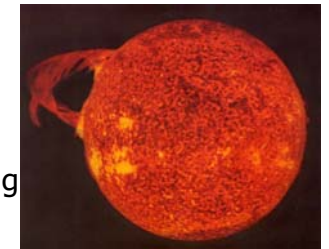
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41

Spectroheliograph

- Produces monochromatic images of the Sun.
- An image of the Sun is focused on a plane.
- A narrow slit lets light into a spectrograph.
- Spectrograph produces a spectrum
 - of the portion of the solar disk imaged on the entrance slit
 - at the same image scale as input image.
- Capture spectrum within an exit slit.
- Image scanned across entrance slit.
- Moving detector behind exit slit at same speed as the image is moving monochromatic image is recorded.



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42

Slitless spectrograph

- Gives co-temporal spectra of all parts of extended objects.
- Concave grating produces image at all wavelengths.
 - Spectral resolution given by reflective grating.
 - Angular resolution given by telescope optics.
- X and λ are same direction.
- Images at different λ overlap.
- Spectral and spatial information mixed into complicated image.



Part of an "overlappograph" image taken by the S082A EUV spectroheliograph on the Skylab Apollo Telescope Mount, January 1, 1974 (JSC Digital Image Collection)

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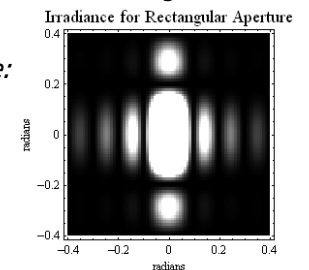
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43

Diffraction by rectangular aperture

- An $a \times b$ aperture gives two orthogonal 1-D diffraction patterns:
- $$I(\theta, \varphi | a, b, \lambda) = \left(\frac{\sin[\pi a(\sin \theta) / \lambda]}{\pi a(\sin \theta) / \lambda} \right)^2 \left(\frac{\sin[\pi b(\sin \varphi) / \lambda]}{\pi b(\sin \varphi) / \lambda} \right)^2 I_0$$
- The widths of the central bright band are inversely proportional to the ratio of the size of the aperture (a, b) to the wavelength λ .

Diffraction pattern of rectangular aperture:
 Horizontal aperture $a = 10 \lambda$
 Vertical aperture $b = 5 \lambda$
 For a ± 0.4 radians range of θ and φ .



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44

Diffraction profile of circular aperture

- Diffraction limits resolution and image quality.
- Diffraction by circular aperture will image point source as a small bright disc (Airy-disc)
 - surrounded by dark and bright rings,
 - intensity of the decrease rapidly
- Angular point spread function (PSF) of circular aperture given by:

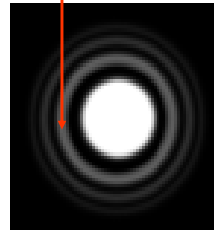
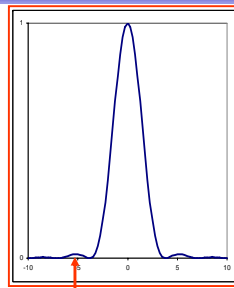
$$F_0(\theta) = 4 \left[\frac{J_1\left(\pi \frac{\theta}{\beta_0}\right)}{\pi \frac{\theta}{\beta_0}} \right]^2$$

θ given in radians,

$\beta_0 = \lambda/D$

J_1 = the first order Bessel function

$F_0(0) = 1$.



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45

Airy disc and diffraction rings

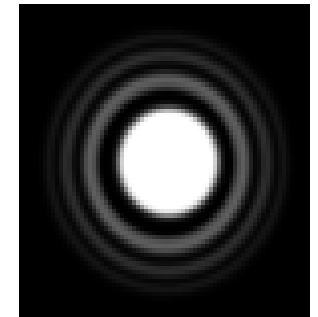
- Angular radii of first **dark** rings:

$$\sin \theta_1 = 1.22 \frac{\lambda}{D} \quad \sin \theta_2 = 2.23 \frac{\lambda}{D} \quad \sin \theta_3 = 3.24 \frac{\lambda}{D}$$

- Angular radii of **bright** rings:

$$\sin \theta = 1.63 \frac{\lambda}{D} \quad \sin \theta = 2.68 \frac{\lambda}{D} \quad \sin \theta = 3.70 \frac{\lambda}{D}$$

- About 85% of energy within Airy disk
- peak intensity of 1. ring $\approx 1.7\%$
- peak intensity of 2. ring $\approx 0.4\%$



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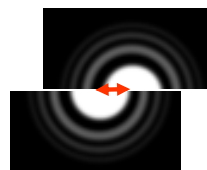
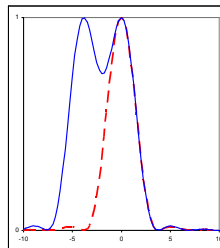
46

Rayleigh criterion

- Cross-section through diffraction-limited image
 - two equally bright point sources at infinity
 - angular separation between point sources:

$$\sin \theta = 1.22 \frac{\lambda}{D}$$

- Corresponds to overlaying two patterns
- maximum of first on first minimum of second.
- 27 % "dip" between the peaks.



- This is the "Rayleigh criterion".
 - For small angles, $\sin \theta = \theta$.

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47

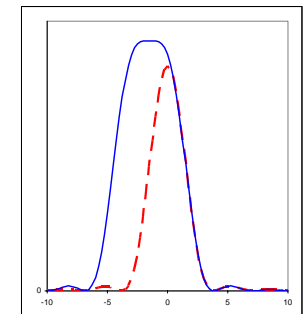
Sparrow criterion

- If the point sources are moved closer than the Rayleigh criterion, the "dip" will become shallower, until it becomes a flat plateau.

- This angular separation is the Sparrow criterion.

- The limit when two point sources "melt together".

- Sparrow: $\theta = 0.952 \lambda/D$.



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48

Circular aperture with central obstruction

- The diffraction profile of an annular aperture is given by the continuous circular symmetric function

$$F(\theta | \delta) = \frac{1}{(1 - \delta^2)^2} \left[\left(\frac{2J_1(v)}{v} \right) - \delta^2 \left(\frac{2J_1(\delta v)}{\delta v} \right) \right]^2$$

J_1 is the first order Bessel function

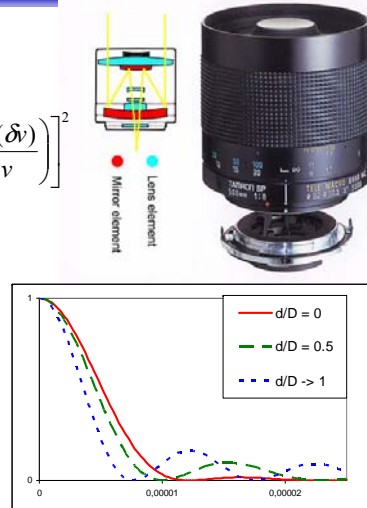
$v = \pi \theta D / \lambda$

θ is given in radians

$\delta = d/D$

d = diameter of central obstruction
 D = diameter of the aperture

- Central obstruction gives
 - a slightly narrower central disk
 - more energy in the bright rings
 - some energy is blocked out.



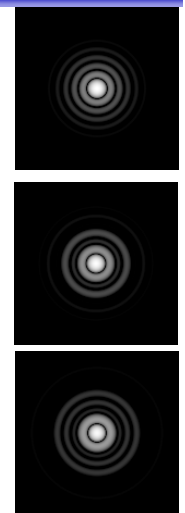
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49

Airy patterns with central obstruction

- Airy pattern for obstructions of 0 %, 20 %, 33 %
- From 0 to 1/3 radial obstruction, maximum intensity of first diffraction ring increases from 1.7 % to 5.4 %.
- Obstruction affects both
 - the position of the minima
 - intensity of maxima in the ring pattern.



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50

Smallest detail visible ... to the eye

- Pupil diameter ≈ 2 mm in bright light.
- Angular resolution of human eye ≈ 60 lines per degree.
 - 120 alternating black and white lines of equal thickness.
- A4 "landscape" paper at 30 cm distance covers $50^\circ \times 40^\circ$.
 - 3 000 black and 3 000 white vertical stripes should be resolved
 - 2 400 black and 2400 white horizontal stripes should be resolved.
- Rayleigh criterion for $D=2.5$ mm at $\lambda=550$ nm:

$$\sin \theta = \frac{1.22 \lambda}{D} = \frac{1.22 \times 550 \times 10^{-9}}{2.5 \times 10^{-3}} = 2.7 \times 10^{-4}$$

- Converting from radians to degrees:

$$\theta = 2.7 \times 10^{-4} \frac{180}{\pi} \approx \frac{1}{60}^\circ$$

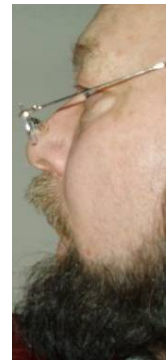
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51

Accommodation distance ... Near point

- The closest distance we may focus sharply with the unaided eye.
 - 7 cm for a 10 year old,
 - 10 cm for a 20 year old,
 - 14 cm for a 30 year old,
 - 22 cm for a 40 year old,
 - 40 cm for a 50 year old,
 - 100 cm for a 60 year old.
- 6 000 dots / 11 inches, ≈ 550 dpi.
- 47 yrs $\Rightarrow s \approx 30$ cm.
 - printer better than 600 dpi is a waste.
- 20 yrs $\Rightarrow s = 10$ cm.
 - can inspect the printout at 10 cm distance
 - Will need 1 200 dpi (common for printing high quality images).



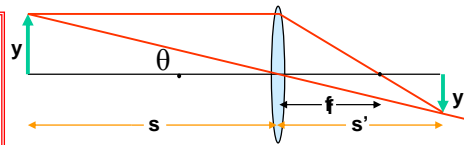
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52

Smallest detail visible by a camera

- Focal length $f = 35$ mm.
- $f/D = 3 \Rightarrow D = f/3$.
- Object is $s = 3.5$ meters away.
- Wavelength is $\lambda = 500$ nm.

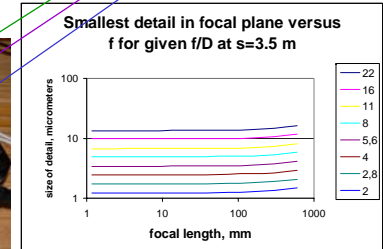
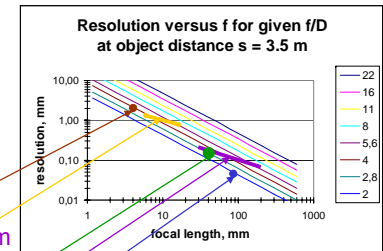


- Just resolvable angle, as seen from center of lens, given by Rayleigh criterion:
 $\sin(\theta) = 1.22 \lambda/D = 1.22 \cdot 500 \cdot 10^{-9} \cdot 3 / 35 \cdot 10^{-3} = 5.23 \cdot 10^{-3}$.
- Q:** What is distance y between two just resolved points on object ?
- A:** Given by: $\tan(\theta) = (y/s)$.
 For small angles $\sin(\theta) = \tan(\theta) = \theta$, when θ given in radians.
 $y = 3.5 \cdot 5.23 \cdot 10^{-3} \text{ m} = 1.83 \cdot 10^{-4} \text{ m} \approx \mathbf{0.2 \text{ mm}}$.
- Q:** What is the corresponding distance y' in the focal plane?
- A:** $y' = y \cdot f / (s-f)$
 $y' = 0.2 \cdot 35 / (3500-35) \text{ mm} \approx 0.002 \text{ mm} = \mathbf{2 \mu\text{m}}$.

Resolution and detail – medium distance

- Object and image resolution depends on aperture and distance.
- At $s = 3.5$ meters:

- Mobile phone camera: $\Delta y = 2.1$ mm
- Compact digital zoom: $\Delta y = 0.6-1.1$ mm
- Compact 35 mm: $\Delta y = 0.2$ mm
- Digital SLR zoom: $\Delta y = 0.06-0.3$ mm
- "old fashioned" 85 mm: $\Delta y = 0.05$ mm



Object resolution versus distance

- For a given lens and f/D , size of smallest resolvable object detail increases linearly with distance.
- Keeping f/D constant, size of smallest resolvable object detail decreases linearly with focal length.
- Keeping focal length constant, size of smallest resolvable object detail increases linearly with f/D .

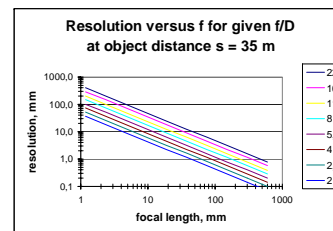
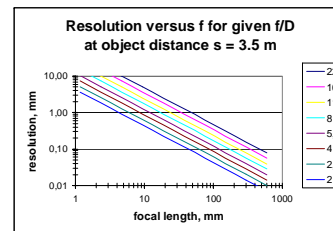
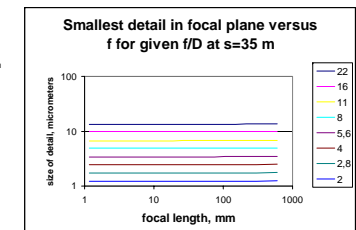
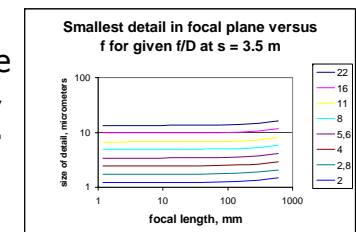


Image resolution versus distance

- For given lens and f/D , size of smallest resolvable detail in image is independent of object distance, except when f is comparable to s .
- Keeping f/D constant, size of smallest resolvable object is independent of focal length, except when f is comparable to s .
- Keeping focal length constant, size of smallest resolvable object detail increases linearly with f/D .



Diffraction limited depth of focus

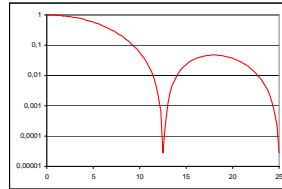
- Intensity of diffraction pattern near focus:

$$I(u,v) = \left(\frac{2}{u}\right)^2 \left[1 + V_0^2(u,v) + V_1^2(u,v) - 2V_0(u,v) \cos\left\{\frac{1}{2}\left(u + \frac{v^2}{u}\right)\right\} - 2V_1(u,v) \sin\left\{\frac{1}{2}\left(u + \frac{v^2}{u}\right)\right\} \right] I_0$$

u is along optical axis, v orthogonal to optical axis.

- Along optical axis:

$$I(u,0) = \left(\frac{2}{u}\right)^2 \left[2 - 2\cos\left\{\frac{u}{2}\right\} \right] I_0 = \left(\frac{\sin(u/4)}{u/4}\right)^2 I_0$$



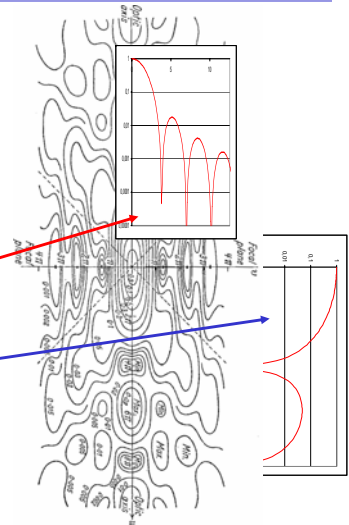
- Symmetric about the focal plane.

- First minimum at $u/4 = \pm \pi$, $\Rightarrow z = \pm 2\lambda(f/D)^2$ from focal plane.
- $f/D = 5.6$, $\lambda = 500 \text{ nm} \Rightarrow z = \pm 0.032 \text{ mm}$
- $f/D = 22$, $\lambda = 500 \text{ nm} \Rightarrow z = \pm 0.5 \text{ mm}$
- $f/D = 5.6$, $\lambda = 1000 \text{ nm} \Rightarrow z = \pm 0.25 \text{ mm}$.

- Depth of focus is independent of physical dimensions of the optics.**

Intensity along optical axis

- A contour plot of the intensity $I(u,v)$ in a meridional plane near the focus of a converging spherical wave diffracted by a spherical aperture.



- Vertical u-axis is the optical axis, and the horizontal v-axis is in the focal plane.

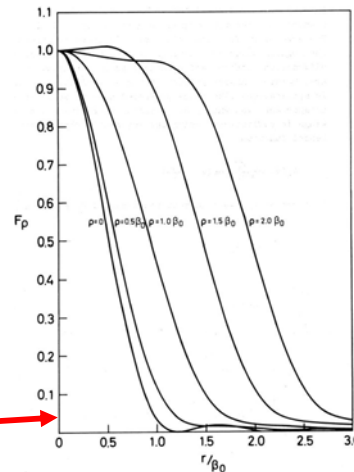
- The maxima and minima along the v-axis correspond to bright and dark rings of focal plane diffraction pattern.

- Maxima and minima along the u-axis illustrate "depth-of-focus".

- Contour plot from M. Born and E. Wolf: "Principles of Optics", Pergamon Press, 4th. Ed., 1970.

Convolving PSF and sampling aperture

- PSF determines resolution.
- Ideal:
 - Point sampling in focal plane
 - Sampling theorem \Rightarrow density
- Reality:
 - Extended sampling aperture
 - Fixed, non-overlapping
 - Movable, overlapping
 - Rectangles
 - Circles
 - Sampling aperture * PSF



PSF convolved with circular aperture

Sampling aperture smooths and broadens PSF

Geometrical Optics: Scattering

- What is scattering?
- Some effects of scattering
 - Atmospheric blurring and straylight in images
 - Turbidity in liquids
 - Subsurface scattering in non-metallic materials and in tissues
- Doppler-shifted straylight

What is scattering?

- Scattering causes radiation to deviate from a straight trajectory.
 - microscopic irregularities in surfaces
 - non-uniformities in transparent media
- *Elastic scattering* : no (or a very small) loss or gain of energy
- *Inelastic scattering* : some change in energy
- *Absorption* : substantial or complete loss of energy
- *Single scattering* : one localized scattering center.
 - treated as a random phenomenon, described by probability distribution.
- *Multiple scattering* : radiation is scattered several times.
 - randomness of interaction averaged out by large number of events
 - => deterministic angular distribution of intensity - PSF.
- Observed blurred image:
 - *convolution* of true image with PSF (diffraction + scattering).

Inverse scattering problem

- A difficult challenge!
- Observe blurred object + scattering around it.
- Determine
 - scattering parameters (PSF)
 - distribution of radiation before scattering (true object).
- In general, the inverse is not unique.
- PSF can be determined by observing image of some well-known object through the same scattering medium.
- PSF then used to deconvolve image of unknown object.

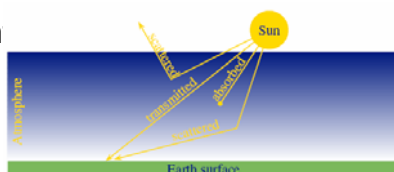
Wavelength dependence of scattering

- **Rayleigh scattering**
 - in transparent solids, liquids and gases.
 - wavelength dependence $\sim \lambda^{-4}$
- **Blue sky**: Blue light is scattered much more than red light.
 - We observe blue light coming from all directions of the sky.
 - At higher altitudes, less scattering particles => sky is much darker.
- **Red sunset**: Sunlight must pass through greater air mass.
 - More scattering of blue light, little scattering of red light => red-hued sky.
- **Mie scattering**
 - scattering by spheres larger than Rayleigh range.
 - wavelength dependence $\sim 1/\lambda$.
 - shape of scattering center significant
 - theory only applies well to spheres, spheroids and ellipsoids.



Atmospheric scattering

- Object is illuminated by sunlight
- Incident radiation detector:
 - Specular reflection
 - Diffuse reflection
 - radiation scattered in the air:
 - scattered before reaching object
 - specular and diffuse reflection, scattered onto detector.
- Important to shield detector to minimize straylight.
- Even with shielding, scattering will be present.
- Corrections important
 - In high precision measurements (e.g., astrophysics, ...)
 - Remote sensing (radiation passing twice through atmosphere)



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65

High density of scatterers

- Vapors:
 - an object that is seen through mist or fog will look blurred.
 - at some distance the object will disappear into the background fog.
- Water:
 - Particles / organisms act as scatterers
 - cause haziness that indicate water quality
 - **turbidity** can be measured using *Secchi disk*
 - lowered into water until it can no longer be seen.
- Translucent solids:
 - light penetrates non-metallic surface and scatters inside material
 - either absorbed or leaving the material at a different location.
 - This phenomenon is called **subsurface scattering (SSS)**.
 - The effect is a “softer” image than a metallic surface would give.
- Tissues:
 - human skin, salmon fillets, etc show **subsurface scattering**
 - may depend on wavelength, condition of tissue, etc.
 - Thus, measuring SSS may be useful for
 - quality inspection of e.g., fish and meat
 - medical diagnostic work.



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66

Straylight integral

- Straylight causes a blurring of the image (here the Sun).
- Given
 - circular symmetric PSF, $\Psi(r)$,
 - true intensity $\Phi_c(p')$
- Observed intensity $I(p)$ given by integral equation:

$$I(p) = \int_{\oplus} \Phi_c(p') \Psi(r) d\omega$$

- p and p' are directions in the sky
- r is the angle between them.
- Integration is performed over the solid angle of the Sun.

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67

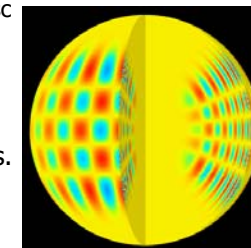
Doppler shifted straylight

- Different parts of the Sun have different line-of-sight velocities.
- => Observed intensity contain straylight with different Doppler velocities.

$$I(d, \lambda) = 2 \int_{\rho_0}^{\rho_1} \int_0^{\alpha_0} \Phi_c(a) (1 - I_c(a)) \exp\left[-(\lambda - \Delta\lambda)^2 / w^2(d)\right] \Psi(\rho) d\rho da$$

- Φ_c = true continuum intensity distribution across the solar disc
- w = Doppler width of Gaussian absorption line profile
- I_c = central intensity of absorption line
- d = distance from the centre of the solar disc
- λ = wavelength within a spectral line
- $\Psi(r)$ = circular symmetric PSF, :

- Straylight introduce errors = 0.1 – 1.0 m/s,
- Amplitudes of global solar oscillations \approx 0.1 m/s.
- => Error \approx velocity oscillation signal.



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68