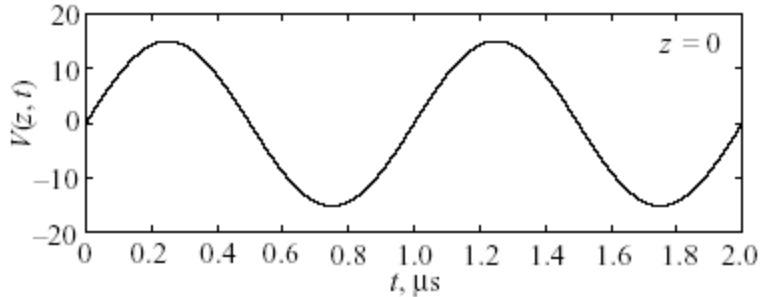
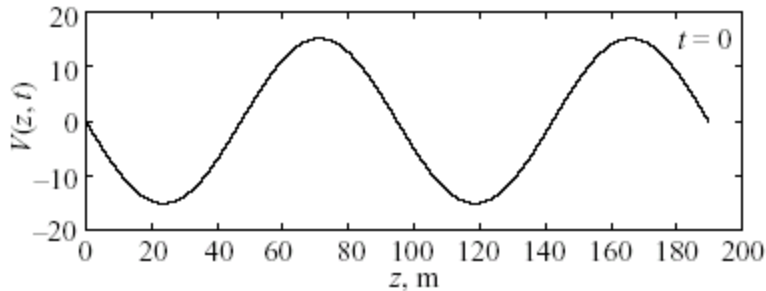


Ch. 2. Transmission Line Analysis



Traveling voltage wave

$$V(z, t) = \frac{E_{0x}}{\beta} \sin(\omega t - \beta z)$$

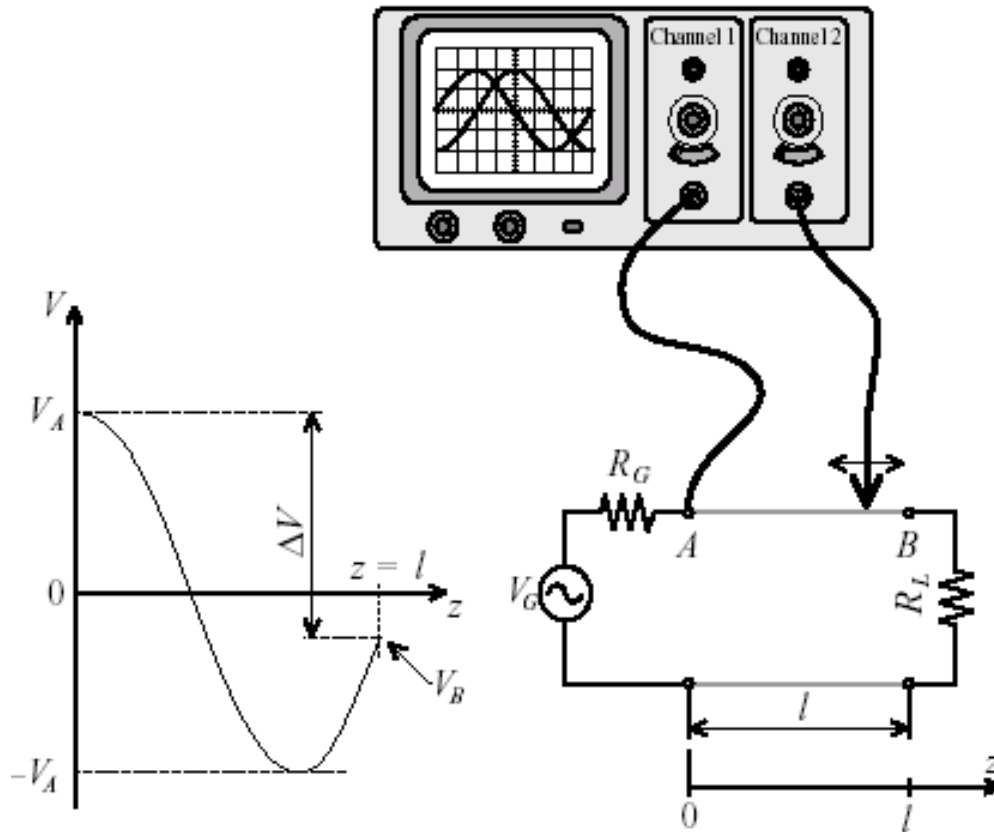


Phase velocity

$$v_p = \omega / \beta = 1 / \sqrt{\mu \epsilon} = c / \sqrt{\epsilon_r}$$

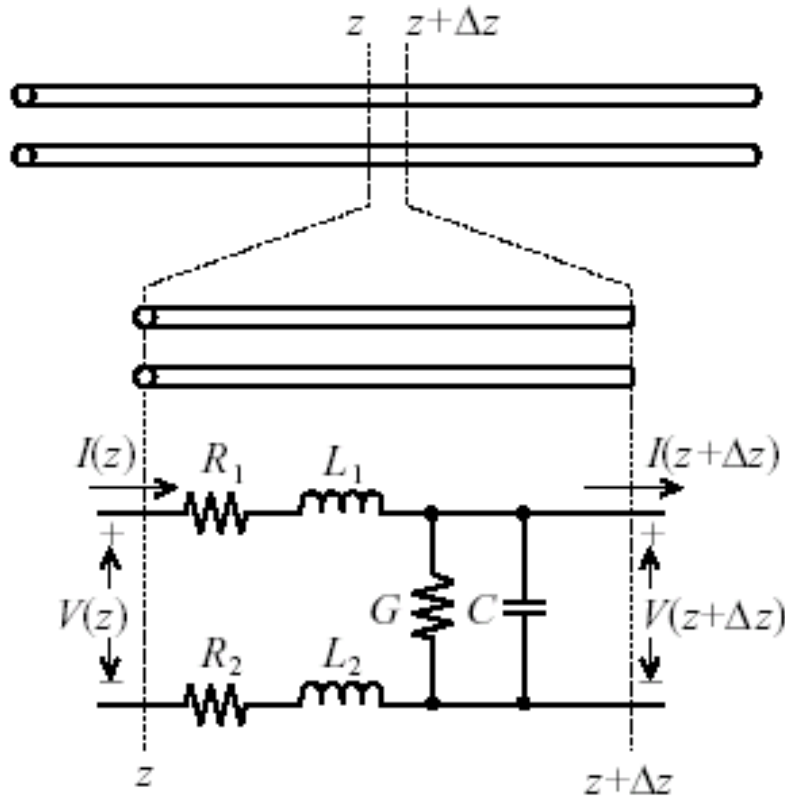
- Voltage has a time and space variation
- Space is neglected for low frequency applications
- For RF there can be a large spatial variation

Consequences of spatial voltage variations



- For low frequency (1MHz) Kirchhoff's laws apply
- For high frequency (1GHz) Kirchhoff's laws do **not apply** anymore
- **Solution:** Consider elements of infinitesimal length

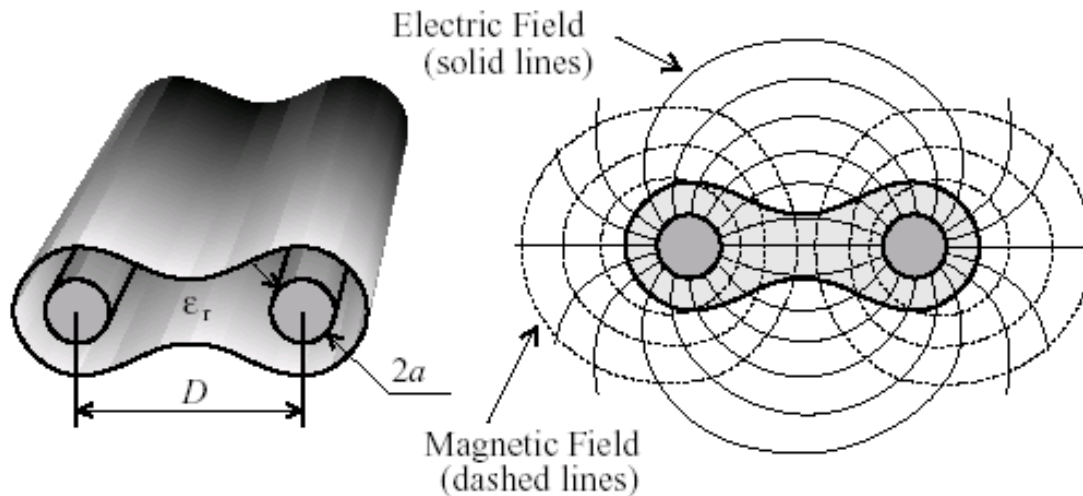
Kirchhoff's laws on a microscopic level



- Over a differential section we can again use basic circuit theory
- Model takes into account line losses and dielectric losses
- Ideal line (lossless) involves only L and C

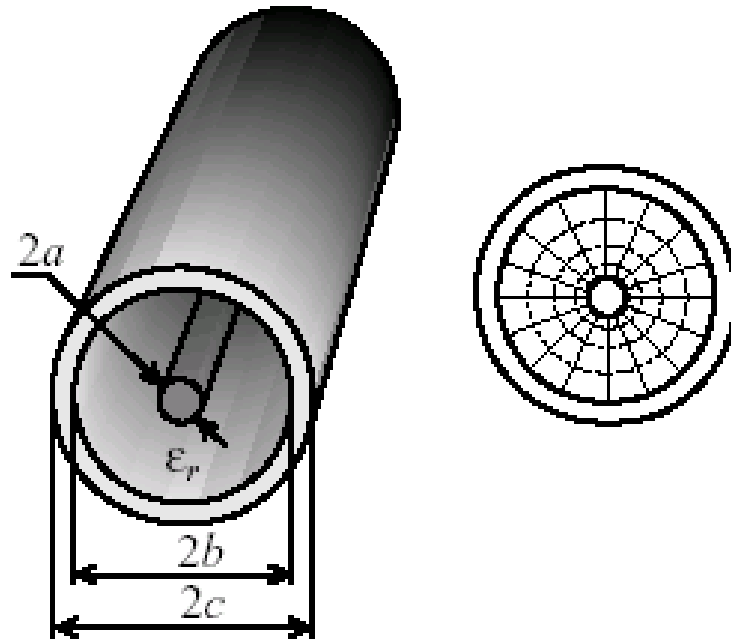
Distributed parameters R, G, L and C

Two-wire transmission line



- Alternating electric field between conductors
- Alternating magnetic field surrounding conductors
- dielectric medium tends to confine field inside material

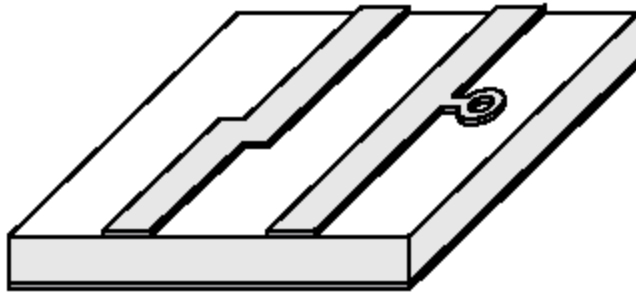
Coaxial line



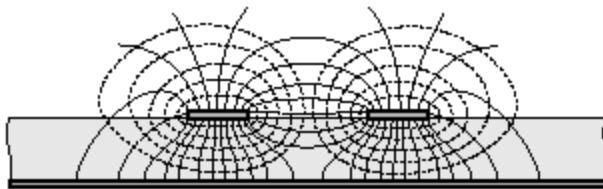
Always used for externally connected RF systems or measuring equipment.
Also LAN.

- Electric field contained between conductors
- Perfect shielding of magnetic field
- TEM mode up to a certain cutoff frequency

Microstrip lines

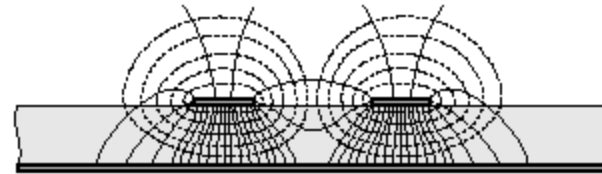


Printed circuit board (PCB) section with ground plane to prevent excessive field leakage, interference, and radiation loss



(a) Teflon epoxy ($\epsilon_r = 2.55$)

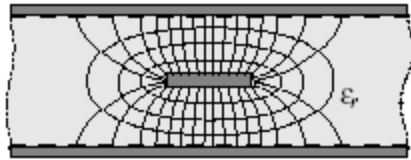
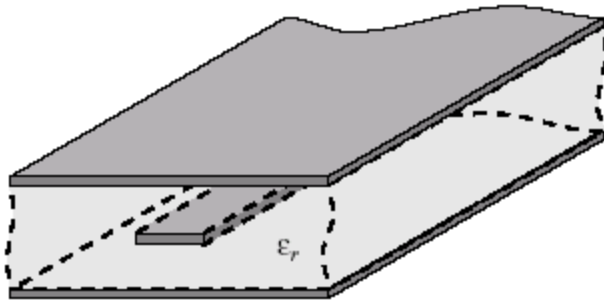
Low dielectric medium



(b) Alumina ($\epsilon_r = 10.0$)

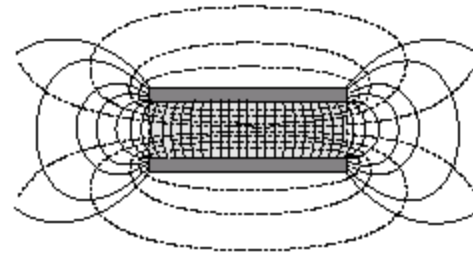
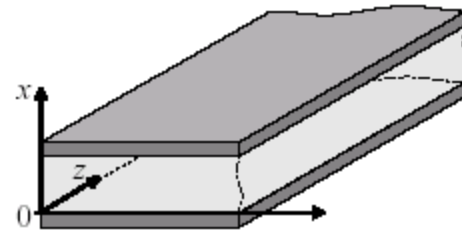
High dielectric medium

Other TEM configurations



Triple-layer line

Reduced radiation losses

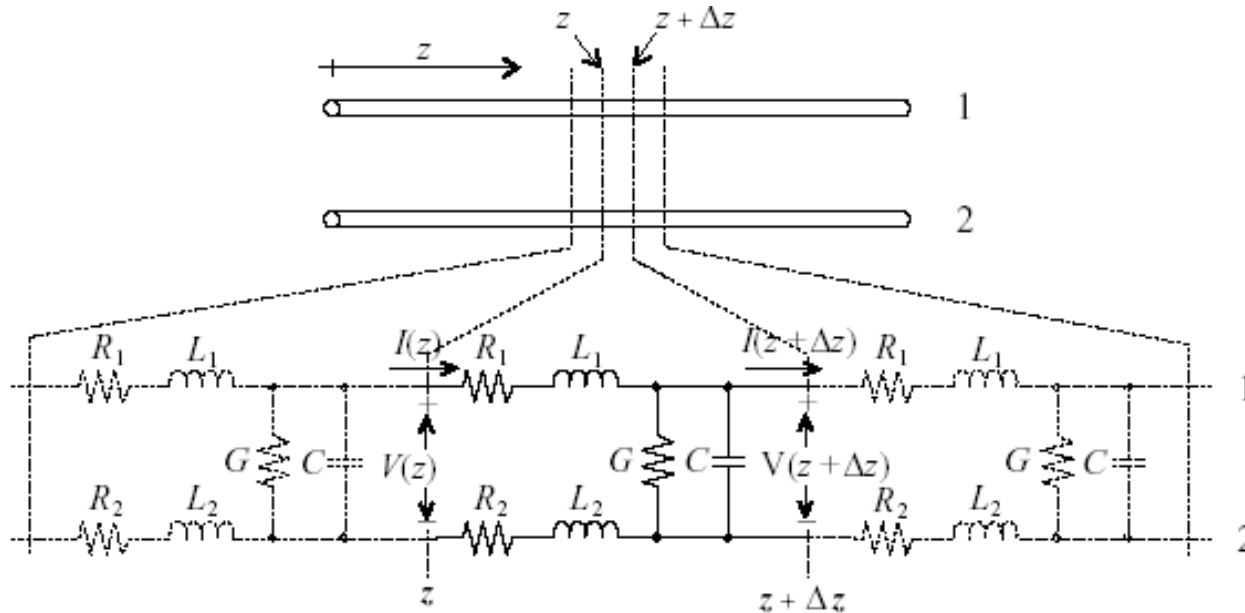


Parallel plate line

Low impedance, high power

Transmission line representation

- Detailed analysis is based on differential section



- Analysis applies to many types of transmission lines such as coax cables, two-wire, microstrip, etc.

Pros and cons of electric circuit representation

- Clear intuitive physical picture
- Yields a standardized two-port network representation
- Serves as building blocks to go from microscopic to macroscopic forms
- Basically a one-dimensional representation (cannot take into account interferences)
- Material nonlinearities, hysteresis, and temperature effects are not accounted for

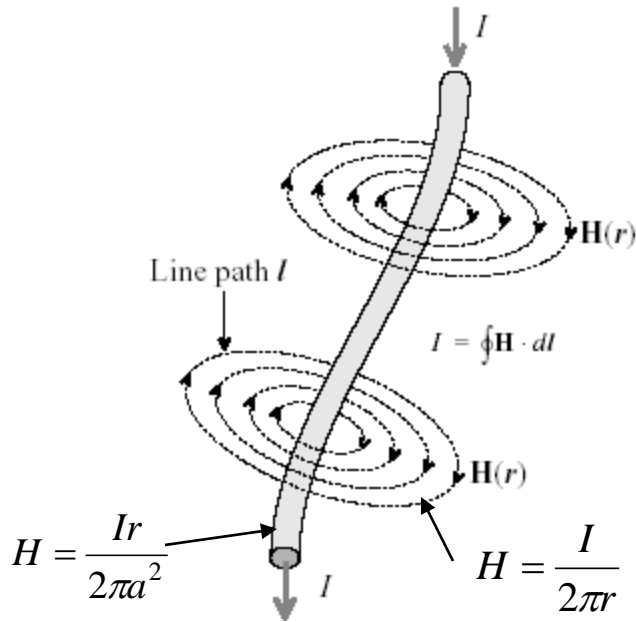


Basic electromagnetism

- Ampère's law

$$\oint \mathbf{H} \cdot d\mathbf{l} = \iint \mathbf{J} \cdot d\mathbf{S}$$

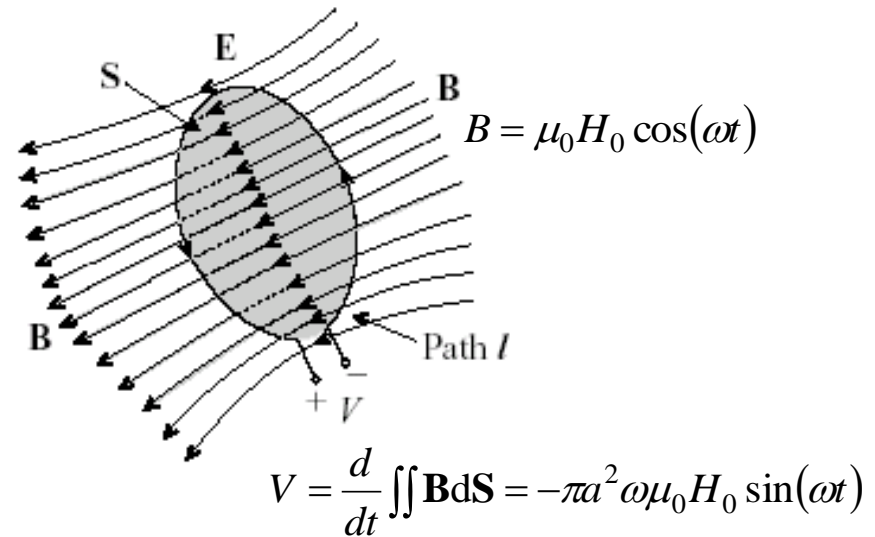
$$\Rightarrow \nabla \times \mathbf{H} = \mathbf{J}$$



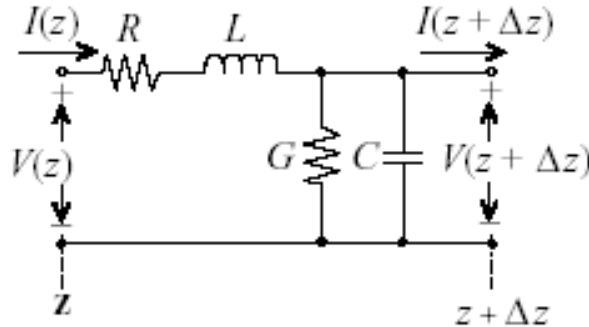
- Faraday's law

$$\oint \mathbf{E} \cdot d\mathbf{l} = -\frac{d}{dt} \iint \mathbf{B} \cdot d\mathbf{S}$$

$$\Rightarrow \nabla \times \mathbf{E} = -\frac{d\mathbf{B}}{dt}$$



Line parameters for specific cases



Generic electric equivalent circuit representation

Parameter	Two-Wire Line	Coaxial Line	Parallel-Plate Line
R Ω/m	$\frac{1}{\pi a \sigma_{\text{cond}} \delta}$	$\frac{1}{2\pi \sigma_{\text{cond}} \delta} \left(\frac{1}{a} + \frac{1}{b} \right)$	$\frac{2}{w \sigma_{\text{cond}} \delta}$
L H/m	$\frac{\mu}{\pi} \cosh^{-1} \left(\frac{D}{2a} \right)$	$\frac{\mu}{2\pi} \ln \left(\frac{b}{a} \right)$	$\mu \frac{d}{w}$
G S/m	$\frac{\pi \sigma_{\text{diel}}}{\cosh^{-1} (D/(2a))}$	$\frac{2\pi \sigma_{\text{diel}}}{\ln (b/a)}$	$\sigma_{\text{diel}} \frac{w}{d}$
C F/m	$\frac{\pi \epsilon}{\cosh^{-1} (D/(2a))}$	$\frac{2\pi \epsilon}{\ln (b/a)}$	$\epsilon \frac{w}{d}$

Check out example 2.3 in text book to get feeling of numbers!

General transmission line equation

KVL:

$$V(z) = (R + j\omega L)I(z)\Delta z + V(z + \Delta z)$$

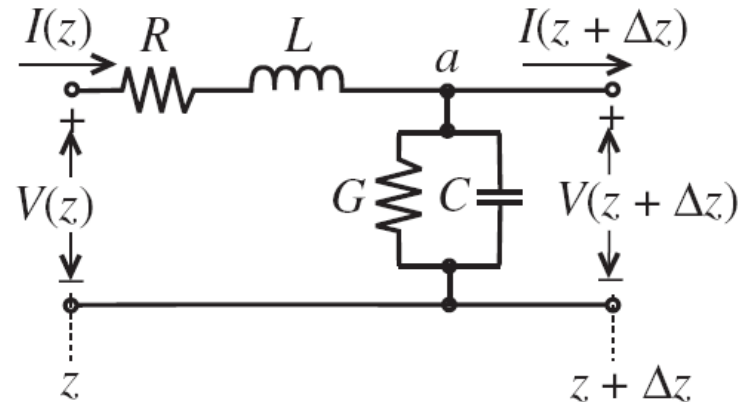
$$\xrightarrow{\Delta z \rightarrow 0} \frac{dV(z)}{dz} = -(R + j\omega L)I(z)$$

KCL:

$$I(z) = V(z + \Delta z)(G + j\omega C)\Delta z + I(z + \Delta z)$$

$$\xrightarrow{\Delta z \rightarrow 0} \frac{dI(z)}{dz} = -(G + j\omega C)V(z)$$

⇒ Coupled first-order differential equations



Traveling voltage and current waves

$$\frac{d^2V(z)}{dz^2} - \gamma^2V(z) = 0$$

$$\frac{d^2I(z)}{dz^2} - \gamma^2I(z) = 0$$

Complex propagation constant:

$$\gamma = \alpha + j\beta = \sqrt{(R + j\omega L)(G + j\omega C)}$$

Solutions:

$$V(z) = V^+e^{-\gamma z} + V^-e^{+\gamma z}$$

$$I(z) = I^+e^{-\gamma z} + I^-e^{+\gamma z}$$



General line impedance definition

$$\frac{dV(z)}{dz} = -(R + j\omega L)I(z) \Rightarrow I(z) = \frac{\gamma}{R + j\omega L} (V^+ e^{-\gamma z} - V^- e^{+\gamma z})$$

Characteristic line impedance:

$$Z_0 \equiv \frac{R + j\omega L}{\gamma} = \sqrt{\frac{R + j\omega L}{G + j\omega C}} = \frac{V^+}{I^+} = -\frac{V^-}{I^-}$$

$$\Rightarrow I(z) = \frac{1}{Z_0} (V^+ e^{-\gamma z} - V^- e^{+\gamma z})$$

Note: Z_0 is not a conventional impedance, but is a characteristic of the positive and negative traveling waves



Lossless transmission line

Lossless implies:
 $R = 0$ and $G = 0$

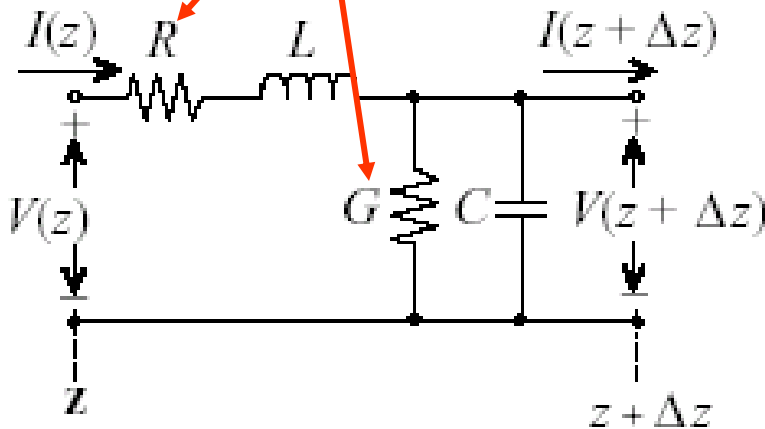
Consider lossless parallel plate transmission line with:

$$L = \mu \frac{d}{w} \quad \text{and} \quad C = \epsilon \frac{w}{d}$$

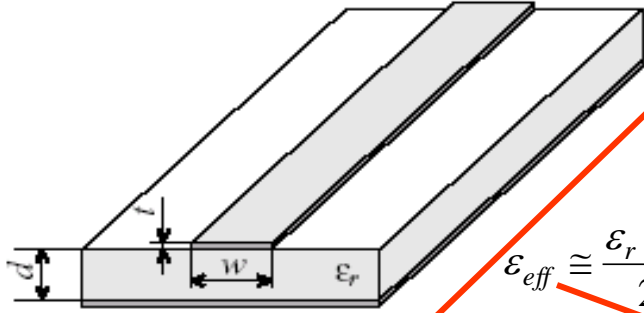
Characteristic impedance:

$$Z_0 = \sqrt{\frac{L}{C}} = \sqrt{\frac{\mu}{\epsilon} \frac{d}{w}} \cong \frac{377 \Omega}{\sqrt{\epsilon_r}} \frac{d}{w}$$

$Z_f = \sqrt{\mu_0 / \epsilon_0} \cong 377 \Omega$ is the wave impedance of free space

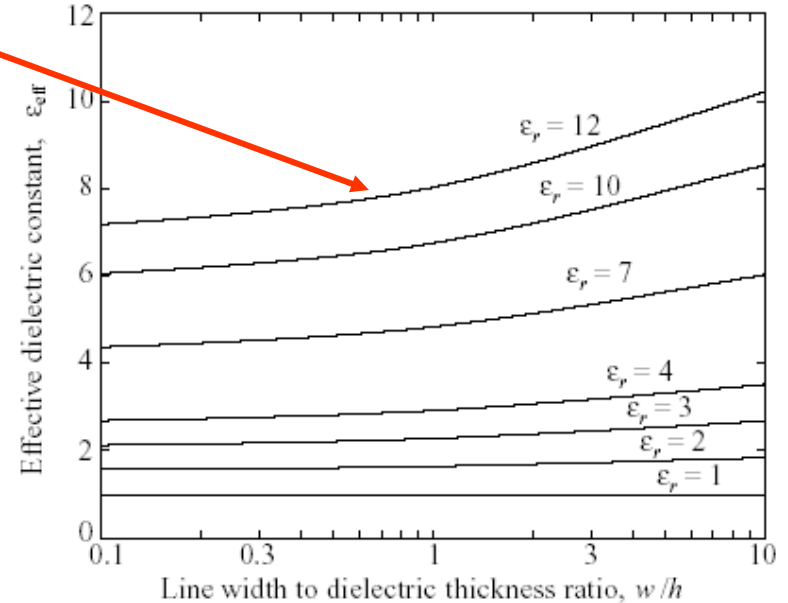
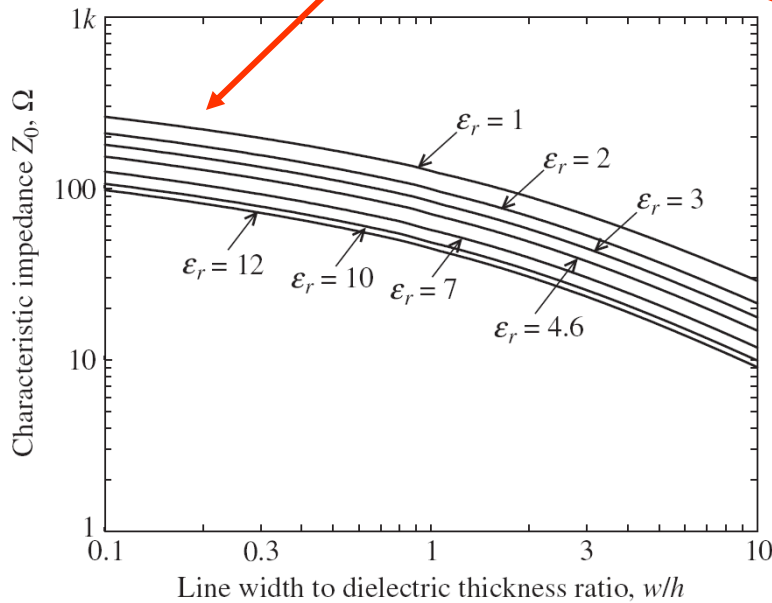


Microstrip transmission line



$$Z_0 \cong \frac{Z_f}{2\pi\sqrt{\epsilon_{ff}}} \ln\left(\frac{8h}{w} + \frac{w}{4h}\right) \Big|_{w < h} \quad \text{or} \quad \frac{Z_f}{\sqrt{\epsilon_{ff}} \left(1.393 + \frac{w}{h} + \frac{2}{3} \ln\left(1.444 + \frac{w}{h}\right)\right)} \Big|_{w > h}$$

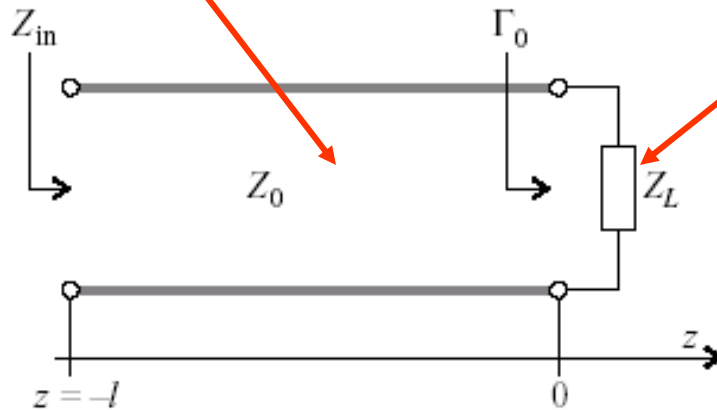
$$\epsilon_{eff} \cong \frac{\epsilon_r + 1}{2} + \frac{\epsilon_r - 1}{2} \left[\left(1 + \frac{12h}{w}\right)^{-\frac{1}{2}} + 0.04 \left(1 - \frac{w}{h}\right)^2 \right] \Big|_{w < h} \quad \text{or} \quad \frac{\epsilon_r + 1}{2} + \frac{\epsilon_r - 1}{2} \left(1 + \frac{12h}{w}\right)^{-\frac{1}{2}} \Big|_{w > h}$$



Terminated lines - Voltage reflection coefficient

$$V(z) = V^+ e^{-\gamma z} + V^- e^{+\gamma z}$$

$$I(z) = (V^+ e^{-\gamma z} - V^- e^{+\gamma z}) / Z_0$$



Load impedance:

$$Z(0) = Z_L = \frac{V(0)}{I(0)} = Z_0 \frac{V^+ + V^-}{V^+ - V^-} = Z_0 \frac{1 + \Gamma_0}{1 - \Gamma_0}$$

Reflection coefficient:

$$\Gamma_0 \equiv \frac{V^-}{V^+} = \frac{Z_L - Z_0}{Z_L + Z_0}$$

Open line: $Z_L \rightarrow \infty$, $\Gamma_0 = 1$ Wave fully reflected with same polarity as incident wave

Short circuit: $Z_L = 0$, $\Gamma_0 = -1$ Wave fully reflected with opposite polarity of incident wave

Load match: $Z_L = Z_0$, $\Gamma_0 = 0$ No reflection when load matches line impedance

Lossless transmission line

For lossless line ($R = G = 0$): $\gamma = \alpha + j\beta = j\omega\sqrt{LC}$

$$\alpha = 0 \quad \beta = \omega\sqrt{LC}$$

Voltage and current waves:

$$V(z) = V^+ \left(e^{-j\beta z} + \Gamma_0 e^{+j\beta z} \right) \quad I(z) = \frac{V^+}{Z_0} \left(e^{-j\beta z} - \Gamma_0 e^{+j\beta z} \right)$$

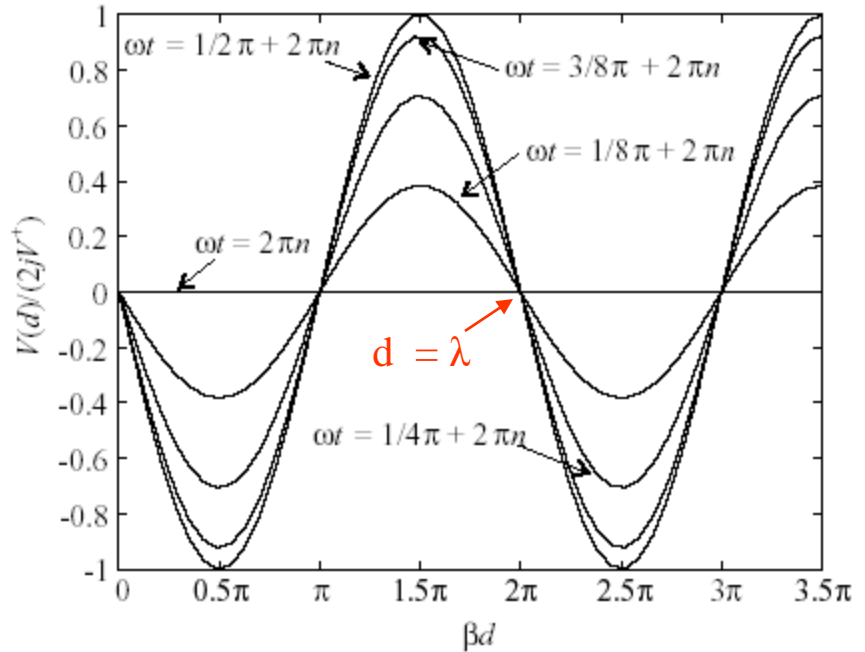
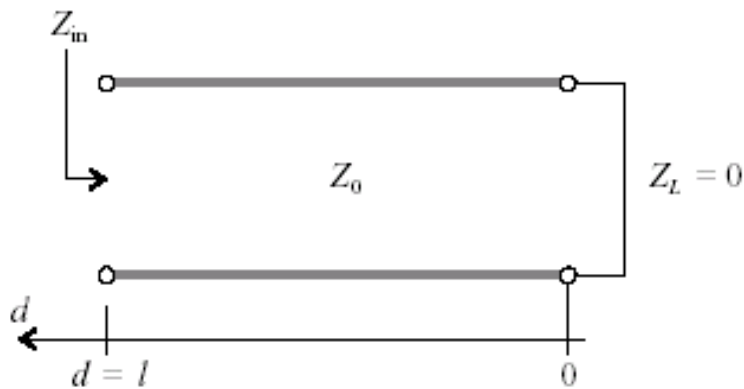
Phase velocity:

$$v_p = \lambda f = \frac{\omega}{\beta} = \frac{1}{\sqrt{LC}}$$



Standing waves

Short circuit: $Z_L = 0$, $\Gamma_0 = -1$
 Wave fully reflected with opposite polarity of incident wave



Note: $d = -z$

\Rightarrow Standing wave pattern:

$$V(d) = V^+ (e^{+j\beta d} - e^{-j\beta d}) = 2jV^+ \sin(\beta d)$$

$$v(d, t) = \text{Re}\{V e^{j\omega t}\} = 2V^+ \sin(\beta d) \cos(\omega t + \pi/2)$$

Standing wave ratio (SWR)

Generally: $V(d) = V^+ e^{+j\beta d} (1 + \Gamma_0 e^{-2j\beta d}) = A(d)[1 + \Gamma(d)]$

$$I(d) = \frac{A(d)}{Z_0} [1 - \Gamma(d)]$$

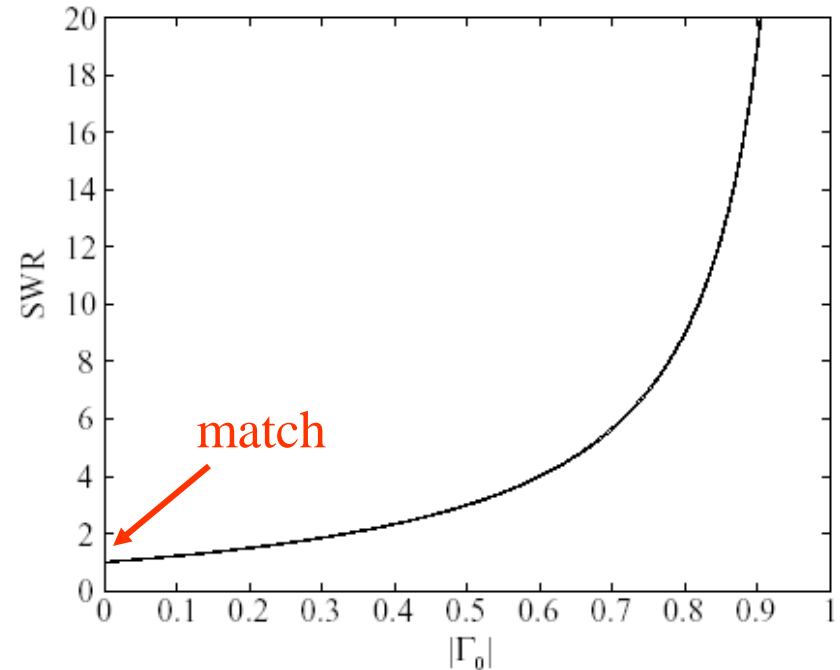
Reflection coefficient: $\Gamma(d) = \Gamma_0 e^{-2j\beta d}$

$$SWR = \frac{|V_{\max}|}{|V_{\min}|} = \frac{|I_{\max}|}{|I_{\min}|} = \frac{1 + |\Gamma_0|}{1 - |\Gamma_0|}$$

SWR is a measure of mismatch of the load to the line

SWR=1 (matched)

SWR $\rightarrow \infty$ (total mismatch)



Note: SRW applies to lossless lines, but also works well in low-loss cases

Transformation of load impedance

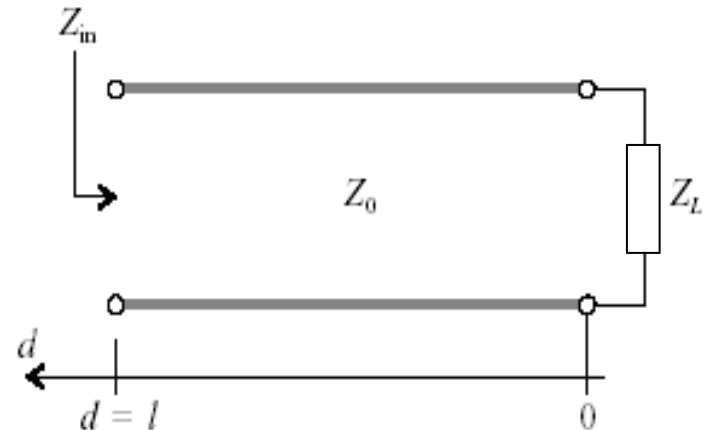
- Terminated lossless transmission line ($Z_0 = \sqrt{L/C}$)

- Input impedance:

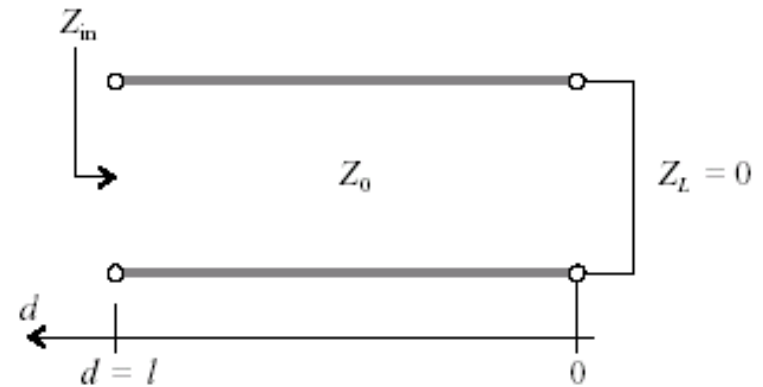
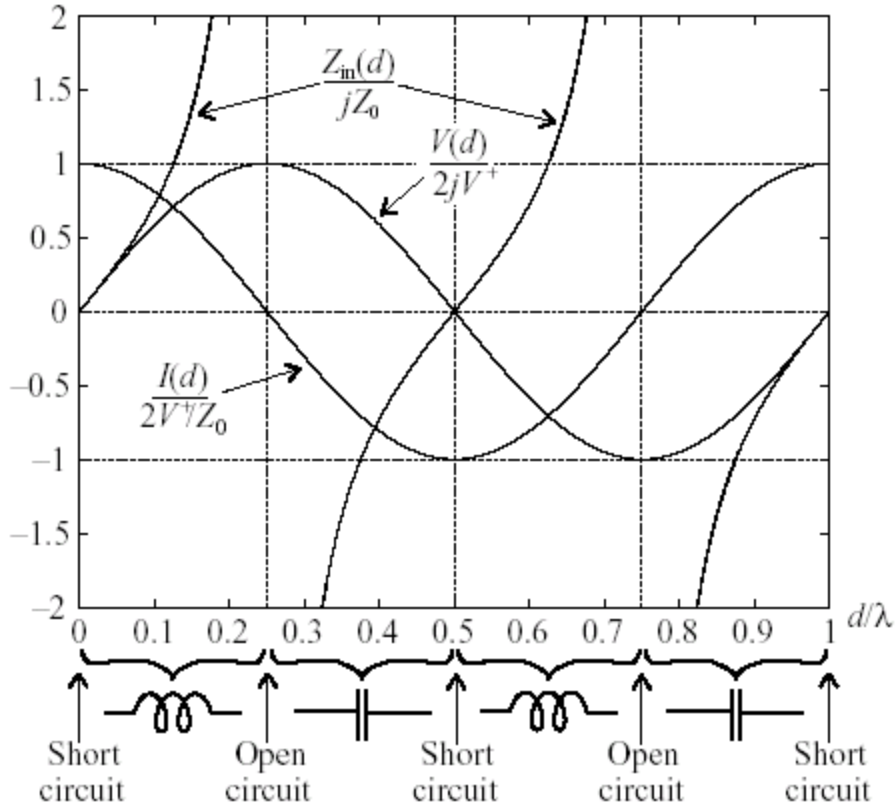
$$Z_{in}(d=l) = \frac{V(l)}{I(l)} = Z_0 \frac{e^{j\beta l} + \Gamma_0 e^{-j\beta l}}{e^{j\beta l} - \Gamma_0 e^{-j\beta l}} = Z_0 \frac{(Z_L + Z_0)e^{j\beta l} + (Z_L - Z_0)e^{-j\beta l}}{(Z_L + Z_0)e^{j\beta l} - (Z_L - Z_0)e^{-j\beta l}}$$

- Used: $\Gamma_0 \equiv \frac{V^-}{V^+} = \frac{Z_L - Z_0}{Z_L + Z_0}$

$$\Rightarrow Z_{in}(d) = Z_0 \frac{Z_L + jZ_0 \tan(\beta d)}{Z_0 + jZ_L \tan(\beta d)}$$



Short circuit transmission line

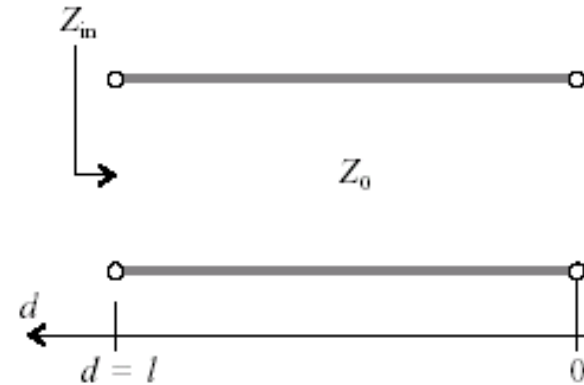
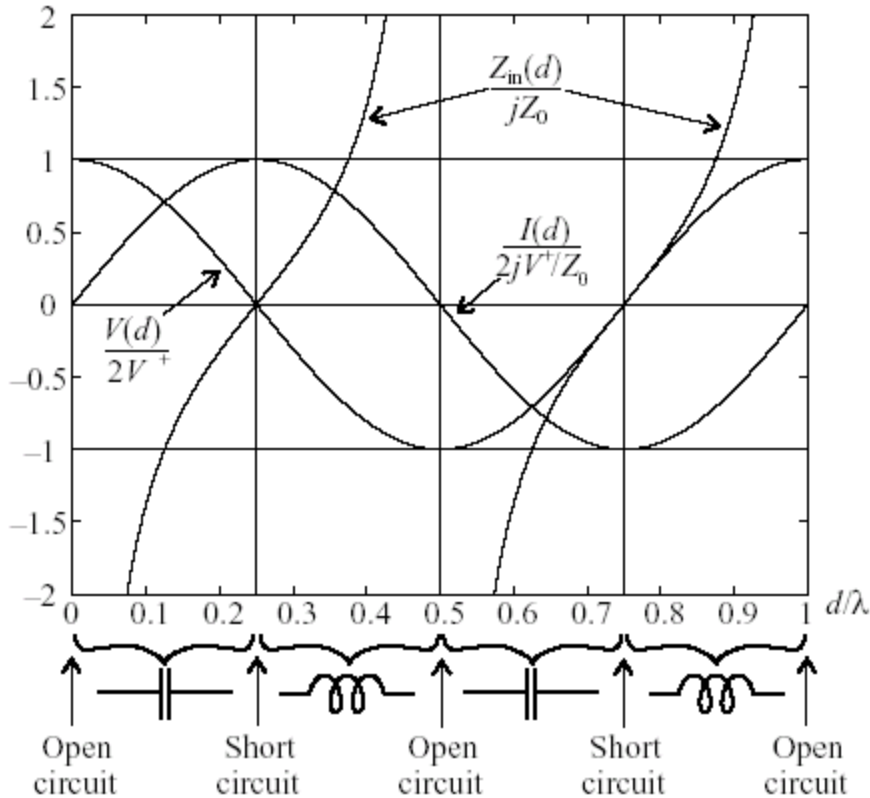


$$Z_{in}(d) = jZ_0 \tan(\beta d)$$

$$V(d) = 2jV^+ \sin(\beta d)$$

$$I(d) = \frac{2V^+}{Z_0} \cos(\beta d)$$

Open circuit transmission line

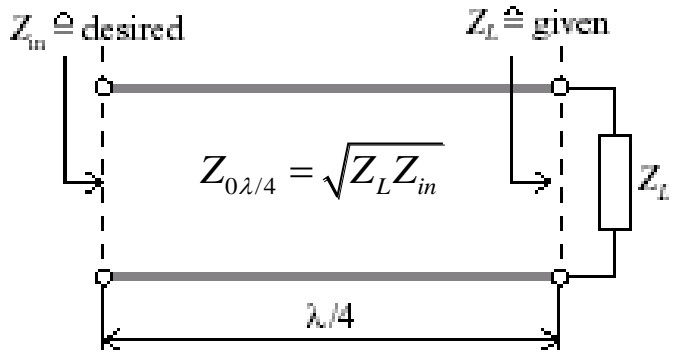


$$Z_{in}(d) = -jZ_0 \cot(\beta d)$$

$$V(d) = 2V^+ \cos(\beta d)$$

$$I(d) = \frac{2jV^+}{Z_0} \sin(\beta d)$$

Quarter-wave transmission line

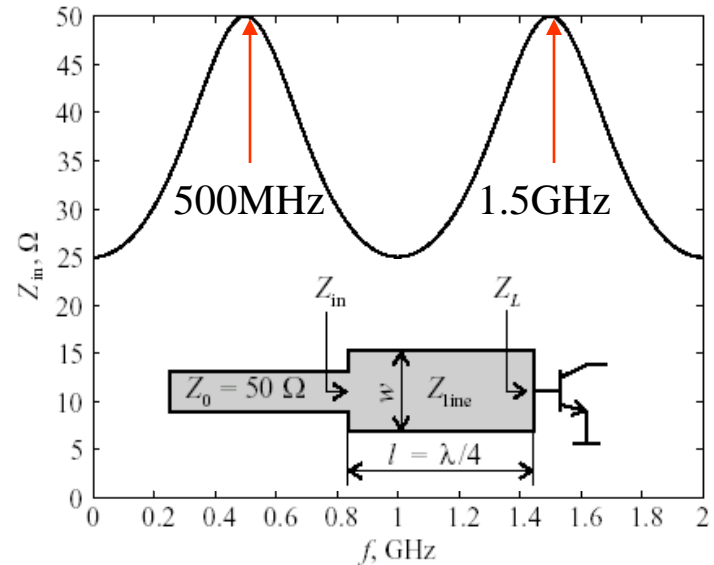


For $d = \lambda/4$ we have:

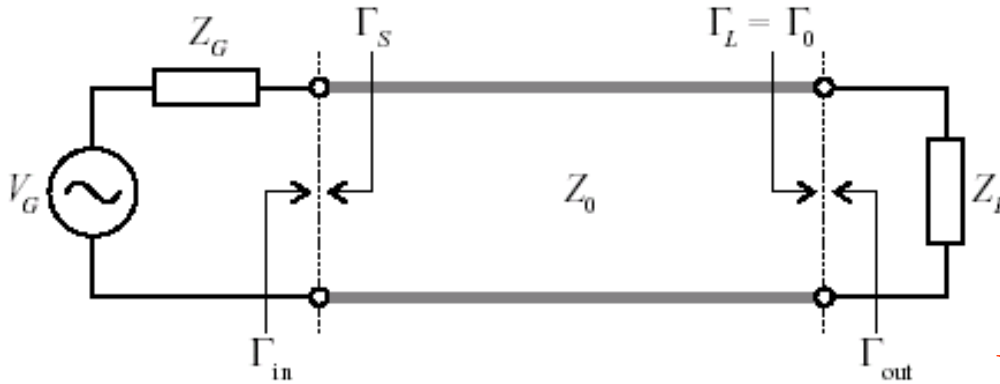
$$Z_{in} \left(d = \frac{\lambda}{4} \right) = Z_{0\lambda/4} \frac{Z_L + jZ_{0\lambda/4} \tan \left(\frac{2\pi}{\lambda} \frac{\lambda}{4} \right)}{Z_{0\lambda/4} + jZ_L \tan \left(\frac{2\pi}{\lambda} \frac{\lambda}{4} \right)} = \frac{Z_{0\lambda/4}^2}{Z_L}$$

Lambda-quarter transformer
 matches given input and output impedances by choosing a line with characteristic impedance (narrowband matching, $Z_{in} = Z_0$):

$$Z_{0\lambda/4} \equiv Z_{line} = \sqrt{Z_L Z_0}$$



Transmission line with source and load



Also have to consider the impedance matching at the source!

Input voltage ($d = l$):

$$V_{in} = V_{in}^+ + V_{in}^- = V_{in}^+(1 + \Gamma_{in}) = V_G \frac{Z_{in}}{Z_{in} + Z_G}$$

Transmission coefficients:

$$T_{in} = 1 + \Gamma_{in} = \frac{2Z_{in}}{Z_{in} + Z_0}$$

$$T_0 = 1 + \Gamma_0 = \frac{2Z_L}{Z_L + Z_0}$$

Reflection coefficients:

$$\Gamma_{in} = \Gamma(d = l) = \frac{Z_{in} - Z_0}{Z_{in} + Z_0} = \Gamma_0 e^{-2j\beta l}$$

$$\Gamma_S = \frac{Z_G - Z_0}{Z_G + Z_0} \quad \Gamma_{out} = \Gamma_S e^{-2j\beta l}$$

Input Power

$$P_{in} = \frac{1}{2} \operatorname{Re}\{V_{in} I_{in}^*\} = \frac{1}{2} \frac{|V_{in}^+|^2}{Z_0} (1 - |\Gamma_{in}|^2)$$

$$V_{in}^+ = \frac{V_{in}}{1 + \Gamma_{in}} = \frac{V_G}{1 + \Gamma_{in}} \frac{Z_{in}}{Z_{in} + Z_G} \quad Z_{in} = Z_0 \frac{1 + \Gamma_{in}}{1 - \Gamma_{in}} \quad Z_G = Z_0 \frac{1 + \Gamma_S}{1 - \Gamma_S}$$

$$P_{in} = \frac{1}{8} \frac{|V_G|^2}{Z_0} \frac{|1 - \Gamma_S|^2}{|1 - \Gamma_S \Gamma_{in}|^2} (1 - |\Gamma_{in}|^2)$$

Lossless TL:

$$P_{in} = \frac{1}{8} \frac{|V_G|^2}{Z_0} \frac{|1 - \Gamma_S|^2}{|1 - \Gamma_S \Gamma_0 e^{-2j\beta l}|^2} (1 - |\Gamma_0|^2)$$



Lossless TL - special cases

Load and source matched to line:

$$\Gamma_0 = \Gamma_s = 0 \Rightarrow P_{in} = \frac{1}{8} \frac{|V_G|^2}{Z_0}$$

Maximum available power
provided by the source

Match at load and mismatch at source:

$$\Gamma_0 = 0 \Rightarrow P_{in} = \frac{1}{8} \frac{|V_G|^2}{Z_0} |1 - \Gamma_s|^2$$

Power usually measured in dBm: $P[dBm] = 10 \log \left(\frac{P[W]}{1mW} \right)$



Return and insertion losses

$$RL = -10\log\left(\frac{P_r}{P_i}\right) = -10\log|\Gamma_{in}|^2 = -20\log|\Gamma_{in}|$$

$$IL = -10\log\left(\frac{P_t}{P_i}\right) = -10\log\left(\frac{P_i - P_r}{P_i}\right) = -10\log(1 - |\Gamma_{in}|^2)$$

