

# Excitability

## Excitable Cells 5.1

Unlike other cells, excitable cells can be triggered to set off an action potential.

During the action potential the transmembrane potential departs from its resting potential, reaches a peak potential and returns to the resting potential after some time.

Nerve cells and cardiac cells uses the action potential as a signal to neighboring cells.

The trigger must be of a certain size, if below the threshold the cell will not “fire” .

# The Hodgkin-Huxley Model

Developed to study the action potential of the squid nerve cells.

Assumed three different current  $I_{Na}$ ,  $I_K$  and  $I_L$

Assumed also linear current-voltage relationship:

$$C_m \frac{dv}{dt} = -I_{ion} + I_{app} = -g_{Na}(v - v_{Na}) - g_K(v - v_K) - g_L(v - v_L) + I_{app}$$

Can collect the current terms due to linearity:

$$C_m \frac{dv}{dt} = -g_{\text{eff}}(v - v_{\text{eq}})$$

where

$$g_{\text{eff}} = g_{\text{Na}} + g_{\text{K}} + g_{\text{L}}$$

and

$$v_{\text{eq}} = \frac{g_{\text{Na}} v_{\text{Na}} + g_{\text{K}} v_{\text{K}} + g_{\text{L}} v_{\text{L}}}{g_{\text{eff}}}$$

$v_{\text{eq}}$  is a weighted average of the individual equilibrium potentials.  
The weighing factors are time and voltage dependent.

A steady applied current  $I_{\text{app}}$  moves the membrane potential to different equilibrium.

$$C_m \frac{dv}{dt} = -g_{\text{eff}}(v - v_{\text{eq}}) + I_{\text{app}} = 0$$

Implies

$$v = v_{\text{eq}} + \frac{1}{C_m g_{\text{eff}}} I_{\text{app}}$$

The applied current will be compensated by an ionic current going the opposite way, thus the net current will be zero.

For a sufficiently large  $I_{\text{app}}$ ,  $v$  will pass the threshold potential and an action potential is triggered. The conductivities will vary greatly.

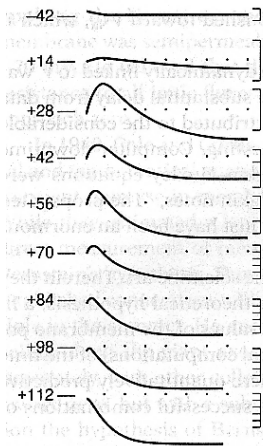
# Voltage Clamp measurements

The transmembrane potential is forced by an applied current to a fixed value.

Since  $I_{\text{ion}} = -I_{\text{app}}$  for a fixed  $v$ , we can measure  $I_{\text{ion}}$  as a function of time for a given level of  $v$ .

Since  $v$  is fixed the observed variations must be due to temporal variation in the conductivities.

## Total membrane current for different steps, 5.1.2



**Figure 4.2** Experimental results describing the total membrane current in response to a step depolarization. The numbers on the left give the final value of the membrane potential, in mV. The interval between dots on the horizontal scale is 1 ms, while one division on the vertical scale represents 0.5 mA/cm<sup>2</sup>. (Hodgkin and Huxley, 1952a, Fig. 2a.)

# From measurements to models

Initially, Hodgkin and Huxley assumed  $I_{\text{ion}} = I_{\text{Na}} + I_{\text{K}}$ . Two kind of experiments conducted:

- 1: Normal concentrations
- 2:  $[\text{Na}]_e$  replaced by cohline  $\Rightarrow$  affects  $I_{\text{Na}}$  but not  $I_{\text{K}}$ .

Assumed further:

- Initially  $I_{\text{K}} = 0$
- $I_{\text{Na}}^1 / I_{\text{Na}}^2 = C$ , constant
- $I_{\text{K}}^1 = I_{\text{K}}^2$

Once  $I_{\text{ion}}^1$  and  $I_{\text{ion}}^2$  is recorded we can determine  $C$  from the first and the second assumptions.



Expressions for the currents in terms of measurable quantities can now be obtained:

$$I_{\text{Na}}^1 = \frac{C}{C-1} (I_{\text{ion}}^1 - I_{\text{ion}}^2)$$

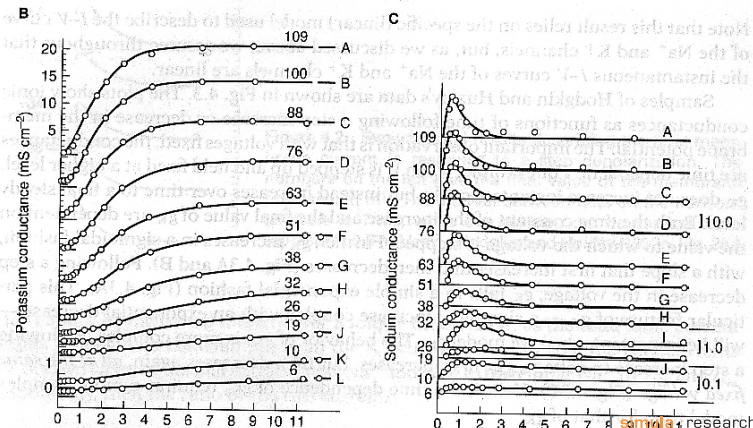
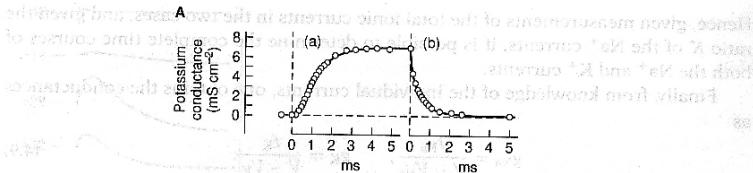
$$I_{\text{K}} = \frac{1}{1-C} (I_{\text{ion}}^1 - CI_{\text{ion}}^2)$$

Assuming linear current-voltage relationships we get expressions for the conductivities:

$$g_{\text{Na}} = \frac{I_{\text{Na}}}{V - V_{\text{Na}}}, \quad g_{\text{K}} = \frac{I_{\text{K}}}{V - V_{\text{K}}}$$

For each pair of voltage clamp experiment (with a given voltage step), we now have a time course for  $g_{\text{Na}}$  and  $g_{\text{K}}$ .

# Potassium and Sodium conductance



# Model for the Potassium conductance

Assumed  $\frac{dg_K}{dt} = f(v, t)$ .

Ended up with introducing a second variable:

$$g_K = \bar{g}_K n^4, \text{ with } \frac{dn}{dt} = \alpha(v)(1 - n) - \beta(v)n$$

and  $\bar{g}$  is the maximum conductance. Forth power was chosen to get the correct shape of the solution.

The solution of

$$\tau_n \frac{dn}{dt} = n_\infty - n$$

with constant coefficients is

$$n(t) = n_\infty + (n(0) - n_\infty)e^{-t/\tau_n}$$

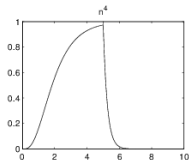
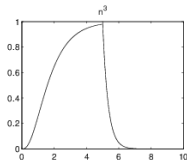
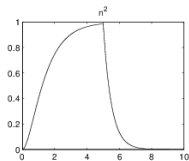
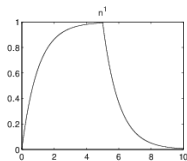
If we assume that  $n_\infty(0) = 0$  a step from from 0 to  $v$  yields:

$$\begin{aligned} n(t) &= n_\infty(v) + (n_\infty(0) - n_\infty(v))e^{-t/\tau_n(v)} \\ &= n_\infty(v)(1 - e^{-t/\tau_n(v)}) \end{aligned}$$

A step in the other direction gives:

$$\begin{aligned} n(t) &= n_\infty(0) + (n_\infty(v) - n_\infty(0))e^{-t/\tau_n(0)} \\ &= n_\infty(v)e^{-t/\tau_n(0)} \end{aligned}$$

# Gating variable raised to different powers



# Sodium conductance model

H&H realized that two different sub units were at work. Ended up with

$$g_{\text{Na}} = \bar{g}_{\text{Na}} m^3 h$$

Values for  $m_{\tau}$ ,  $m_{\infty}$ ,  $h_{\tau}$  and  $h_{\infty}$  obtained by fitting the solution to plots of  $g_{\text{Na}}$ .

# The Hodgkin-Huxley model

Introduces a third current, not time dependent:

$$C_m \frac{dv}{dt} = -\bar{g}_K n^4 (v - v_K) - \bar{g}_{Na} m^3 h (v - v_{Na}) - \bar{g}_L (v - v_L)$$

with

$$\frac{dg}{dt} = \alpha_g(v)(1 - g) - \beta_g(v)g, \quad g = m, h, n$$

Model based on voltage clamp measurement. How will it behave under normal conditions?

The model will predict the action potential.

# Analysis of the Hodgkin-Huxley model



## Qualitative analysis, 5.1.3

Would like to reduce the number of state variables to simplify analysis.

One way is to treat the slowest variables as constants. Of the three gating variables  $m$  has the fastest dynamics. (Controls the activation of the Na-current).

Reduced model:

$$C_m \frac{dv}{dt} = -\bar{g}_K n_0^4 (v - v_K) - \bar{g}_{Na} m^3 h_0 (v - v_{Na}) - \bar{g}_L (v - v_L)$$

## Equilibria in the reduced HH-model

The nullclines  $\frac{dv}{dt} = 0$  and  $\frac{dm}{dt} = 0$  form curves in the  $(v, m)$ -plane. Their intersections are the equilibria.

Initially three steady states  $v_r$ ,  $v_s$  and  $v_e$ .  $v_r$  and  $v_e$  are stable and  $v_s$  unstable.

As  $n_0$  and  $h_0$  changes, the  $\frac{dv}{dt} = 0$  line will shift.  $v_e$  will decrease, coincide with  $v_s$  and disappear.

$v_r$  will become the only stable equilibrium.

# Phase plot for the fast sub-system

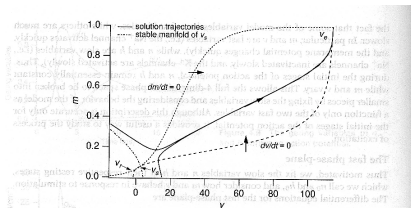


Figure 4.10 The Hodgkin-Huxley fast phase-plane, showing the nullclines  $dv/dt = 0$  and  $dm/dt = 0$  (with  $h_0 = 0.596$ ,  $n_0 = 0.3176$ ), two sample trajectories and the stable manifold of the saddle point  $v_s$ .

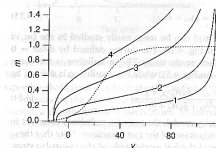


Figure 4.11 The Hodgkin-Huxley fast phase-plane as a function of the slow variables, showing the  $m$  nullcline (dashed) and the movement of the  $v$  nullcline (solid) and the disappearance of the steady states. For these curves, parameter values are (1)  $h_0 = 0.596$ ,  $n_0 = 0.3176$ ; (2)  $h_0 = 0.4$ ,  $n_0 = 0.5$ ; (3)  $h_0 = 0.2$ ,  $n_0 = 0.7$ ; and (4)  $h_0 = 0.1$ ,  $n_0 = 0.8$ .

Alternative reduction:

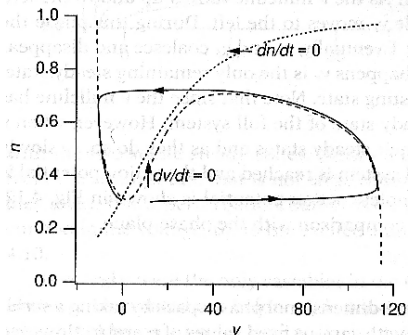
- $m$  is very fast, almost in equilibrium:  $m = m_\infty(v)$
- $h + n$  almost constant:  $h = 0.8 - n$

We then have

$$C_m \frac{dv}{dt} = -\bar{g}_K n^4 (v - v_K) - \bar{g}_{Na} m_\infty^3(v) \overbrace{(0.8 - n)}^h (v - v_{Na}) - \bar{g}_L (v - v_L)$$

Equilibria found by looking at the crossing of the nullclines  $\frac{dv}{dt} = 0$  and  $\frac{dn}{dt} = 0$  in the  $(v, n)$ -plane.

# Phase plot for the fast-slow reduced system

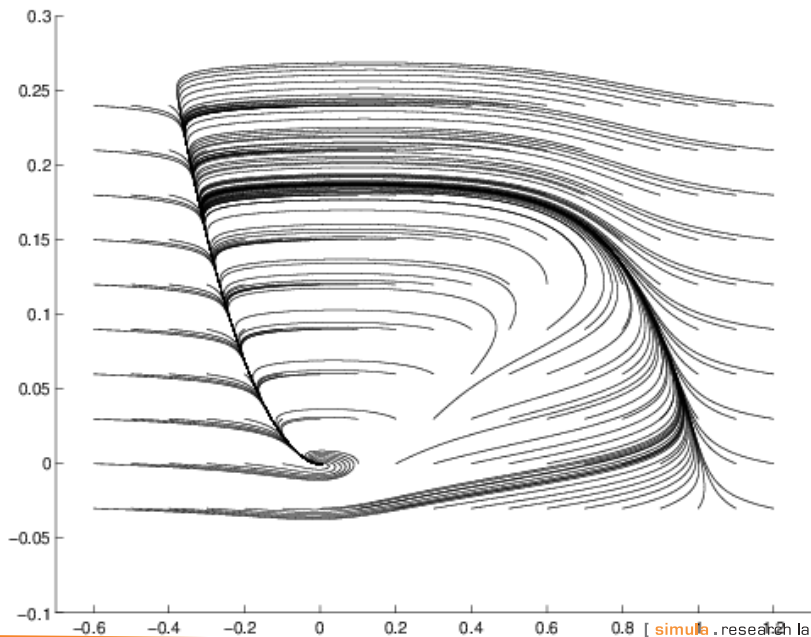


**Figure 4.13** Fast-slow phase-plane of the Hodgkin-Huxley model.

## Properties of the phase plot

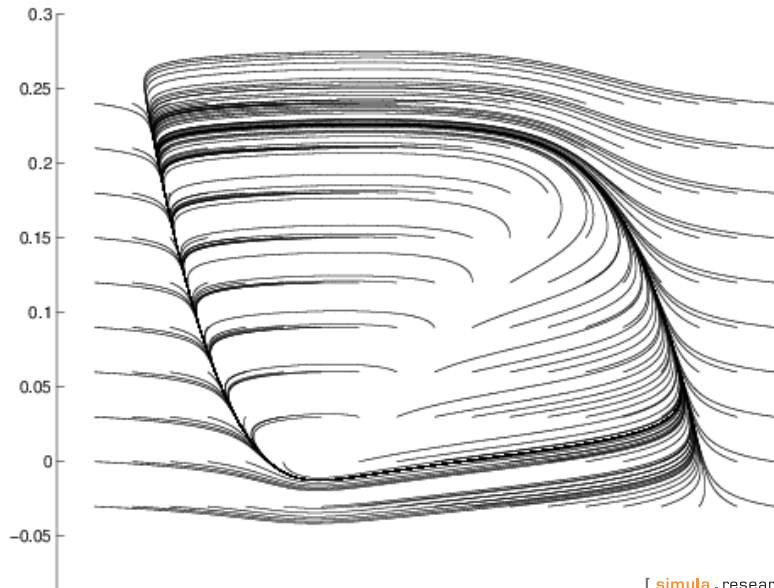
- $\frac{dv}{dt} = 0$  cubic form, with two stable and one unstable branch
- $\frac{dn}{dt} = 0$  sigmoid form
- One crossing with default parameters
- Trajectories horizontal due faster dynamic of  $v$
- Starting points to the left of the unstable branch converges to equilibrium without crossing the unstable branch
- Starting points to the right of the unstable branch crosses this branch, reaches the rightmost branch, follows this branch and the trajectory continues to rise until  $\frac{dn}{dt} = 0$  is crossed. The trajectory finally reaches the leftmost branch and follows it to the equilibrium points.

## Simulations with different initial conditions



## Modified model

The point  $(0,0)$  is no longer a stable equilibrium.





## Other models of the action potential

## The FitzHugh-Nagumo model, 5.2

Purpose:

Keep the qualitative behavior of the Hodgkin-Huxley system, but in a simplified form. Derivation based on an electrical circuit model.

On dimensionless form:

$$\begin{aligned}\epsilon \frac{dv}{dt} &= f(v) - w - w_0 \\ \frac{dw}{dt} &= v - \gamma w - v_0\end{aligned}$$

The variable  $w$  is called the recovery variable.

Typically  $f$  is chosen to be “cubic”, i.e. with three zeros,  $f(0) = f(\alpha) = f(1)$  and  $0 < \alpha < 1$ . Some choices:

$$f(v) = Av(v - \alpha)(1 - v)$$

$$f(v) = \begin{cases} -v, & v < \alpha \\ 1 - v, & v > \alpha \end{cases}$$

$$f(v) = \begin{cases} -v, & v < \alpha/2 \\ v - \alpha, & \alpha/2 < v < (1 + \alpha)/2 \\ 1 - v, & v > (1 + \alpha)/2 \end{cases}$$

# Cardiac cells

Excitable like neurons, display great variability

- SA node cells: Pace maker cells, controls the heart rate, self depolarizing
- AV node cells: Transmit signal from atria to ventricles with a delay
- Purkinje cells: Very high conductivity, propagate signal from AV out to the ventricles.
- Myocardial cells: Muscle cells (can contract)

These cells have different action potentials.

The HH-model was based on neurons. Other models necessary for cardiac cells.

# The Beeler-Reuter model

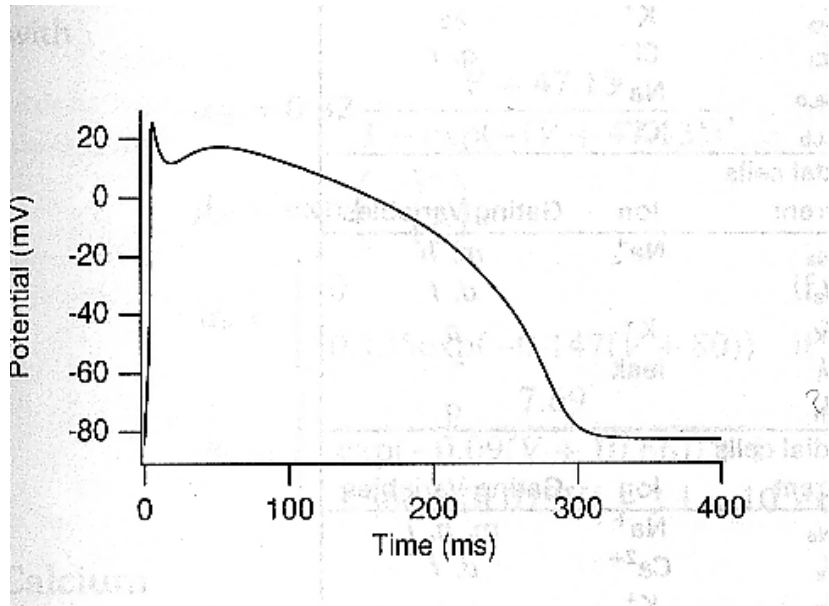
A model for ventricular cells, includes three currents, six gates and one ionic concentration.

$$-C_m \frac{dV}{dt} = I_{\text{Na}}(V, m, h, j) + I_{\text{K}}(V, x) + I_{\text{S}}(V, f, g, [\text{Ca}]_i)$$

Here  $m, h, j, x, f, g$  are gating variables and  $[\text{Ca}]_i$  is the intra cellular Calcium concentration

The action potential is much longer then for HH. Early repolarization (notch).

## Action potential produced by the Beeler-Reuter model



# Currents of the Beeler-Reuter model

Sodium current:

Third gating variable included to model the slow recovery (long refractory period). The model also include an ungated “leakage” current:

$$I_{\text{Na}} = (4m^3hj + 0.003)(V - 50)$$

Potassium:

One singled gated (with  $x$ ) and one ungated component:

$$I_{\text{K}} = f(v) + xg(v)$$

Calcium:

To gates,  $d$  activates,  $f$  inactivates:

$$I_S = 0.09fg(V - V_{Ca})$$

In addition the  $[Ca]_i$  is updated:

$$\frac{dc}{dt} = 0.07(1 - c) - I_S$$

where  $c = 10^7[Ca]_i$