

TEK9540 – Quantum computation and quantum information  
Problem set 3 – Solutions

**Exercise 2.40**

$$[X, Y] = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} - \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} = \begin{pmatrix} 2i & 0 \\ 0 & -2i \end{pmatrix} = 2iZ$$

$$[Y, Z] = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} - \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} = \begin{pmatrix} 0 & 2i \\ 2i & 0 \end{pmatrix} = 2iX$$

$$[Z, X] = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} - \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} = \begin{pmatrix} 0 & 2 \\ -2 & 0 \end{pmatrix} = 2iY$$

We can write

$$[\sigma_j, \sigma_k] = 2i \sum_{l=1}^3 \varepsilon_{jkl} \sigma_l$$

**Exercise 2.41**

Calculating

$$\{X, Y\} = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} + \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} = \begin{pmatrix} 0 & 0 \\ 0 & 0 \end{pmatrix}$$

$$\{X, Z\} = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} + \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} = \begin{pmatrix} 0 & 0 \\ 0 & 0 \end{pmatrix}$$

$$\{Y, Z\} = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} + \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} = \begin{pmatrix} 0 & 0 \\ 0 & 0 \end{pmatrix}$$

and

$$X^2 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

$$Y^2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

$$Z^2 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

Therefore we can write

$$\{\sigma_j, \sigma_k\} = 2\delta_{jk}I$$

### Exercise 2.42

We have that

$$AB = \frac{AB + AB}{2} + \frac{BA - BA}{2} = \frac{AB - BA}{2} + \frac{AB + BA}{2} = \frac{[A, B] + \{A, B\}}{2}$$

### Exercise 2.43

From the previous tasks we see that

$$\sigma_j \sigma_k = \frac{[\sigma_j, \sigma_k] + \{\sigma_j, \sigma_k\}}{2} = \frac{2\delta_{jk}I + 2i \sum_{l=1}^3 \varepsilon_{jkl} \sigma_l}{2} = \delta_{jk}I + i \sum_{l=1}^3 \varepsilon_{jkl} \sigma_l$$

### 2.51

$H$  is unitary if  $H^\dagger = H^{-1}$  i.e. if  $H^\dagger H = I$ . It is easy to see that  $H^\dagger = H$  so we have,

$$H^\dagger H = \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} = I$$

### 2.52

Since we have shown that  $H^\dagger H = I$  and  $H^\dagger = H$  we have  $H^2 = I$

### 2.53

The eigenvalues  $\lambda$  we find by solving the characteristic equation  $|H - \lambda I| = 0$ .

$$\left(-\lambda + \frac{1}{\sqrt{2}}\right)\left(-\lambda - \frac{1}{\sqrt{2}}\right) - \frac{1}{\sqrt{2}}\frac{1}{\sqrt{2}} = \lambda^2 - 1 = 0 \implies \lambda_1 = 1, \quad \lambda_2 = -1$$

We find the eigenvectors  $v$  by solving the equations  $Hv - \lambda v = 0$ . The unnormalized eigenvectors are

$$v_1 = \begin{pmatrix} 1 \\ \sqrt{2} - 1 \end{pmatrix} \quad v_2 = \begin{pmatrix} 1 \\ -\sqrt{2} - 1 \end{pmatrix}$$

### 2.54

Assume that  $A$  and  $B$  are commuting, Hermetian operators. Show that  $e^A e^B = e^{A+B}$ .

We know that two commuting Hermitian operators may be diagonalized in the same basis.

$$A = \sum_i a_i |i\rangle \langle i| \quad B = \sum_i b_i |i\rangle \langle i|.$$

Since  $A$  and  $B$  are Hermitian we have  $e^A = \sum_i e^{a_i} |i\rangle\langle i|$  and  $e^B = \sum_i e^{b_i} |i\rangle\langle i|$ . Then we get

$$\begin{aligned} e^A e^B &= \sum_i e^{a_i} |i\rangle\langle i| \sum_j e^{b_j} |j\rangle\langle j| = \sum_{i,j} e^{a_i} e^{b_j} |i\rangle\langle i| \underbrace{|j\rangle\langle j|}_{\delta_{ij}} \\ &= \sum_i e^{a_i+b_i} |i\rangle\langle i| = e^{A+B}. \end{aligned}$$

## 2.57

Assume that we initially have a system in a state  $|\psi\rangle$ . After a measurement with measurement operators  $\{L_l\}$  the system is in the state

$$|\phi_l\rangle = \frac{L_l|\psi\rangle}{\sqrt{\langle\psi|L_l^\dagger L_l|\psi\rangle}}$$

with probability

$$p_\phi(l) = \langle\psi|L_l^\dagger L_l|\psi\rangle.$$

If we then do a measurement with  $\{M_m\}$  the system will be in the state

$$\begin{aligned} |\xi_{ml}\rangle &= \frac{M_m|\phi_l\rangle}{\langle\phi_l|M_m^\dagger M_m|\phi_l\rangle} = \frac{M_m L_l|\psi\rangle}{\sqrt{\langle\psi|L_l^\dagger L_l|\psi\rangle} \sqrt{\frac{\langle\psi|L_l^\dagger}{\sqrt{\langle\psi|L_l^\dagger L_l|\psi\rangle}} M_m^\dagger M_m \frac{L_l|\psi\rangle}{\sqrt{\langle\psi|L_l^\dagger L_l|\psi\rangle}}}} \\ &= \frac{M_m L_l|\psi\rangle}{\sqrt{\langle\psi|L_l^\dagger M_m^\dagger M_m L_l|\psi\rangle}} = \frac{N_{lm}|\psi\rangle}{\sqrt{\langle\psi|N_{lm}^\dagger N_{lm}|\psi\rangle}} \end{aligned}$$

with probability

$$p_\xi(l, m) = \langle\phi_l|M_m^\dagger M_m|\phi_l\rangle p_\phi(l) = \langle\psi|N_{lm}^\dagger N_{lm}|\psi\rangle.$$

This shows that the two measurements  $\{M_m\}$  and  $\{L_l\}$  is equivalent to the measurement  $\{N_{lm}\}$ .

## 2.59

Expectation value:

$$\langle X \rangle = \langle 0|X|0\rangle = \begin{pmatrix} 1 & 0 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = 0$$

Standard deviation:

$$\sigma = \langle X^2 \rangle - \langle X \rangle^2 = \begin{pmatrix} 1 & 0 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} - 0 = 1$$

## 2.63

By singular value decomposition  $M_m = UDV$  and  $M_m^\dagger = V^\dagger D U^\dagger$  with  $D$  diagonal and non-negative and  $U$  and  $V$  unitary. We then have

$$\begin{aligned} \sqrt{E_m} &= \sqrt{M_m^\dagger M_m} = \sqrt{V^\dagger D U^\dagger U D V} = \sqrt{V^\dagger D^2 V} \\ &= V^\dagger D V = V^\dagger U^\dagger U D V = U_m^\dagger M_m \end{aligned}$$

with  $U_m = UV$ . Since  $U$  and  $V$  are unitary  $U_m$  is also unitary.

This shows that if we know a POVM-element, we know the corresponding measurement operator only up to a unitary transformation. Accordingly measurement operators only distinguished by a unitary transformation corresponds to the same POVM-element.

Note that the different POVM-elements of a measurement generally is related to the measurement operators by different unitary transformations.

## 2.66

We choose to use matrix notation, it is also possible to use outer product notation.

$$\frac{\langle 00| + \langle 11|}{\sqrt{2}} X \otimes Z \frac{|00\rangle + |11\rangle}{\sqrt{2}} = \frac{1}{2} \begin{pmatrix} 1 & 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \\ 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \end{pmatrix} \begin{pmatrix} 1 \\ 0 \\ 0 \\ 1 \end{pmatrix} = 0$$