

# Trial exam 2023 MATMEK-4270

1) a) Find  $u_N \in V_N$  such that

$$\textcircled{1} \quad (u - u_N, v)_{L^2} = 0, \quad \forall v \in V_N$$

The  $L^2$  inner product is here defined as

$$(a, b)_{L^2([-1, 1])} = \int_{-1}^1 a b dx,$$

where  $a(x)$  and  $b(x)$  are real functions.

1) b) Insert for  $v = P_i$  and  $u_N = \sum_{j=0}^N \hat{u}_j P_j$  in  $\textcircled{1}$

$$(u - \sum_{j=0}^N \hat{u}_j P_j, P_i) = 0, \quad \forall i = 0, 1, \dots, N$$

$$\sum_{j=0}^N (P_j, P_i) \hat{u}_j = (u, P_i)$$

7b) Use  $(P_j, P_i) = \frac{2}{2i+1} \delta_{ij}$  and solve for  $\hat{u}$

$$\hat{u}_i = \frac{2i+1}{2} (u, P_i), \quad i=0, 1, \dots, N$$

Insert for  $u = x^3$

$$\textcircled{2} \quad \hat{u}_i = \frac{2i+1}{2} (x^3, P_i), \quad i=0, 1, \dots, N$$

Since  $x^3$  is an odd function all even coefficients  $\hat{u}_{2i} = 0, i=0, 1, \dots$

Since  $x^3 \in P_3$   $N=3$ . So need to compute  $\hat{u}_1$  and  $\hat{u}_3$  from  $\textcircled{2}$ .

$$\hat{u}_1 = \frac{2 \cdot 1 + 1}{2} \int_{-1}^1 x^3 \cdot x \, dx = \frac{3}{2} \left[ \frac{1}{5} x^5 \right]_{-1}^1$$

$$\hat{u}_1 = \underline{\underline{\frac{3}{5}}}$$

For  $\hat{u}_3$  we need  $P_3$ , which can be computed from the recursion in Eq. 1.

$$7b) P_3 = \frac{1}{2}(5x^3 - 3x)$$

$$\hat{u}_3 = \frac{2 \cdot 3 + 1}{2} \int_{-1}^1 x^3 \cdot \frac{1}{2}(5x^3 - 3x) dx$$

$$= \frac{7}{4} \int_{-1}^1 (5x^6 - 3x^4) dx$$

$$= \frac{7}{4} \left[ \frac{5}{7} x^7 - \frac{3}{5} x^5 \right]_{-1}^1 = \frac{7}{4} \left( \frac{5}{7} - \frac{3}{5} \right) \cdot 2$$

$$= \frac{1}{2} \left( 5 - \frac{21}{5} \right) = \underline{\underline{\frac{2}{5}}}$$

$$\underline{\underline{\hat{u} = \left\{ 0, \frac{3}{5}, 0, \frac{2}{5} \right\}}}}$$

Note that

$$u_N = \sum_{j=0}^3 \hat{u}_j P_j = x^3$$

but the expansion coefficients are here required for full score, since these are requested.

7c) The approximation is exact

$$\underline{\underline{u_N = x^3}}$$

$\{x^n\}_{n=0}^N$  is a poor basis since the basis functions are not orthogonal in the  $L^2$  inner product space, or any other space. The basis functions are almost linearly dependent and the mass matrix

$$M = (m_{ij})_{i,j=0}^N, \quad m_{ij} = (x^j, x^i)$$

is ill-conditioned, leading to large round-off errors for numerical computations with finite precision.

2a) The Galerkin method is to find  $u_N \in V_N$ , such that

$$\textcircled{a} \quad (u_N'', v) + \alpha^2 (u_N, v) = (f, v)$$

for all  $v \in V_N$ . Alternatively, replace  $\textcircled{a}$  by  $\textcircled{b}$ :

$$\textcircled{b} \quad -(u_N', v') + \alpha^2 (u_N, v) = (f, v) - [u_N' v]_{x=-1}^{x=1}$$

where  $[u_N' v]_{-1}^1 = 0$ , since  $v(\pm 1) = 0$ .

2b) Insert for  $v = \psi_i$  and  $u_N = \sum_{j=0}^N \hat{u}_j \psi_j$  in either  $\textcircled{a}$  or  $\textcircled{b}$ . Using  $\textcircled{a}$ :

$$\sum_{j=0}^N (\psi_j'', \psi_i) \hat{u}_j + \sum_{j=0}^N \alpha^2 (\psi_j, \psi_i) \hat{u}_j = (f, \psi_i)$$

for  $i = 0, 1, \dots, N$ .

Specific bases that can be used:

$$\left\{ P_i - P_{i+2} \right\}_{i=0}^N$$

$$\left\{ \sin \left( (i+1) \pi \frac{x+1}{2} \right) \right\}_{i=0}^N$$

Both satisfy the boundary conditions.

2b) Note that this problem also can be solved with the finite element method. However, since it is given that  $\psi_j(\pm 1) = 0$ , this means that you cannot use Lagrange polynomials, since  $l_0(-1) = 1$  and  $l_1(1) = 1$ .

2c) Define

$$u = \tilde{u} + B,$$

where  $B(-1) = u(-1) = a$  and

$$B(1) = u(1) = b:$$

$$B = \frac{a}{2}(1-x) + \frac{b}{2}(1+x)$$

Insert  $u = \tilde{u} + B$  into Eq. (2):

$$(\tilde{u} + B)'' + \alpha^2(\tilde{u} + B) = f$$

$$\tilde{u}'' + \alpha^2\tilde{u} = f - \alpha^2 B$$



2c) The Galerkin problem becomes:  
Find  $\tilde{u}_N \in V_N$  such that

$$(\tilde{u}_N'', v) + \alpha^2 (\tilde{u}_N, v) = (f - \alpha^2 B, v)$$

for all  $v \in V_N$ .

$$\underline{\underline{\text{Set } u_N = \tilde{u}_N + B}}$$

2d) The basis functions do not satisfy  $\psi_j(\pm 1) = 0$ , or  $\psi_j'(\pm 1) = 0$ .

Hence we need to set the boundary conditions weakly using integration by parts.

From 2a) (b):

$$-(u_N', v') + \alpha^2 (u_N, v) = (f, v) - [u_N' v]_{x=-1}^{x=1}$$

Since  $u_N'(\pm 1) = 0$ , we can solve

$$-(u_N', v') + \alpha^2 (u_N, v) = (f, v)$$

2d) Insert for  $v = P_i$  and  $u_N = \sum_{j=0}^N \hat{u}_j P_j$   
to obtain the equation

$$-\sum_{j=0}^N (P_j', P_i') \hat{u}_j + \alpha^2 \sum_{j=0}^N (P_j, P_i) \hat{u}_j = (f, P_i)$$

for  $i = 0, 1, \dots, N$ .

Note that it is not required to  
derive the linear algebra problem.



3a) The diffusion equation is discretized using  $N$  equal intervals of size  $\Delta x = \frac{1}{N}$ , and  $N_t$  time steps of length  $\Delta t = \frac{T}{N_t}$ .

A mesh function is defined as

$$u_j^n = u(x_j, t_n) = u(j\Delta x, n\Delta t)$$

Initialize  $u^0 = \{u_j^0\}_{j=0}^N$  as

$$u_j^0 = I(x_j), \quad j = 0, 1, \dots, N$$

Discretization:

$$\frac{u_j^{n+1} - u_j^n}{\Delta t} = c \frac{u_{j+1}^{n+1} - 2u_j^{n+1} + u_{j-1}^{n+1} + u_{j+1}^n - 2u_j^n + u_{j-1}^n}{2\Delta x^2}$$

Or in matrix form:

$$\frac{u^{n+1} - u^n}{\Delta t} = \frac{c}{2} D^{(2)} (u^{n+1} + u^n) \quad *$$

3a) where  $D^{(2)}$  is the central, second order accurate, second differentiation matrix.  $D^{(2)} \in \mathbb{R}^{N \times N}$

We get from \*

$$\left(I - \frac{c\Delta t}{2} D^{(2)}\right) u^{n+1} = \left(I + \frac{c\Delta t}{2} D^{(2)}\right) u^n$$

where  $I \in \mathbb{R}^{N \times N}$  is the identity matrix.

Writing further

$$A u^{n+1} = B u^n$$

where  $A = I - \frac{c\Delta t}{2} D^{(2)}$ ,  $B = I + \frac{c\Delta t}{2} D^{(2)}$ ,

we get

$$u^{n+1} = A^{-1} B u^n, \quad n = 0, 1, \dots, N_t - 1.$$

The complete algorithm becomes:

Initialize  $u^0$

For  $n = 0, 1, \dots, N_t - 1$ , find

$$u^{n+1} = A^{-1} B u^n$$

Set  $u_0^{n+1} = 0$  and  $u_N^{n+1} = 0$

3b) Define the approximation

$$u_N^n = \sum_{j=0}^N \hat{u}_j^n \psi_j(x) \approx u(x, n\Delta t)$$

and the function space  $V_N = \text{span} \{ \psi_j \}_{j=0}^N$ .

Let  $\psi_j(0) = \psi_j(1) = 0$  for all  $j \geq 0$ .

Initialize  $u_N^0$  by projection: Find  $u_N^0 \in V_N$ , such that

$$(I(x) - u_N^0, v) = 0, \quad \forall v \in V_N$$

For remaining time steps:

For  $n = 0, 1, \dots, N_t - 1$  find  $u_N^{n+1} \in V_N$  such that

$$(u_N^{n+1} - u_N^n, v) = \frac{c\Delta t}{2} \left( \frac{\partial^2}{\partial x^2} (u_N^{n+1} + u_N^n), v \right)$$

for all  $v \in V_N$ .

3c) We have the equation

$$\left(I - \frac{c\Delta t}{2} D^{(2)}\right) u^{n+1} = \left(I + \frac{c\Delta t}{2} D^{(2)}\right) u^n \quad (*)$$

Assume that the solution can be written as (ansatz)

$$u^{n+1} = g^{n+1} u^0,$$

where  $g$  is the amplification factor, independent of  $x$  and  $t$ .

Insert the ansatz into (\*) and divide by  $g^n$ :

$$g\left(I - \frac{c\Delta t}{2} D^{(2)}\right) u^0 = \left(I + \frac{c\Delta t}{2} D^{(2)}\right) u^0$$

We can now use that

$$D^{(2)} u^0 = \lambda u^0,$$

where  $\lambda$  are the eigenvalues of  $D^{(2)}$ .

We get

$$g\left(I - \frac{c\Delta t}{2} \lambda I\right) u^0 = \left(I + \frac{c\Delta t}{2} \lambda I\right) u^0$$

3c) Simplify by eliminating  $u^0$

$$g\left(1 - \frac{c\Delta t}{2}\lambda\right) = \left(1 + \frac{c\Delta t}{2}\lambda\right),$$

such that

$$g = \frac{1 + \frac{c\Delta t}{2}\lambda}{1 - \frac{c\Delta t}{2}\lambda}$$

It is reported that  $\lambda < 0$  and real, and also that  $c > 0$ . Hence

$|g| \leq 1$  for all  $\Delta t$ , and thus unconditionally stable.