

SUGGESTED SOLUTIONS TO
PARTS OF MANDATORY
ASSIGNMENT: MAT 4770

1a

We only do the argument with Itô's Formula.
The other approach has been done in previous
exercises.

Choose a function $g(t, x) = e^{\alpha t} \cdot x$, and
use $x = \mathcal{X}(t)$, where $d\mathcal{X}(t) = -\alpha \mathcal{X}(t) dt + \sigma dB(t)$:

$$dg(t, \mathcal{X}(t)) = \frac{\partial g}{\partial t}(t, \mathcal{X}(t)) dt + \frac{\partial g}{\partial x}(t, \mathcal{X}(t)) d\mathcal{X}(t) \\ + \frac{1}{2} \frac{\partial^2 g}{\partial x^2}(t, \mathcal{X}(t)) (d\mathcal{X}(t))^2$$

$$\left(\frac{\partial g}{\partial t} = \alpha e^{\alpha t} x = \alpha g, \quad \frac{\partial g}{\partial x} = e^{\alpha t}, \quad \frac{\partial^2 g}{\partial x^2} = 0 \right)$$

$$= \alpha g(t, \mathcal{X}(t)) dt + e^{\alpha t} d\mathcal{X}(t) \\ = \underbrace{\alpha g(t, \mathcal{X}(t)) dt + e^{\alpha t} (-\alpha \mathcal{X}(t)) dt}_{= 0}$$

$$+ e^{\alpha t} \sigma dB(t)$$

$$= \sigma e^{\alpha t} dB(t)$$

Integrating both sides from 0 to t yields

$$e^{\alpha t} \bar{X}(t) - e^{\alpha \cdot 0} \bar{X}(0) = \int_0^t \sigma e^{\alpha v} dB(v)$$

$$\Rightarrow \bar{X}(t) = e^{-\alpha t} \bar{X}(0) + \sigma e^{-\alpha t} \int_0^t e^{\alpha v} dB(v)$$

For $s \geq t$, we have

$$\begin{aligned} \bar{X}(s) &= e^{-\alpha s} \bar{X}(0) + \sigma e^{-\alpha s} \int_0^s e^{\alpha v} dB(v) \\ &= e^{-\alpha(s-t)} \cdot e^{-\alpha t} \bar{X}(0) + \sigma e^{-\alpha(s-t)} e^{-\alpha t} \int_0^s e^{\alpha v} dB(v) \end{aligned}$$

$$\begin{aligned} \int_0^s e^{\alpha v} dB(v) &= \int_0^t e^{\alpha v} dB(v) + \int_t^s e^{\alpha v} dB(v) \\ &= e^{-\alpha(s-t)} \left(e^{-\alpha t} \bar{X}(0) + \sigma e^{-\alpha t} \int_0^t e^{\alpha v} dB(v) \right) \\ &\quad + \sigma e^{-\alpha(s-t)} \int_t^s e^{\alpha v} dB(v) \end{aligned}$$

$$= e^{-\alpha(s-t)} \bar{X}(t) + \sigma e^{-\alpha s} \int_t^s e^{2\alpha v} dB(v)$$

Note that $\int_t^s e^{2\alpha v} dB(v)$ is an integral of a deterministic function $e^{2\alpha v}$ with respect to Brownian motion. We know that such integrals are Gaussian random variables. We also know that

$$E\left[\int_t^s e^{2\alpha v} dB(v)\right] = 0$$

$$E\left[\left(\int_t^s e^{2\alpha v} dB(v)\right)^2\right] = \int_t^s e^{2\alpha v} dv$$

Hence $\int_t^s e^{2\alpha v} dB(v) \sim N\left(0, \int_t^s e^{2\alpha v} dv\right)$

$$\Rightarrow \bar{X}(s) | \bar{X}(t) \sim N\left(e^{-\alpha(s-t)} \bar{X}(t), \sigma^2 e^{-2\alpha s} \int_t^s e^{2\alpha v} dv\right)$$

from the rules of Gaussian random variables

$$\left(Z \sim N(0, b^2) \Rightarrow \mu + cZ \sim N(\mu, c^2 b^2) \right)$$

Since $\bar{X}(t)$ by definition is given by

$\int_0^t e^{\alpha v} dB(v)$, it depends on Brownian motion $B(v)$, for $0 \leq v \leq t$, i.e., on B only up to time t . As $\int_t^s e^{\alpha v} dB(v) \sim \sum e^{\alpha v_i} \Delta B(v_i)$

this integral depends on the increments of B from t to s , i.e., increments after time t . Brownian motion has independent increments by definition, and hence $\int_t^s e^{\alpha v} dB(v)$ is independent of $X(t)$ for $s \geq t$.

1b

Choose the function $g(t, x) = e^x$, and $x = X(t)$, $dX(t) = -\alpha X(t)dt + \sigma dB(t)$. Using Itô's Formula we find:

$$\begin{aligned} dS(t) &= dg(t, X(t)) \\ &= \frac{\partial g}{\partial t}(t, X(t)) + \frac{\partial g}{\partial x}(t, X(t)) dX(t) \\ &\quad + \frac{1}{2} \frac{\partial^2 g}{\partial x^2}(t, X(t)) (dX(t))^2 \end{aligned}$$

$$\frac{\partial g}{\partial t} = 0, \quad \frac{\partial g}{\partial x} = e^x = g, \quad \frac{\partial^2 g}{\partial x^2} = g$$

$$(d\bar{X}(t))^2 = (-\alpha \bar{X}(t))^2 (dt)^2 + 2(-\alpha \bar{X}(t)) dt \cdot \sigma dB(t) + \sigma^2 (dB(t))^2$$

$$= \sigma^2 dt, \text{ using } (dt)^2 = 0$$

$$dt \cdot dB(t) = 0$$

$$(dB(t))^2 = dt$$

RULES.

$$= \underbrace{g(t, \bar{X}(t))}_{= S(t)} \left((-\alpha \bar{X}(t)) dt + \sigma dB(t) \right)$$

$$+ \frac{1}{2} \underbrace{g(t, \bar{X}(t))}_{= S(t)} \sigma^2 dt$$

$$= \left(\frac{1}{2} \sigma^2 - \alpha \underbrace{\bar{X}(t)}_{\ln S(t)} \right) S(t) dt + \sigma S(t) dB(t)$$

$$\underline{\underline{dS(t) = \left(\frac{1}{2} \sigma^2 - \alpha \ln S(t) \right) S(t) dt + \sigma S(t) dB(t)}}$$

Since $\bar{X}(t)$ is Gaussian, $S(t) = e^{\bar{X}(t)}$ is lognormal by definition.

Recall that if $Z \sim N(a, b^2)$, then

$$\underline{E[e^Z]} = \underline{e^{a + \frac{1}{2}b^2}}$$

Moreover $2Z \sim N(2a, 4b^2)$, so

$$E[(e^Z)^2] = E[e^{2Z}] = e^{2a + \frac{1}{2}4b^2} = e^{2a + 2b^2}$$

Therefore ;

$$\begin{aligned} \underline{\text{Var}(e^Z)} &= E[e^{2Z}] - E[e^Z]^2 \\ &= e^{2a + 2b^2} - e^{2a + b^2} = \underline{e^{2a + b^2} (e^{b^2} - 1)} \end{aligned}$$

From 1a, we know that $X(t) \sim N(a, b^2)$
with $a = e^{-\alpha t} X(0)$ and $b^2 = \sigma^2 e^{-2\alpha t} \int_0^t e^{2\alpha v} dv$

Inserting into above, yields $E[S(t)]$ and
 $\text{Var}(S(t))$!

1c.

Giřsenov's theorem says that W is a
Brownian motion on (any) given interval

$0 \leq t \leq \tau$ with respect to the probability Q given by

$$\frac{dQ}{dP} = \exp\left(-\theta B(\tau) - \frac{1}{2}\theta^2\right)$$

With respect to Q , \bar{X} has the form

$$\begin{aligned} \bar{X}(s) &= e^{-\alpha(s-t)} \bar{X}(t) + \sigma e^{-\alpha s} \int_t^s e^{\alpha v} (dW(v) + \theta dv) \\ &= e^{-\alpha(s-t)} \bar{X}(t) + \sigma e^{-\alpha s} \theta \int_t^s e^{\alpha v} dv + \sigma e^{-\alpha s} \int_t^s e^{\alpha v} dW(v) \\ &= e^{-\alpha(s-t)} \bar{X}(t) + \frac{\sigma\theta}{\alpha} (1 - e^{-\alpha(s-t)}) \\ &\quad + \sigma e^{-\alpha s} \int_t^s e^{\alpha v} dW(v) \end{aligned}$$

Hence, (typo in arrangement, $T \rightarrow \tau$!)

$$f(t, \tau) = E_Q \left[\exp(\bar{X}(\tau)) \mid \mathcal{F}_t \right]$$

$$= E_Q \left[\exp \left(e^{-\alpha(\tau-t)} \bar{X}(t) + \frac{\sigma\theta}{\alpha} (1 - e^{-\alpha(\tau-t)}) + \sigma e^{-\alpha\tau} \int_t^\tau e^{\alpha v} dW(v) \right) \mid \mathcal{F}_t \right]$$

\uparrow \mathcal{F}_t -measurable \uparrow deterministic \uparrow \mathcal{F}_t -independent

$$= \exp\left(e^{-\alpha(\tau-t)} \mathcal{I}(t) + \frac{\sigma\theta}{\alpha}(1-e^{-\alpha(\tau-t)})\right) \\ \cdot E_q\left[\exp\left(\sigma e^{-\alpha\tau} \int_t^\tau e^{\alpha\nu} dW(\nu)\right)\right]$$

$$\sim N\left(0, \sigma^2 e^{-2\alpha\tau} \int_t^\tau e^{2\alpha\nu} d\nu\right)$$

$$(*) \quad \underline{\underline{= \exp\left(e^{-\alpha(\tau-t)} \mathcal{I}(t) + \frac{\sigma\theta}{\alpha}(1-e^{-\alpha(\tau-t)}) + \frac{1}{2} \sigma^2 e^{-2\alpha\tau} \int_t^\tau e^{2\alpha\nu} d\nu\right)}} \\ = \frac{1}{2\alpha} (e^{2\alpha\tau} - e^{2\alpha t})$$

Since $\mathcal{I}(t) = \ln S(t)$, we get

$$\exp\left(e^{-\alpha(\tau-t)} \mathcal{I}(t)\right) = \exp\left(e^{-\alpha(\tau-t)} \ln S(t)\right) \\ = \left(\exp(\ln S(t))\right)^{e^{-\alpha(\tau-t)}} \\ = S(t)^{e^{-\alpha(\tau-t)}}$$

Thus

$$\underline{\underline{f(t, \tau) = S(t)^{e^{-\alpha(\tau-t)}} \cdot \exp\left(\frac{\sigma\theta}{\alpha}(1-e^{-\alpha(\tau-t)}) + \frac{\sigma^2}{4\alpha}(1-e^{-2\alpha(\tau-t)})\right)}}$$

By definition of $f(t, \tau)$ as a conditional expectation, the process $t \mapsto f(t, \tau)$ for $t \leq \tau$ is a martingale under \mathbb{Q} . As a martingale under \mathbb{Q} , its dynamics can only depend on $dW(t)$ (the \mathbb{Q} -Brownian motion), and not on dt .

Going back to the expression of $f(t, \tau)$, we can write it as

$$f(t, \tau) = \exp\left(e^{-\alpha(\tau-t)} \bar{X}(t) + h(t, \tau)\right)$$

for a function $h(t, \tau)$ that we can read off from the expression (*)

Consider now Itô's formula used on

$$g(t, x) = \exp\left(e^{-\alpha(\tau-t)} x + h(t, \tau)\right) \text{ and}$$

$$x = \bar{X}(t) \text{ with } d\bar{X}(t) = -\alpha \bar{X}(t) dt + \sigma dW(t) - \sigma \theta dt$$

From Itô we see

$$df(t, \tau) = \frac{\partial g}{\partial t} dt + \frac{\partial g}{\partial x} d\tilde{x}(t) + \frac{1}{2} \frac{\partial^2 g}{\partial x^2} (d\tilde{x}(t))^2$$

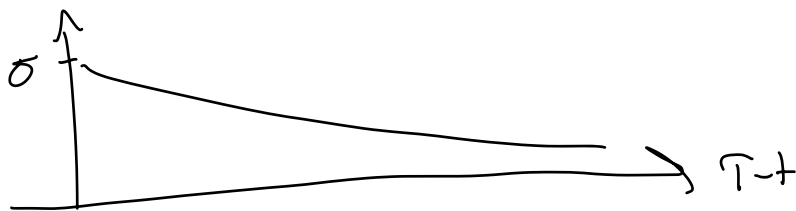
\uparrow
dt
 \uparrow
dt & dW
 \uparrow
dt

But all dt-terms must cancel since we know that we must end up with only a dW-term. Hence, we only need to consider $\frac{\partial g}{\partial x}$ and the dW-part of $d\tilde{x}$!

$$\begin{aligned} \underline{\underline{df(t, \tau)}} &= \frac{\partial g}{\partial x}(t, \tilde{x}(t)) \sigma dW(t) & \frac{\partial g}{\partial x} &= e^{-\alpha(\tau-t)} \cdot g \\ &= e^{-\alpha(\tau-t)} g(t, \tilde{x}(t)) \sigma dW(t) \\ &= \underline{\underline{f(t, \tau) \sigma e^{-\alpha(\tau-t)} dW(t)}} \end{aligned}$$

When $\tau - t \downarrow 0$, then $\sigma e^{-\alpha(\tau-t)} \rightarrow \sigma$

$\tau - t \rightarrow \infty$, then $\sigma e^{-\alpha(\tau-t)} \rightarrow 0$



1d

From above, $f(t, \tau)$ is a Geometric
Brownian motion with time-dependent vol.

We guess that

$$(*) \quad f(t, \tau) = f(b, \tau) \exp\left(\int_0^t \sigma e^{-\alpha(\tau-s)} dW(s) - \frac{1}{2} \int_0^t \sigma^2 e^{-2\alpha(\tau-s)} ds\right)$$

By letting $dZ(t) = \sigma e^{-\alpha(\tau-t)} dW(t)$, and

considering the function

$$g(t, y) = f(b, \tau) \exp\left(y - \frac{1}{2} \int_0^t \sigma^2 e^{-2\alpha(\tau-s)} ds\right)$$

we can show that (*) is indeed the
solution.

Hence, $t = T \leq \tau$

$$f(T, \tau) = f(b, \tau) \exp\left(\int_0^T \sigma e^{-\alpha(\tau-s)} dW(s) - \frac{1}{2} \int_0^T \sigma^2 e^{-2\alpha(\tau-s)} ds\right)$$

Notice that

$$\int_0^T \sigma e^{-\alpha(\tau-s)} dW(s) \sim \mathcal{N}\left(0, \int_0^T \sigma^2 e^{-2\alpha(\tau-s)} ds\right)$$

$$\text{Denote } \Sigma^2 := \int_0^T \sigma^2 e^{-2\alpha(\tau-t)} dt = \frac{\sigma^2}{2\alpha} \begin{pmatrix} -2\alpha(T-T) & -2\alpha T \\ e & -e \end{pmatrix}$$

Then, in distribution, we find

$$f(T, \tau) \stackrel{d}{=} f(0, \tau) \exp\left(\Sigma z - \frac{1}{2} \Sigma^2\right)$$

where $z \sim N(0, 1)$ under Q .

We calculate $E_Q[\max(f(T, \tau) - K, 0)]$.

$$\begin{aligned} E_Q[\max(f(T, \tau) - K, 0)] &= E_Q\left[\max\left(f(0, \tau) e^{\Sigma z - \frac{1}{2} \Sigma^2} - K, 0\right)\right] \\ &= \frac{1}{\sqrt{2\pi}} \int_{\mathbb{R}} \max\left(f(0, \tau) e^{\Sigma z - \frac{1}{2} \Sigma^2} - K, 0\right) e^{-z^2/2} dz \end{aligned}$$

$$f(0, \tau) e^{\Sigma z - \frac{1}{2} \Sigma^2} > K \Leftrightarrow \Sigma z - \frac{1}{2} \Sigma^2 > \ln \frac{K}{f(0, \tau)}$$

$$\Leftrightarrow z > \frac{\ln \frac{K}{f(0, \tau)} + \frac{1}{2} \Sigma^2}{\Sigma} =: -d$$

$$= \frac{1}{\sqrt{2\pi}} \int_{-d}^{\infty} \left(f(0, \tau) e^{\Sigma z - \frac{1}{2} \Sigma^2} - K\right) e^{-z^2/2} dz$$

$$= \frac{1}{\sqrt{2\pi}} f(b, \tau) e^{-\frac{1}{2}\bar{z}^2} \int_{-d}^{\infty} e^{\bar{z}z - \frac{1}{2}z^2} dz$$

$$- K \frac{1}{\sqrt{2\pi}} \int_{-d}^{\infty} e^{-\frac{z^2}{2}} dz$$

$$= -\frac{1}{2}(z^2 - 2\bar{z}z + \bar{z}^2 - \bar{z}^2)$$

$$= -\frac{1}{2}(z^2 - 2\bar{z}z + \bar{z}^2) + \frac{1}{2}\bar{z}^2$$

$$= -\frac{1}{2}(z - \bar{z})^2 + \frac{1}{2}\bar{z}^2$$

$$= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^d e^{-\frac{z^2}{2}} dz = \Phi(d)$$

Φ cumulative normal distr. function.

$$= f(b, \tau) \frac{1}{\sqrt{2\pi}} \int_{-d}^{\infty} e^{-\frac{1}{2}(z - \bar{z})^2} dz - K \Phi(d)$$

change of variable $u := z - \bar{z}$

$$= f(b, \tau) \frac{1}{\sqrt{2\pi}} \int_{-d - \bar{z}}^{\infty} e^{-\frac{1}{2}u^2} du - K \Phi(d)$$

$$= f(b, \tau) \Phi(d + \bar{z}) - K \Phi(d)$$

Hence

$$C(t, \tau) = e^{-r\tau} f(b, \tau) \Phi(d + \bar{z}) - e^{-r\tau} K \Phi(d)$$

$$\text{where } d = \frac{\ln \frac{f(b, \tau)}{K} - \frac{1}{2} \Sigma^2}{\Sigma}$$

$$\text{and } \Sigma^2 = \frac{\sigma^2}{2\alpha} \begin{pmatrix} -2\alpha(\tau - T) & -2\alpha\tau \\ e & -e \end{pmatrix}$$

Black76, almost as famous as Black & Scholes.
 here, Black76 with time-dependent vol!

2a

Comment on the initial values of \bar{X} .

Since $\bar{X}_1(t) = T(t) - A(t)$, you find that

$\bar{X}_1(0) = T(0) - A(0)$. Hence, you read off

today's temperature $\hat{T}(0)$, subtract $A(0)$, and

you have the initial value

$$\bar{X}_1(0) = \hat{T}(0) - A(0).$$

Since

$\hat{X}_2'(t) = T'(t) - \Lambda'(t)$, we approximate $T'(t)$ by its numerical derivative, which is given by (with daily time steps)

$$T'(t) \approx \frac{T(t) - T(t-1)}{1} = T(t) - T(t-1)$$

$$\text{So, } \hat{T}'(t) = \hat{T}(t) - \hat{T}(t-1)$$

\uparrow today \uparrow yesterday.

Hence,

$$\hat{X}_2(t) = \hat{T}(t) - \hat{T}(t-1) - \underbrace{\Lambda'(t)}$$

this you calculate from your seasonality function.

Finally

$$\hat{X}_3(t) = T''(t) - \Lambda''(t).$$

$$T''(t) \approx \frac{T(t) - 2T(t-1) + T(t-2)}{1^2}$$

$$= T(t) - 2T(t-1) + T(t-2)$$

Therefore

$$T''(t) \approx \underbrace{T(t)}_{\text{today}} - 2\underbrace{T(t-1)}_{\text{yesterday}} + \underbrace{T(t-2)}_{\text{2 days ago}}$$

$$\vec{x}_3(t) = T(t) - 2T(t-1) + T(t-2) - \Lambda''(t)$$

By reading off the temperature the last 3 days (today, yesterday and day before yesterday), we can initialize $\vec{x}(t) = (x_1(t), x_2(t), x_3(t))$.