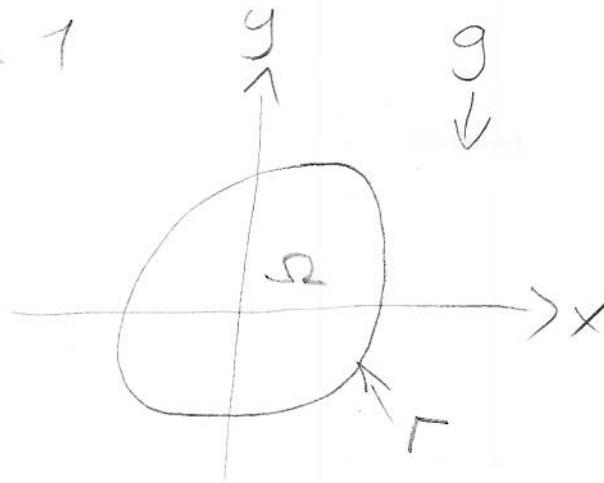


Problem 1



a) $u \cdot \nabla u = -\frac{1}{\rho} \nabla p + \nu \nabla^2 u - gk$

$$\nabla \cdot u = 0$$

$$u = 0 \text{ on } \Gamma$$

b) $u = u_z k = (0, 0, u_z)$

Momentum eq. z-dir:

$$u \cdot \nabla u_z = -\frac{1}{\rho} \frac{\partial p}{\partial z} + \nu \nabla^2 u_z$$

Uniform in z

$$\frac{1}{\rho} \frac{\partial p}{\partial z} = \nu \left(\frac{\partial^2 u_z}{\partial x^2} + \frac{\partial^2 u_z}{\partial y^2} \right)$$

Momentum eq. x-direction:

$$0 = \frac{\partial p}{\partial x}$$

Momentum eq. y-direction:

$$0 = -\frac{\partial p}{\partial y} - \rho g$$

=> $P = -\rho g y + P_0(z)$, P_0 constant

$$\frac{\partial p}{\partial z} = \frac{\partial P_0}{\partial z}$$

$$\frac{\partial P_0}{\partial z} = \mu \left(\frac{\partial^2 u_z}{\partial x^2} + \frac{\partial^2 u_z}{\partial y^2} \right)$$

Left hand side function of z only and
right hand side function of x and y .

$$\Rightarrow \frac{\partial P_0}{\partial z} = \text{constant}$$

We have

$$\frac{\partial^2 u_z}{\partial x^2} + \frac{\partial^2 u_z}{\partial y^2} = -\beta$$

$$\Rightarrow \frac{1}{\mu} \frac{\partial P_0}{\partial z} = \beta$$

$$c) \tau_w = \mu (\nabla u + \nabla u^T) \cdot n \quad (= \mu \nabla u \cdot n)$$

$$\nabla u = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ \frac{\partial u_z}{\partial x} & \frac{\partial u_z}{\partial y} & \cancel{\frac{\partial u_z}{\partial z}} \end{pmatrix}, \quad \nabla u^T = \begin{pmatrix} 0 & 0 & \frac{\partial u_z}{\partial x} \\ 0 & 0 & \frac{\partial u_z}{\partial y} \\ 0 & 0 & 0 \end{pmatrix}$$

$$n = n_x \hat{i} + n_y \hat{j} = (n_x, n_y, 0)$$

$$\tau_w = \mu \begin{pmatrix} 0 & 0 & \frac{\partial u_z}{\partial x} \\ 0 & 0 & \frac{\partial u_z}{\partial y} \\ \frac{\partial u_z}{\partial x} & \frac{\partial u_z}{\partial y} & 0 \end{pmatrix} \cdot \begin{pmatrix} n_x \\ n_y \\ 0 \end{pmatrix} = \underline{\underline{\left(\mu \frac{\partial u_z}{\partial x} n_x + \mu \frac{\partial u_z}{\partial y} n_y \right)}}$$

Show:

(3)

$$\Phi = \mu \left(\left(\frac{\partial u_2}{\partial x} \right)^2 + \left(\frac{\partial u_2}{\partial y} \right)^2 \right)$$

Definition:

$$\Phi = \mu \nabla u : (\nabla u + \nabla u^T)$$

$$= \mu \frac{\partial u_i}{\partial x_j} \left(\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right)$$

$$= \mu \cdot \begin{pmatrix} 0 & 0 & \frac{\partial u_2}{\partial x} \\ 0 & 0 & \frac{\partial u_2}{\partial y} \\ \frac{\partial u_2}{\partial x} & \frac{\partial u_2}{\partial y} & 0 \end{pmatrix} \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ \frac{\partial u_2}{\partial x} & \frac{\partial u_2}{\partial y} & 0 \end{pmatrix}$$

$$= \mu \left(\left(\frac{\partial u_2}{\partial x} \right)^2 + \left(\frac{\partial u_2}{\partial y} \right)^2 \right)$$

d)

$$F(x, y) = \frac{x^2}{a^2} + \frac{y^2}{b^2} - 1$$

$$\nabla^2 F = 2 \left(\frac{1}{a^2} + \frac{1}{b^2} \right) = \text{constant}$$

and $F(x, y) = 0$ on boundary Γ .

\Rightarrow F satisfies both equation and boundary condition

(4)

$u_z = F \cdot C$, where C is constant

We have

$$\nabla^2 u_z = -\beta$$

$$2C \left(\frac{1}{a^2} + \frac{1}{b^2} \right) = -\beta$$

$$C = -\frac{\beta}{2 \left(\frac{1}{a^2} + \frac{1}{b^2} \right)} = -\frac{\beta}{2} \frac{a^2 b^2}{b^2 + a^2}$$

$$u_z = -\frac{\beta}{2} \frac{a^2 b^2}{a^2 + b^2} \left(\frac{x^2}{a^2} + \frac{y^2}{b^2} - 1 \right)$$

e)

$$u_z = -\frac{\beta}{2} \frac{a^{2x}}{2a^2} \left(\frac{x^2 + y^2}{a^2} - 1 \right)$$

$$= -\frac{\beta}{2} a^2 \left(\frac{x^2 + y^2}{a^2} - 1 \right)$$

$$= -\frac{\beta}{4} (x^2 + y^2 - a^2)$$

Cylindrical coordinates

$$u_z = -\frac{\beta}{4} (r^2 - a^2)$$

1e)

$$Q = \int_{\Sigma} u_z dA$$

$$= \int_0^{2\pi} \int_0^a -\frac{3}{4} (r^2 - a^2) r d\theta dr$$

$$= -\frac{3}{4} 2\pi \int_0^a (r^3 - a^2 r) dr$$

$$= -\frac{3\pi}{2} \left(\frac{1}{4} r^4 - \frac{a^2}{2} r^2 \right) \Big|_0^a$$

$$= -\frac{3\pi}{2} \left(\frac{a^4}{4} - \frac{a^4}{2} \right)$$

$$Q = \underline{\underline{\frac{3\pi a^4}{8}}}$$

$$\tau_w = \mu \nabla u \cdot n$$

Cylinder coordinates

$$\tau_w = \mu \frac{\partial u_z}{\partial r} \Big|_{r=a} = -\mu \frac{3}{4} (2r) \Big|_{r=a} = -\mu \underline{\underline{\frac{3a}{2}}}$$

(9)

$$7f) \quad \nabla^2 u_z = -\beta \left(= \frac{1}{\mu} \frac{dP_0}{dz} \right)$$

Multiply by u_z and integrate over A

$$\int u_z \nabla^2 u_z dA = \int -\beta u_z dA$$

$$\text{Gauss theorem} \left(\nabla \cdot (u_z \nabla u_z) - \nabla u_z \cdot \nabla u_z \right) dA = -\beta \int u_z dA$$

$$\downarrow \int u_z \nabla u_z d\Gamma - \int \nabla u_z \cdot \nabla u_z dA = -\beta Q$$

$$- \int \nabla u_z \cdot \nabla u_z dA = \frac{1}{\mu} \frac{dP_0}{dz} Q$$

$$\text{Use } \Phi = \mu \nabla u_z \cdot \nabla u_z$$

$$\Rightarrow - \int \Phi dA = \frac{dP_0}{dz} Q$$

Interpretation - Effect of pressure work goes to dissipation.

Problem 2

a) Navier - Stokes (neglecting body forces)

$$\frac{\partial u}{\partial t} + (u \cdot \nabla)u = -\frac{1}{\rho} \nabla p + \nu \nabla^2 u$$

$$\nabla \cdot u = 0$$

Reynolds decomposition

$$u = \bar{u} + u'$$

$$p = \bar{p} + p'$$

where overline represents an ensemble average and prime is fluctuation about the average.

Take average of momentum eq and use commutation of averaging and differentiation: = div. eq.

$$\frac{\partial \bar{u}}{\partial t} + \overline{(u \cdot \nabla)u} = -\frac{1}{\rho} \nabla \bar{p} + \nu \nabla^2 \bar{u}$$

$$\nabla \cdot \bar{u} = 0$$

Rewrite convection term

$$\overline{(u \cdot \nabla)u} = \nabla \cdot \overline{uu} - \overline{u \nabla u}$$

where uu is the outer product of u times u

Index notation:

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$$\overline{u_j \frac{\partial u_i}{\partial x_j}} = \frac{\partial \overline{u_i u_j}}{\partial x_j} - \overline{u_i \frac{\partial u_j}{\partial x_j}}$$

Introduce Reynolds decomposition:

$$\overline{\nabla \cdot u u} = \overline{\nabla \cdot (\bar{u} + u')(\bar{u} + u')}$$

$$= \overline{\nabla \cdot (\bar{u} \bar{u} + u' u')}$$

$$= (\bar{u} \cdot \nabla) \bar{u} + \nabla \cdot \overline{u' u'}$$

=> Momentum eq:

$$\frac{\partial \bar{u}}{\partial t} + (\bar{u} \cdot \nabla) \bar{u} = -\frac{1}{\rho} \nabla \bar{p} + \nu \nabla^2 \bar{u} - \nabla \cdot \overline{u' u'}$$

where $\overline{u' u'}$ is the Reynolds stress tensor.

b) Convection is zero since u is aligned with the walls. Stationary $\Rightarrow \frac{\partial \bar{u}}{\partial t} = 0$

Momentum eq:

X-dir

$$0 = -\frac{1}{\rho} \frac{\partial \bar{p}}{\partial x} + \nu \frac{\partial^2 \bar{u}}{\partial y^2} - \frac{\partial \overline{u' u'}}{\partial y}$$

$$0 = -\frac{\partial \bar{p}}{\partial x} + \mu \frac{\partial}{\partial y} \left(\frac{\partial \bar{u}}{\partial y} - \rho \overline{u' u'} \right)$$

only change in y
 \downarrow $u' =$ vel. fluct.
in y -dir

y-dir

$$0 = -\frac{1}{\rho} \frac{\partial \bar{P}}{\partial y} - \frac{\partial \overline{u'u'}}{\partial x} - \frac{\partial \overline{v'u'}}{\partial y} - \frac{\partial \overline{w'u'}}{\partial z}$$

$$\frac{\partial \bar{P}}{\partial y} = \rho \frac{\partial \overline{v'u'}}{\partial y}$$

z-dir

$$0 = -\frac{1}{\rho} \frac{\partial \bar{P}}{\partial z} - \frac{\partial \overline{w'u'}}{\partial x} - \frac{\partial \overline{w'u'}}{\partial y} - \frac{\partial \overline{w'u'}}{\partial z}$$

all zero

$\overline{w'u'}$ zero because statistics are homogeneous in x and z-dir

$$\frac{\partial \bar{P}}{\partial z} = 0$$

$$\bar{P} = P(x, y)$$

$$\frac{\partial \bar{P}}{\partial y} = \rho \frac{\partial \overline{v'u'}}{\partial y}$$

$$\bar{P} = \rho \int_{-h}^y \overline{v'u'} dy + P_0(x)$$

$$\bar{P} = \rho \overline{v'u'}(y) + P_0(x)$$

From x-dir.

$$\frac{d\bar{P}}{dx} = \mu \frac{d}{dy} \left(\frac{d\bar{u}}{dy} - \beta \overline{u'v'} \right)$$

} depends only on y

=> $\frac{\partial \bar{P}}{\partial x}$ is constant

$$\frac{\partial \bar{P}}{\partial x} = \frac{\partial P_0}{\partial x} = \text{constant}$$

$$P_0 = \beta \cdot x + d$$

$$\underline{\bar{P}(x,y) = \beta \overline{u'v'} + \beta x + d}$$

c) Very close to the wall the only non-zero term in the momentum equation is the Laplacian:

$$\frac{d\bar{P}}{dx} = \mu \frac{d^2 \bar{u}}{dy^2}$$

$$\frac{d^2 \bar{u}}{dy^2} = -\beta \left(= -\frac{1}{\mu} \frac{d\bar{P}}{dx} \right) = \text{constant}$$

Consider y as the distance to the wall located at y=0.

Hence $\bar{u}(y) = -\frac{1}{2} 3y^2 + 4y + c_2$

As $y \rightarrow 0$ $\bar{u}(y) \approx 4y$

The wall shear stress $\tau_w = \mu \frac{d\bar{u}}{dy}$

$$\bar{u} = \frac{1}{\mu} \int_0^y \tau_w dy$$

$$\bar{u} = \frac{\tau_w y}{\mu} + \cancel{u_0}$$

Problem 3

a) $u \frac{\partial u}{\partial x} = \frac{u^*}{LB} \frac{\partial \frac{u^*}{LB}}{\partial \frac{x^*}{L}} = \frac{L u^*}{L^2 B^2} \frac{\partial u^*}{\partial x^*} = \frac{u^*}{LB^2} \frac{\partial u^*}{\partial x^*}$

=> X-momentum has been multiplied by

$$\frac{1}{LB^2}$$

=> $\frac{1}{\rho} \frac{\partial P^*}{\partial x^*} \cdot \frac{1}{LB^2} = \frac{1}{\rho} \frac{\partial P^*}{\partial \frac{x^*}{L}} \frac{1}{L^2 B^2} = \frac{\partial \frac{P^*}{\rho L^2 B^2}}{\partial \frac{x^*}{L}} = \frac{\partial P}{\partial x}$

=> $P = \frac{P^*}{\rho L^2 B^2}$ and $\hat{P} = \rho L^2 B^2$

b) $u = -F(y)$

Continuity :

$$\frac{\partial u(x,y)}{\partial x} + \frac{\partial u(y)}{\partial y} = 0$$

$$\frac{\partial u}{\partial x} = F'$$

Integrate

$$u(x,y) = F' \int_0^x dx$$

$$u(x,y) = F' x + C_0(y)$$

Symmetry:

$$u(0,y) = F'(y) \cdot 0 + C_0(y) = 0$$

$$\Rightarrow C_0 = 0$$

$$\underline{u(x,y) = F'(y) x}$$

Boundary conditions

$$y=0 \begin{cases} u=0 & \Rightarrow F'(0) = 0 \\ u=0 & \Rightarrow F(0) = 0 \end{cases}$$

$$y \rightarrow \infty \quad u \rightarrow u_{outer}$$

$$u_{outer}^x = Bx^* = BLx$$

$$u_{outer} = \frac{BLx}{BL} = x$$

$$u = F'(\infty)x = x$$

$$\Rightarrow \underline{F'(\infty) = 1}$$

3c) Substitute $u = F'x$
 $u = -F$

into y-eq

$$\gamma \left(u \frac{\partial u}{\partial x} + u \frac{\partial u}{\partial y} \right) = - \frac{\partial p}{\partial y} + \gamma \left(\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} \right)$$

$$\frac{\partial u}{\partial x} = 0, \quad \frac{\partial u}{\partial y} = -F', \quad \frac{\partial^2 u}{\partial x^2} = 0, \quad \frac{\partial^2 u}{\partial y^2} = -F''$$

$$\gamma \left(F'x \cdot 0 + (-F)(-F') \right) = - \frac{\partial p}{\partial y} + \gamma \left(0 + (-F'') \right)$$

$$\gamma FF' = - \frac{\partial p}{\partial y} - \gamma F''$$

$$\frac{\partial p}{\partial y} = - \gamma \left(F'' + FF' \right)$$

3c) Integrate y

$$P = -\gamma \int_0^y F'' + FF' dy + P_0$$

$$P = -\gamma \left(F'(y) - \cancel{F'(0)} + \int_0^y FF' dy \right) + P_0$$

$$\int_0^y FF' dy = F^2 / 2 - \int_0^y F' F dy$$

$$\Rightarrow \int_0^y FF' dy = \frac{1}{2} F^2$$

$$P = -\gamma \left(F'(y) + \frac{1}{2} F^2(y) \right) + P_0$$

↑
1/2 missing in assignment?

3d)

X-component

$$u \frac{du}{dx} + v \frac{du}{dy} = -\frac{\partial p}{\partial x} + \frac{\partial^2 u}{\partial y^2} + \gamma \frac{\partial^2 u}{\partial x^2}$$

$$F' \times F' + (-F)F'' \times = -\frac{\partial p}{\partial x} + F''' \times + \gamma \cdot 0$$

$$\frac{dp}{dx} = X (F''' + FF'' - F'^2)$$

⇒

$$F''' + FF'' - F'^2 = \alpha = \text{constant}$$

As $y \rightarrow \infty$ $F'(\infty) = 1$ (constant)

and $F''(\infty) = F'''(\infty) = 0$

$$\Rightarrow \alpha = -1$$
